

Radiative seesaw and Dark Matter



**A two-loop Radiative seesaw
with multi-component DM
explaining
the Gamma Excess
in the Higgs Decay and at the Fermi LAT**

arXiv:

ON JAN.13 AT THE UNIV. OF TOYAMA

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KANAZAWA UNIVERSITY(KU)**

**IN COLLABORATION WITH
MAYUMI AOKI
(KU&MPI,HEIDELBERG)
HIROSHI TAKANO(KU)**

PLAN

I A brief review on radiative seesaw models and DM

II A two-loop model with two or three DM

III By products:

*** 135 gamma-ray line at the Fermi LAT**

*** Enhancement of $h \rightarrow$ two gammas**

IV Conclusion

I Radiative Seesaw and DM

Radiative Neutrino Mass Generation

(ZEE,80;86; WOLFENSTEIN,80;
BABU,88, ETC)



Radiative Seesaw Mechanism

(KRAUS, NASRI + TRODDEN,02; MA,06; AOKI, KANEMURA+
SETO,08; ETC)

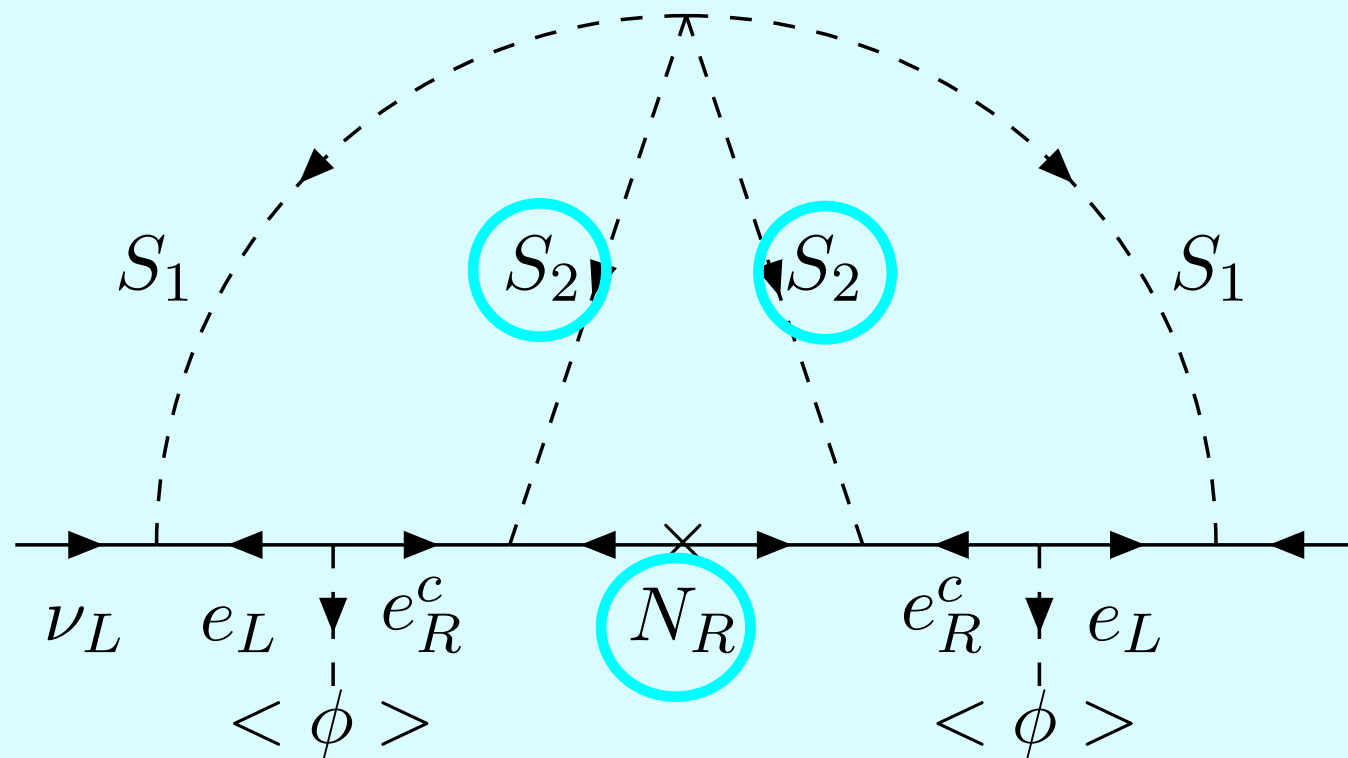
**Unbroken discrete symmetry to
forbid the Dirac masses**



N_R , Inert Higgs, etc are DM candidates.

Three-loop Model

KRAUS, NASRI + TRODDEN, 02



Z_2 -ODD SM SINGLET

N_R CAN BE DM.

$$\mathbf{m}_\nu = \begin{pmatrix} 0 & 0 \\ 0 & M \end{pmatrix} \rightarrow \begin{pmatrix} m^2/M & 0 \\ 0 & M \end{pmatrix}$$

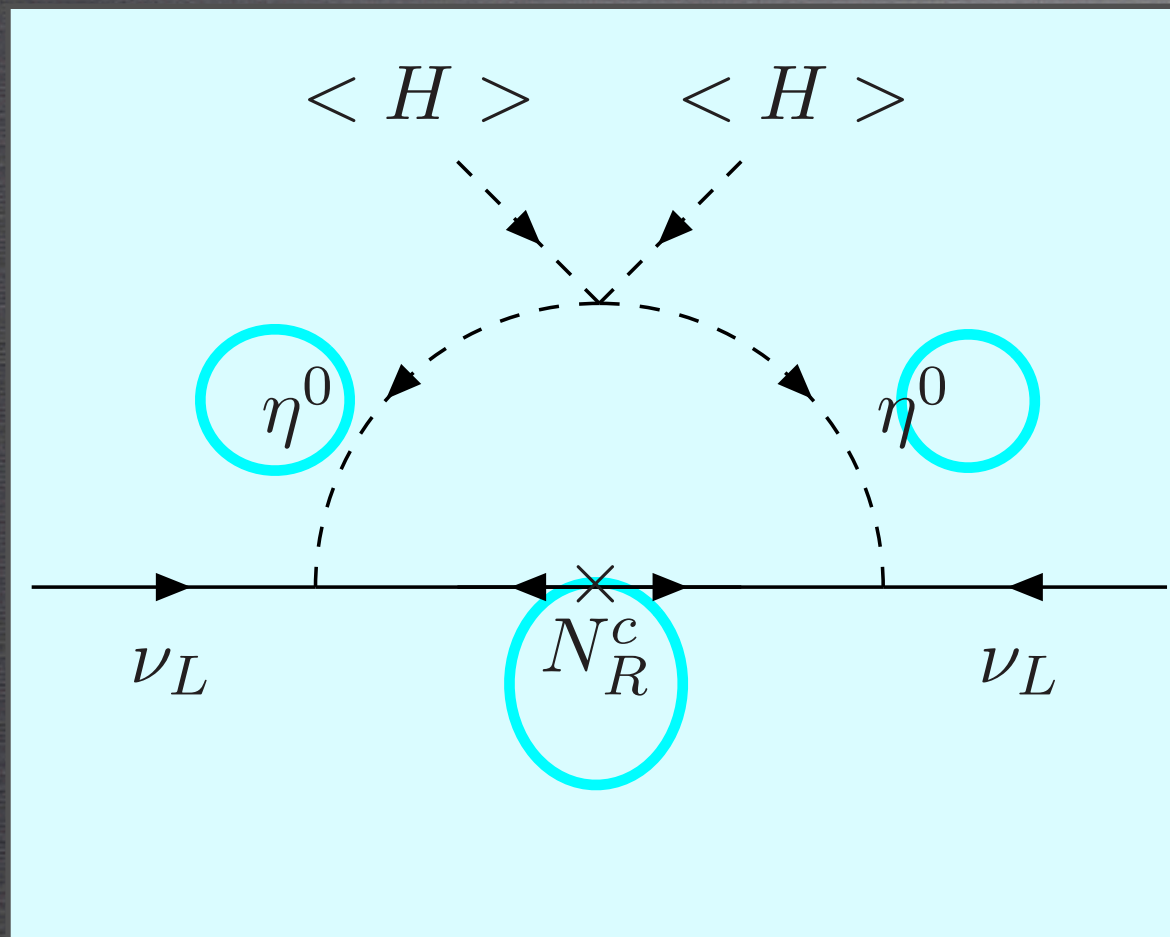
Radiative Seesaw

One-loop Model

MA, 06

Z₂-ODD

INERT SU(2) DOUBLET
HIGGS



$$\langle \eta \rangle = 0$$

N_R and eta are DM candidates.

Eta DM studied by

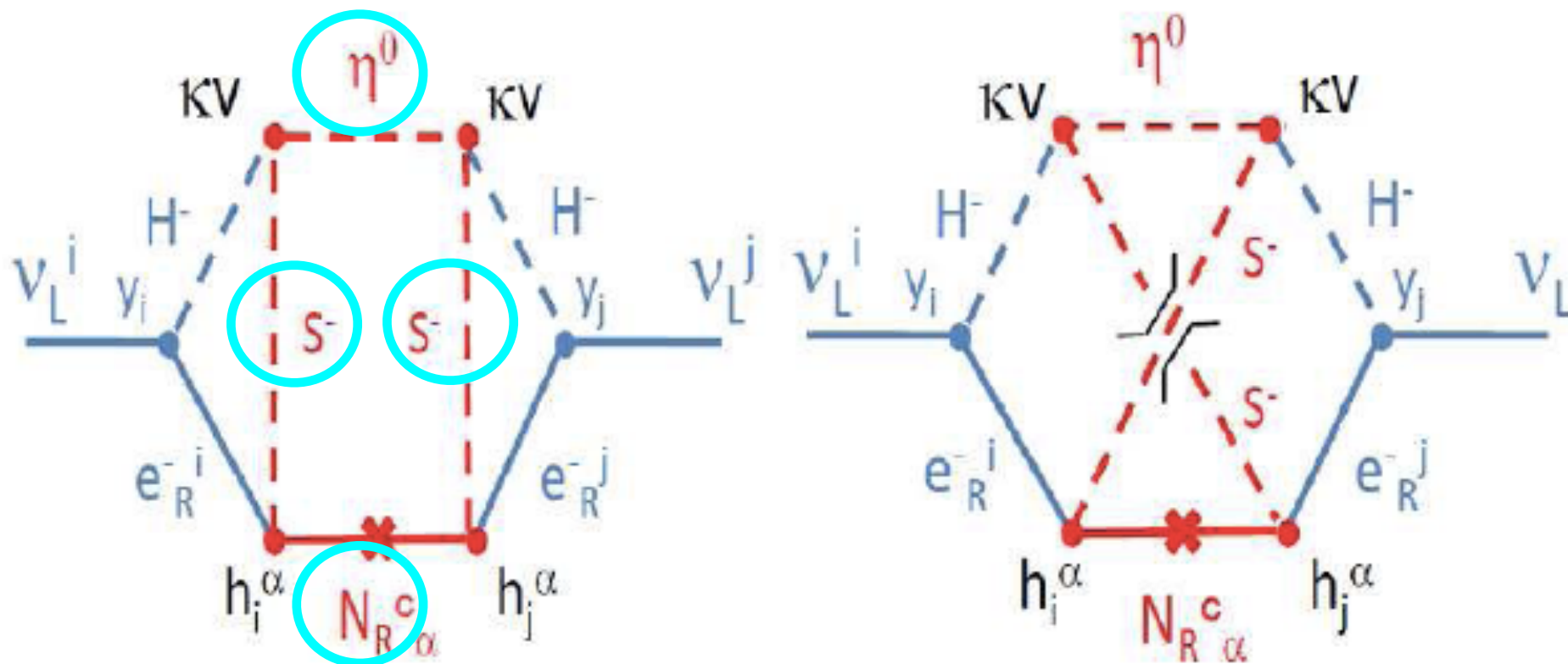
BARBIERI, HALL + RYCHKOV, 06;
LOPEZ, OLIVER + TYTGAT, 06;
DOLLE + SU, 09 ETC

N_R DM studied by

KUBO, MA + SUEMATSU, 06;
ARISTIZABAL SIERRA, KUBO, RESTREPO,
SUEMATSU + ZAPATA, 08;
GELMINI, OSOBA + PALOMARES-RUIZ, 09, ETC

Three-loop Model

AOKI, KANEMURA+SETO, 08



Z_2 ODD

INERT $SU(2)$ DOUBLET
HIGGS

$$\langle \eta \rangle = 0$$

	Q^i	u_R^i	d_R^i	L^i	e_R^i	Φ_1	Φ_2	$S^\pm$	η	N_R^α
Z_2 (exact)	+	+	+	+	+	+	+	-	-	-
$\tilde{Z}_2$ (softly broken)	+	-	-	+	+	+	-	+	-	+

NON-SUSY MODELS OF RADIATIVE GENERATION OF m_ν

Model	L-violating dim.	No. of loops	ν_R	No. of stable DMs
[1]	3 (tri-linear scalar coupling)	1	No	0
[2, 3, 21]	3 (tri-linear scalar coupling)	2	No	0
[15]	3 (tri-linear scalar coupling)	1	No	1
[22]	3 (tri-linear scalar coupling))	2	Yes	2
[4, 5, 11]	3 (Majorana mass)	3	Yes	1
[6–9, 12, 14, 23, 24]	3 (Majorana mass)	1	Yes	1
[13]	3 (Majorana mass)	1	Yes ($SU(2)_L$ triplet)	1
[17]	4 (quartic scalar coupling)	1	No	1
[19]	No L-violation	1	Yes	0
[18]	Spontaneous violation	2	No	1
[20]	Sponataneos violation	2	Yes	1
[16]	Spontaneous violation	1	Yes	1
[10]	2 (scalar mass)	1	Yes	0
Our model	2 (scalar mass)	2	Yes	2 or 3

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- [4] L. M. Krauss, S. Nasri and M. Trodden, Phys. Rev. D **67** (2003) 085002 [hep-ph/0210389].
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- [17] M. Aoki, S. Kanemura and K. Yagyu, Phys. Lett. B **702** (2011) 355 [Erratum-ibid. B **706** (2012) 495] [arXiv:1105.2075 [hep-ph]].
- [18] M. Lindner, D. Schmidt and T. Schwetz, Phys. Lett. B **705** (2011) 324 [arXiv:1105.4626 [hep-ph]].
- [19] S. Kanemura, T. Nabeshima and H. Sugiyama, Phys. Lett. B **703** (2011) 66 [arXiv:1106.2480 [hep-ph]].
- [20] S. Kanemura, T. Nabeshima and H. Sugiyama, Phys. Rev. D **85** (2012) 033004 [arXiv:1111.0599 [hep-ph]].
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THE MODEL

Model	L-violating dim.	No. of loops	ν_R	No. of stable DMs
Our model	2 (scalar mass)	2	Yes	2 or 3

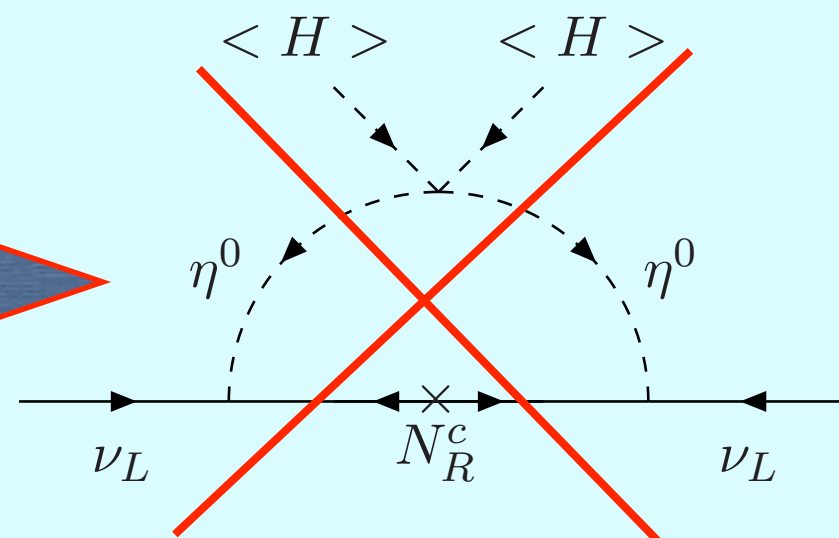
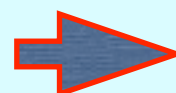
MA

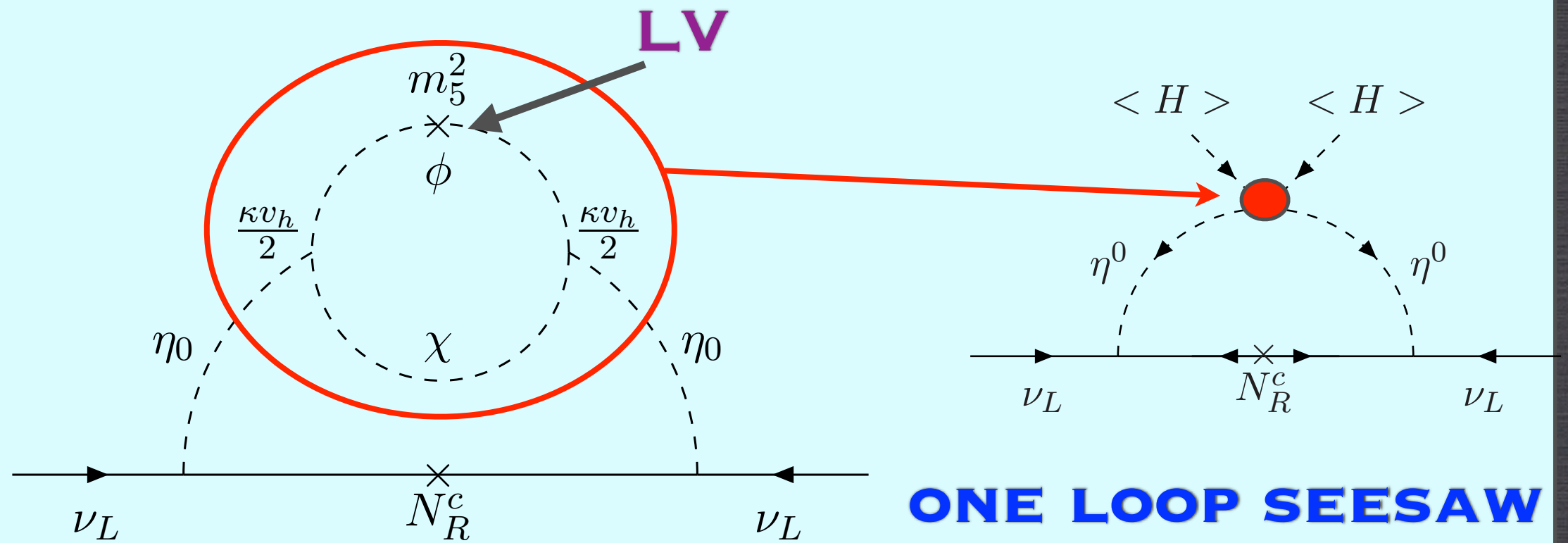
field	statistics	$SU(2)_L$	$U(1)_Y$	L	Z_2	Z'_2	D_{2N}
(ν_L, l)	F	2	$-1/2$	1	+	+	1
l^c	F	1	1	1	+	+	1
N_R^c	F	1	1	0	-	+	1''
$H = (H^+, H^0)$	B	2	$1/2$	0	+	+	1
$\eta = (\eta^+, \eta^0)$	B	2	$1/2$	-1	-	+	1''
χ	B	1	0	0	+	-	1'
ϕ	B	1	0	1	-	-	1'''

MATTER CONTENT (MINIMAL EXTENSION)

NO

“ λ_5 term”, $(1/2)\lambda_5(H^\dagger\eta)^2$





TWO LOOP SEESAW

$$m_\nu \sim \left(\frac{\kappa}{16\pi^2} \right)^2 \frac{m_D^2}{M_R} \sim 10^{-7} \frac{m_D^2}{M_R} \text{ for } \kappa = 0.1$$

$M_R \sim \mathbf{1 \text{ TEV IS A NATURAL SCALE.}}$

THE Z2 x Z2 x **L** INVARIANT QUARTIC COUPLINGS

$$\begin{aligned} V_\lambda = & \lambda_1(H^\dagger H)^2 + \lambda_2(\eta^\dagger \eta)^2 + \lambda_3(H^\dagger H)(\eta^\dagger \eta) + \lambda_4(H^\dagger \eta)(\eta^\dagger H) \\ & + \frac{1}{4}\gamma_1\chi^4 + \gamma_2(H^\dagger H)\chi^2 + \gamma_3(\eta^\dagger \eta)\chi^2 + \gamma_4|\phi|^4 + \gamma_5(H^\dagger H)|\phi|^2 \\ & + \gamma_6(\eta^\dagger \eta)|\phi|^2 + \gamma_7\chi^2|\phi|^2 + \frac{\kappa}{2}[(H^\dagger \eta)\chi\phi + h.c.] . \end{aligned}$$

THE Z2 x Z2 INVARIANT SCALAR MASS TERMS

$$V_m = m_1^2 H^\dagger H + m_2^2 \eta^\dagger \eta + \frac{1}{2}m_3^2 \chi^2 + m_4^2 |\phi|^2 + \frac{1}{2}m_5^2 [\phi^2 + (\phi^*)^2]$$

BREAKS L.

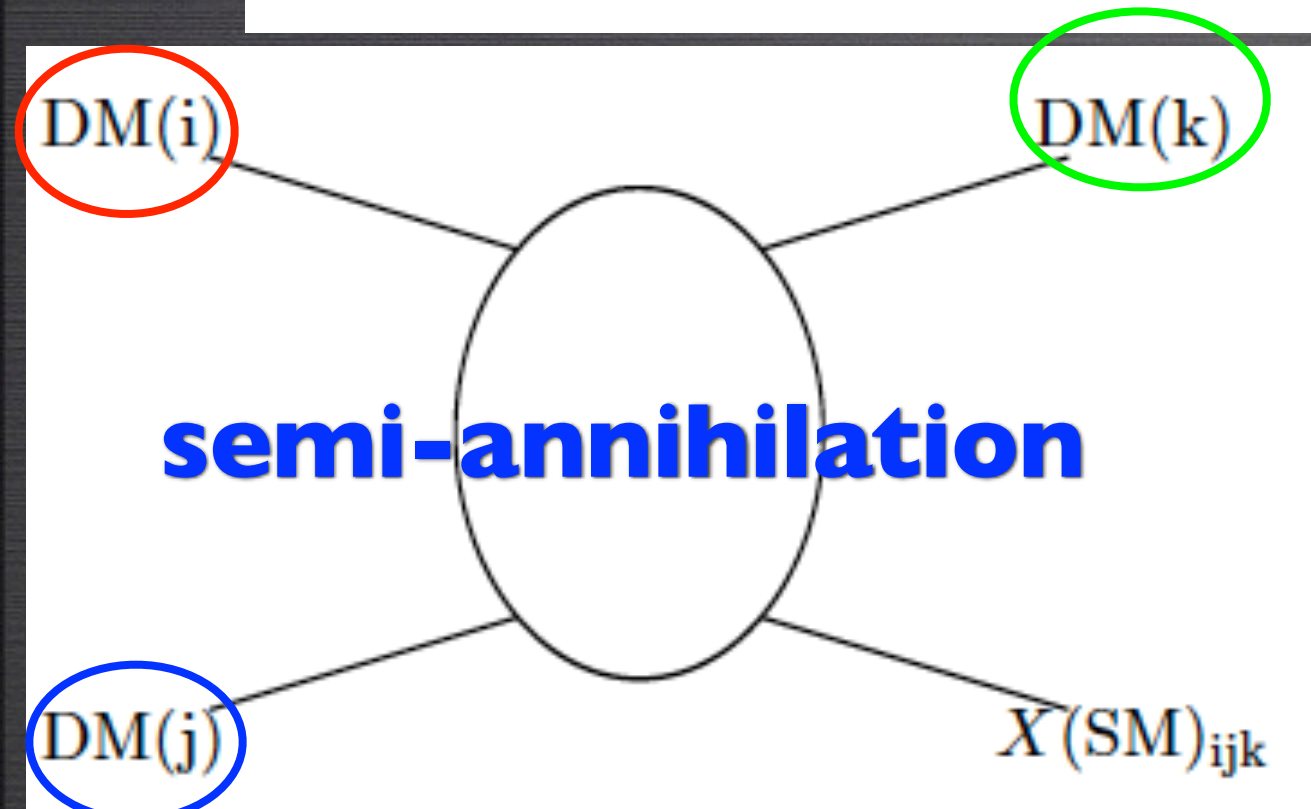
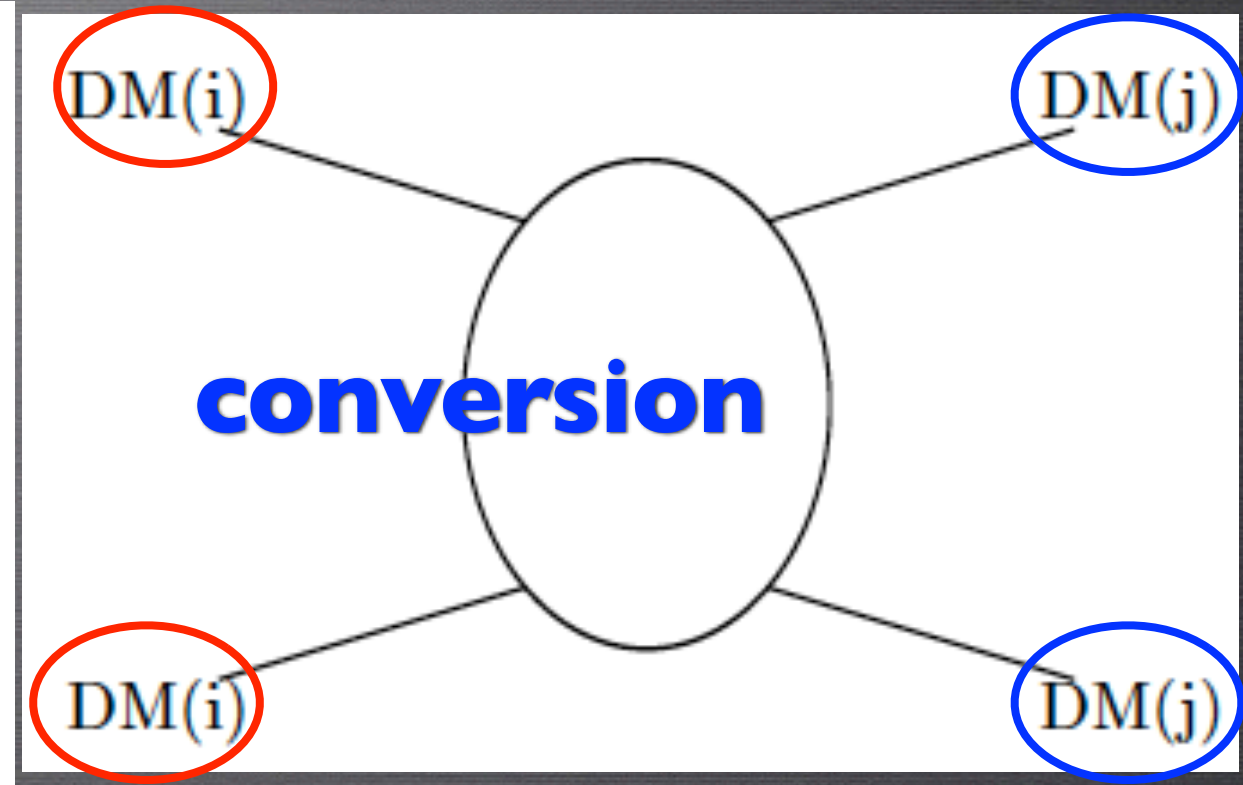
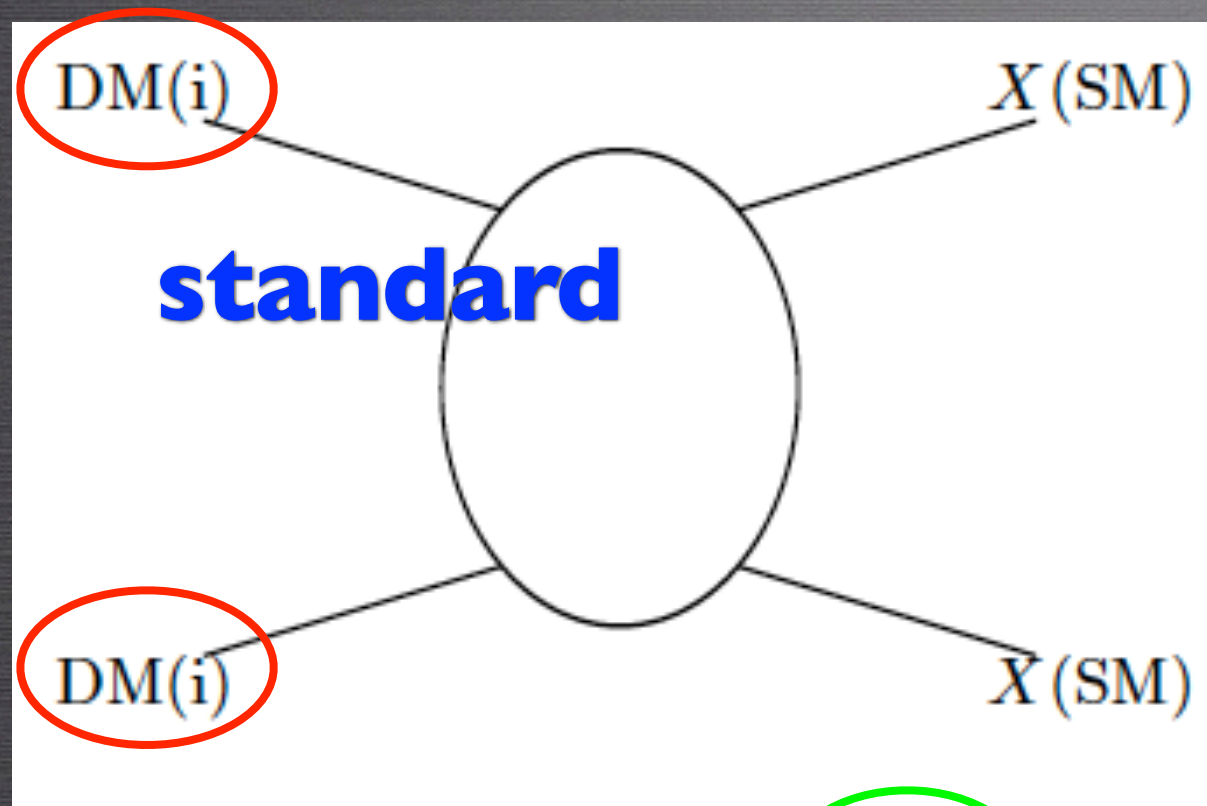
COLD DARK MATTER CANDIDATES

field	statistics	$SU(2)_L$	$U(1)_Y$	L	Z_2	Z'_2	D_{2N}
(ν_L, l)	F	2	$-1/2$	1	+	+	1
l^c	F	1	1	1	+	+	1
N_R^c	F	1	1	0	-	+	1''
$H = (H^+, H^0)$	B	2	$1/2$	0	+	+	1
$\eta = (\eta^+, \eta^0)$	B	2	$1/2$	-1	-	+	1''
χ	B	1	0	0	+	-	1'
ϕ	B	1	0	1	-	-	1'''

OF DMS=3: (η, χ, ϕ) **OR** (N_R^c, χ, ϕ)

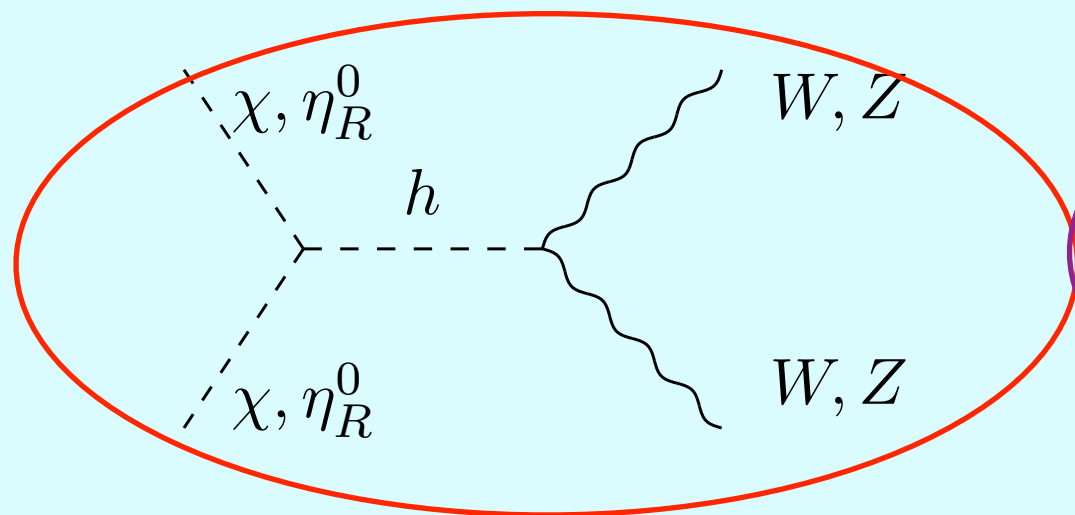
OF DMS=2: (η, χ) **OR** (χ, ϕ) **ETC**

Non-standard annihilations and relic abundance

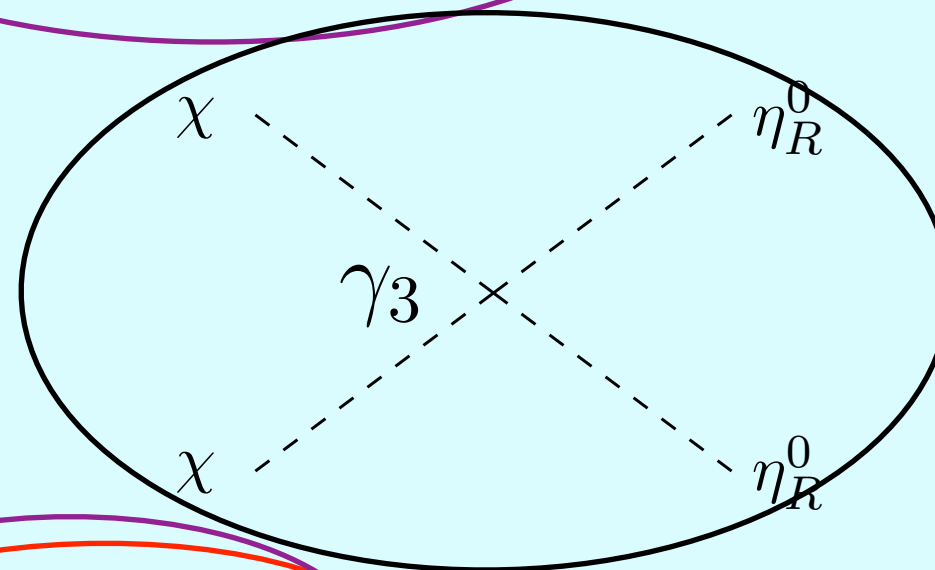
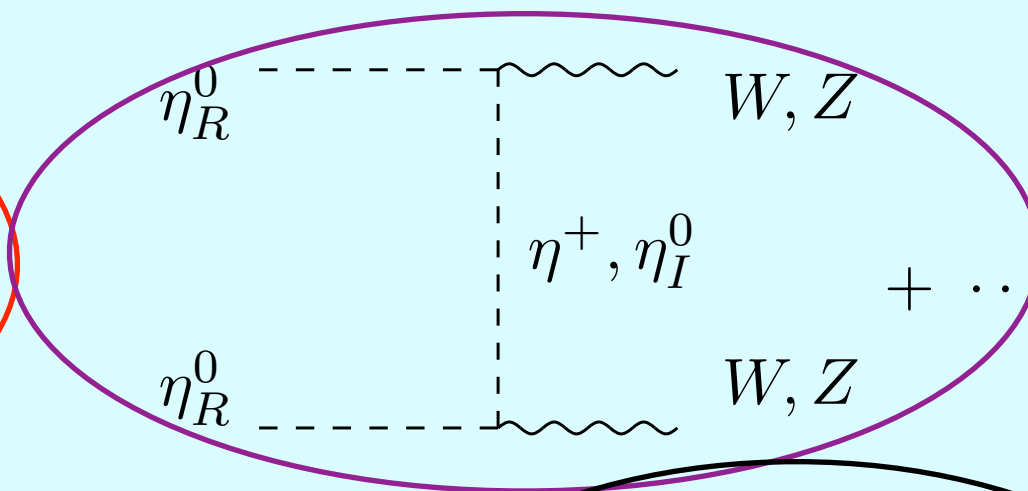


SEE FOR MORE DETAILS:
AOKI, DUERR, KUBO
+TAKANO,
ARXIV:1207.3318

$\eta - \chi$ DM SYSTEM



$$(2\gamma_2, \lambda_3)v_h$$



BOLTZMANN EQS.

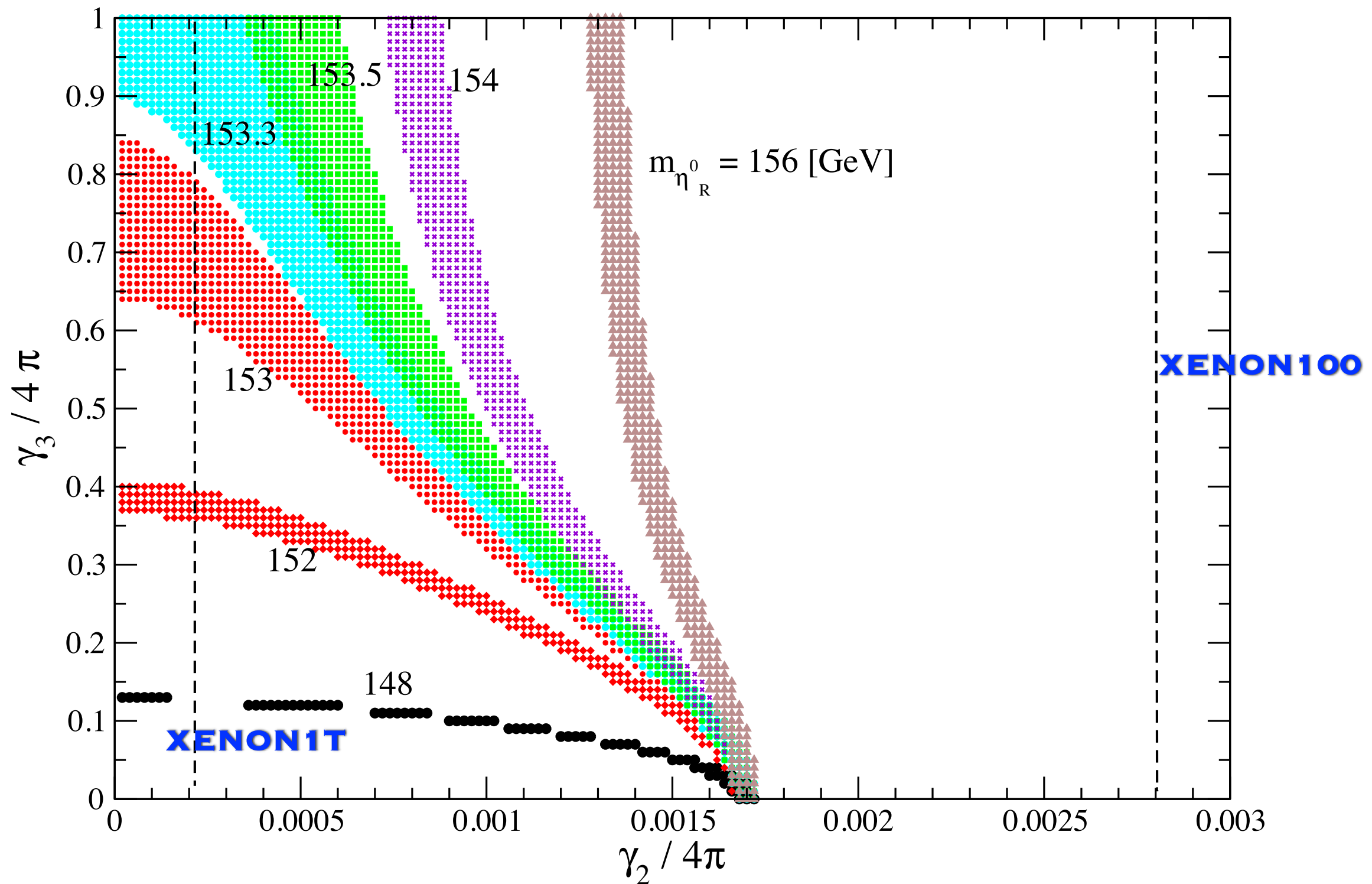
$$\frac{dY_{\eta_R^0}}{dx} = -0.264 g_*^{1/2} \left[\frac{\mu M_{\text{PL}}}{x^2} \right] \left\{ \langle \sigma(\eta_R^0 \eta_R^0; \text{SM}) v \rangle \left(Y_{\eta_R^0} Y_{\eta_R^0} - \bar{Y}_{\eta_R^0} \bar{Y}_{\eta_R^0} \right) \right.$$

$$\left. + \langle \sigma(\eta_R^0 \eta_R^0; \chi\chi) v \rangle \left(Y_{\eta_R^0} Y_{\eta_R^0} - \frac{Y_\chi Y_\chi}{\bar{Y}_\chi \bar{Y}_\chi} \bar{Y}_{\eta_R^0} \bar{Y}_{\eta_R^0} \right) \right\} ,$$

$$\frac{dY_\chi}{dx} = -0.264 g_*^{1/2} \left[\frac{\mu M_{\text{PL}}}{x^2} \right] \left\{ \langle \sigma(\chi\chi; \text{SM}) v \rangle \left(Y_\chi Y_\chi - \bar{Y}_\chi \bar{Y}_\chi \right) \right.$$

$$\left. + \langle \sigma(\eta_R^0 \eta_R^0; \chi\chi) v \rangle \left(Y_{\eta_R^0} Y_{\eta_R^0} - \frac{Y_\chi Y_\chi}{\bar{Y}_\chi \bar{Y}_\chi} \bar{Y}_{\eta_R^0} \bar{Y}_{\eta_R^0} \right) \right\} ,$$

$$\Omega_T h^2 (\simeq \Omega_\chi h^2) = 0.1157 \pm 0.0046 \quad (2\sigma) \quad \textbf{WITH} \quad m_\chi = 135 \text{ GeV}$$

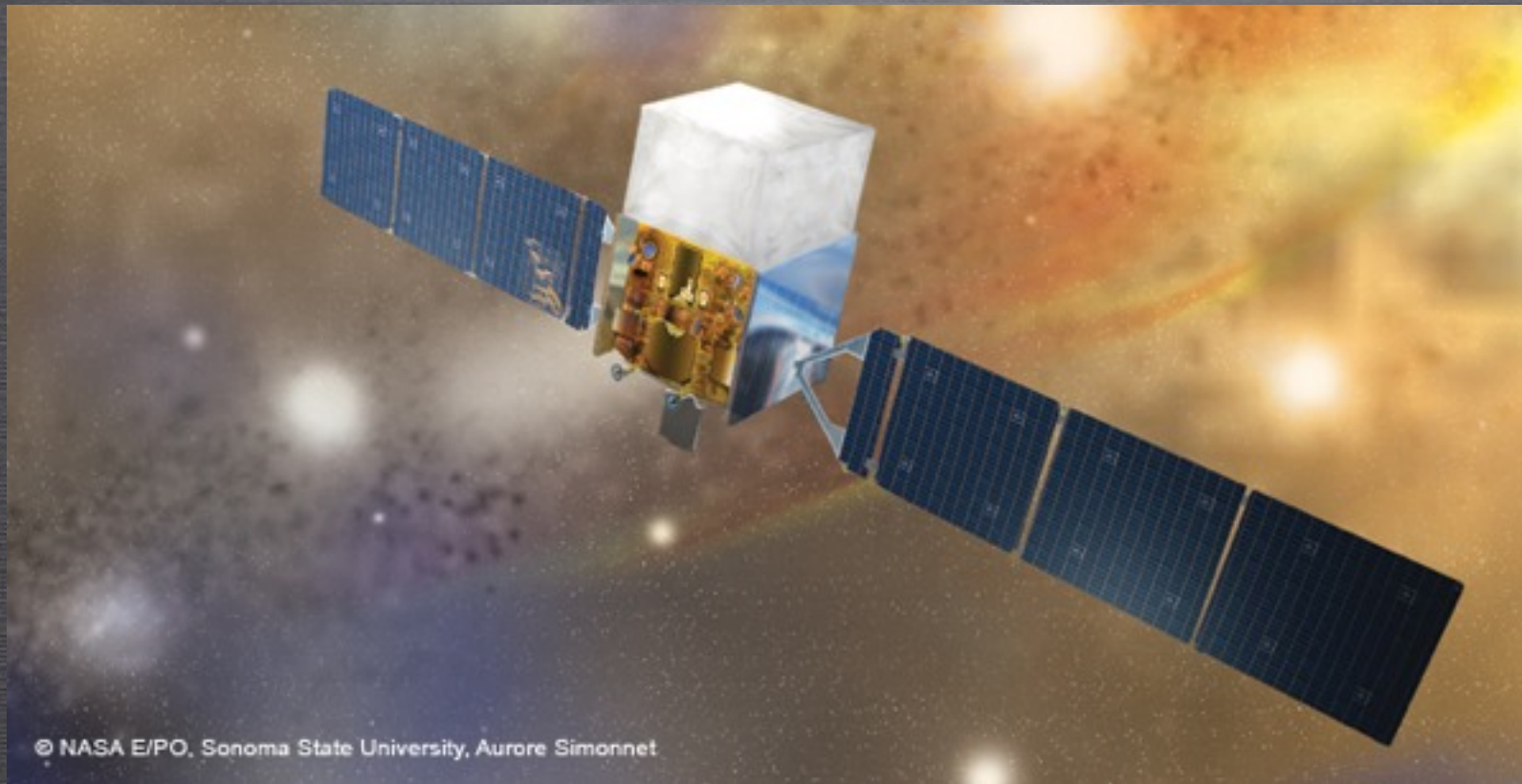


$$\lambda_1 = 0.129, \quad \lambda_3 = -1.26, \quad \lambda_4 = -0.0205$$

III BONUS

*** THE FIRST BONUS:**

FERMI-LAT 135 GEV GAMMA-RAY LINE



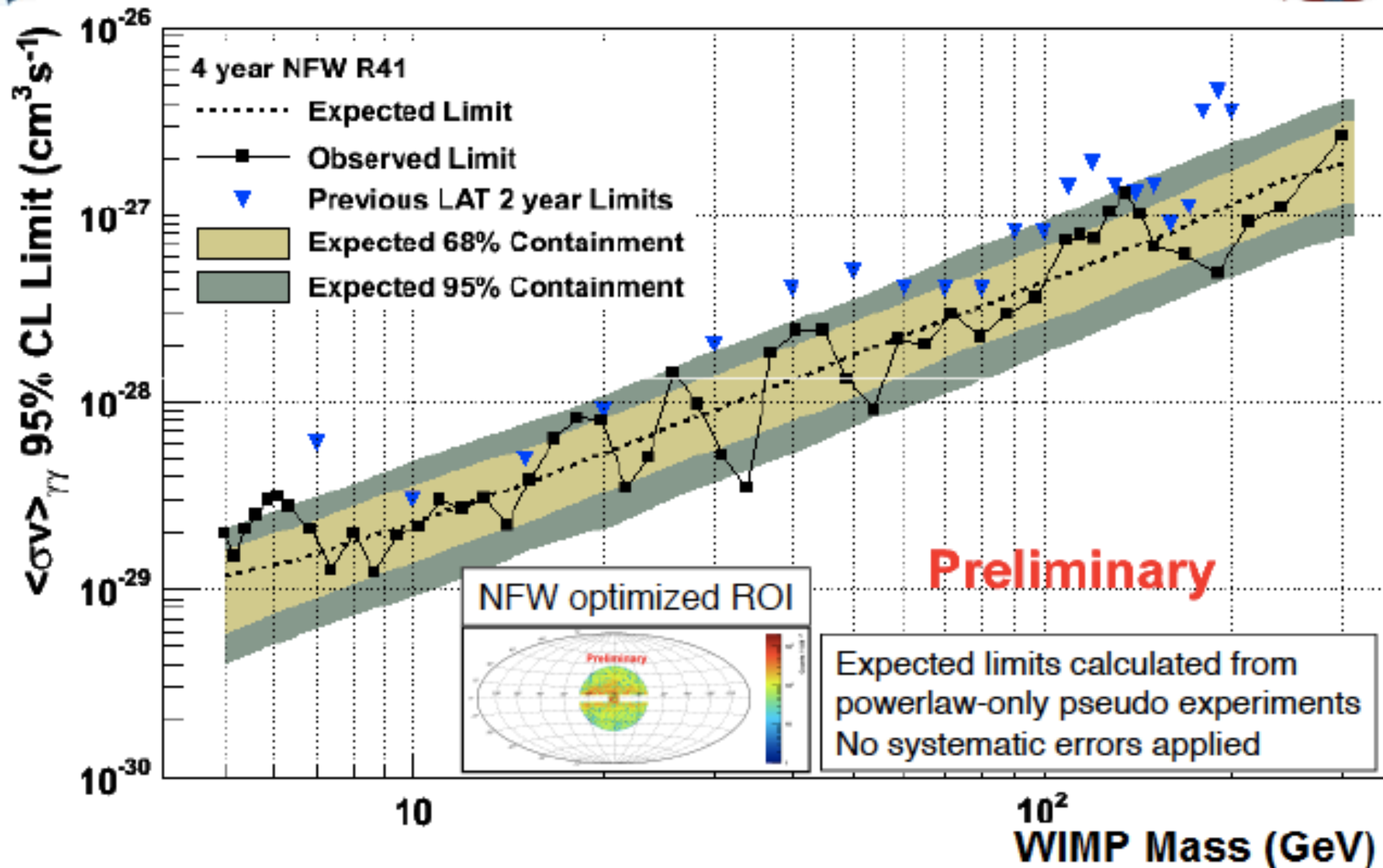
**FERMI-LAT IS MEASURING GAMMA RAYS COMING
FROM THE UNIVERSE SINCE 2008.**

LAT COLLABORATION, ARXIV:1205.2739; ETC

IF THE DATA FROM THE WHOLE SKY IS INCLUDED:



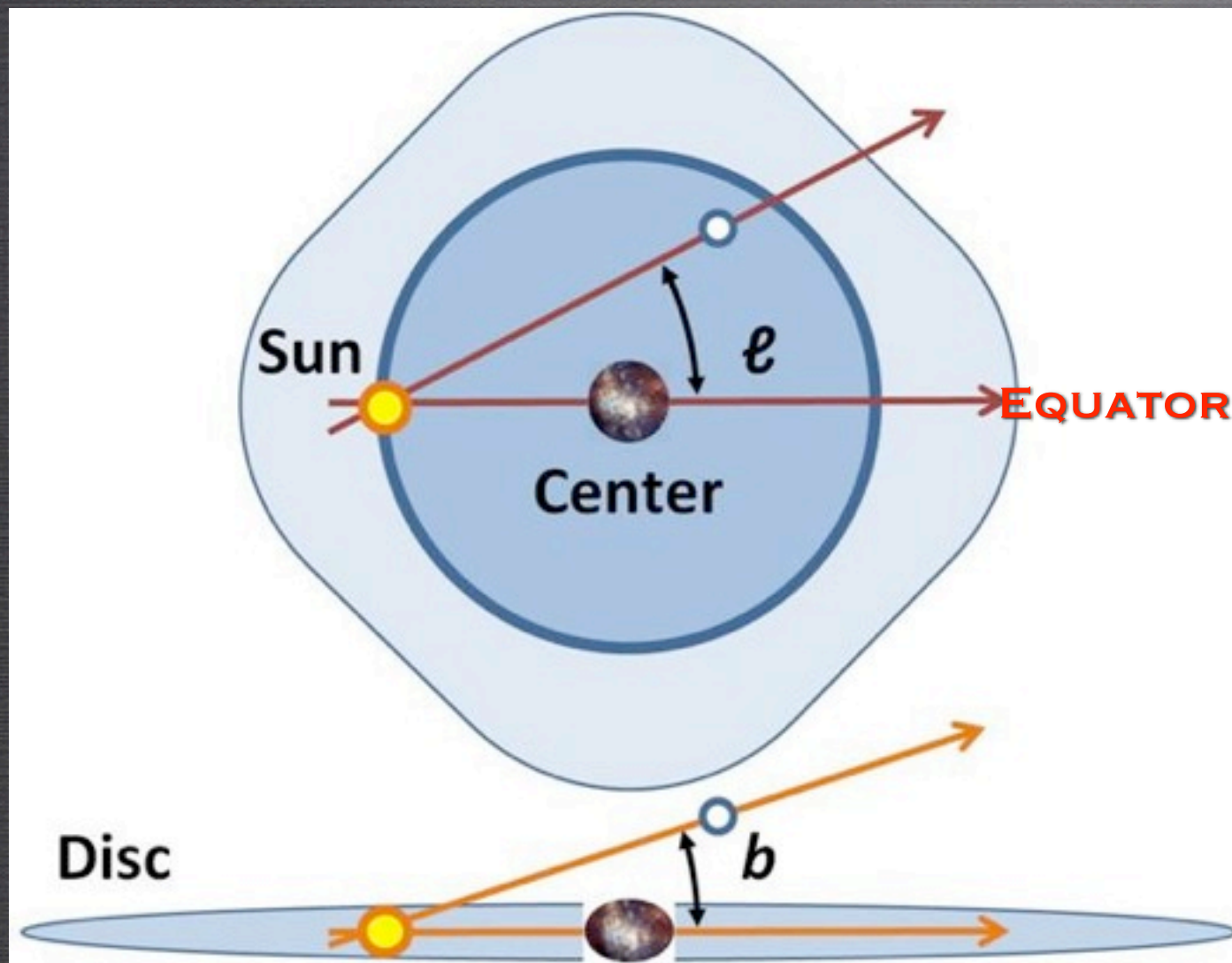
95% CL $\langle\sigma v\rangle_{\gamma\gamma}$ NFW Upper Limit R41



HOWEVER , IF YOU LOOK AT THE CENTER OF THE GALAXY, SOMETHING MORE INTERESTING IS GOING ON.

WENIDER, ARXIV:1204.2797;

BRINGMANN ET AL, ARXIV:1203.1312; ETC

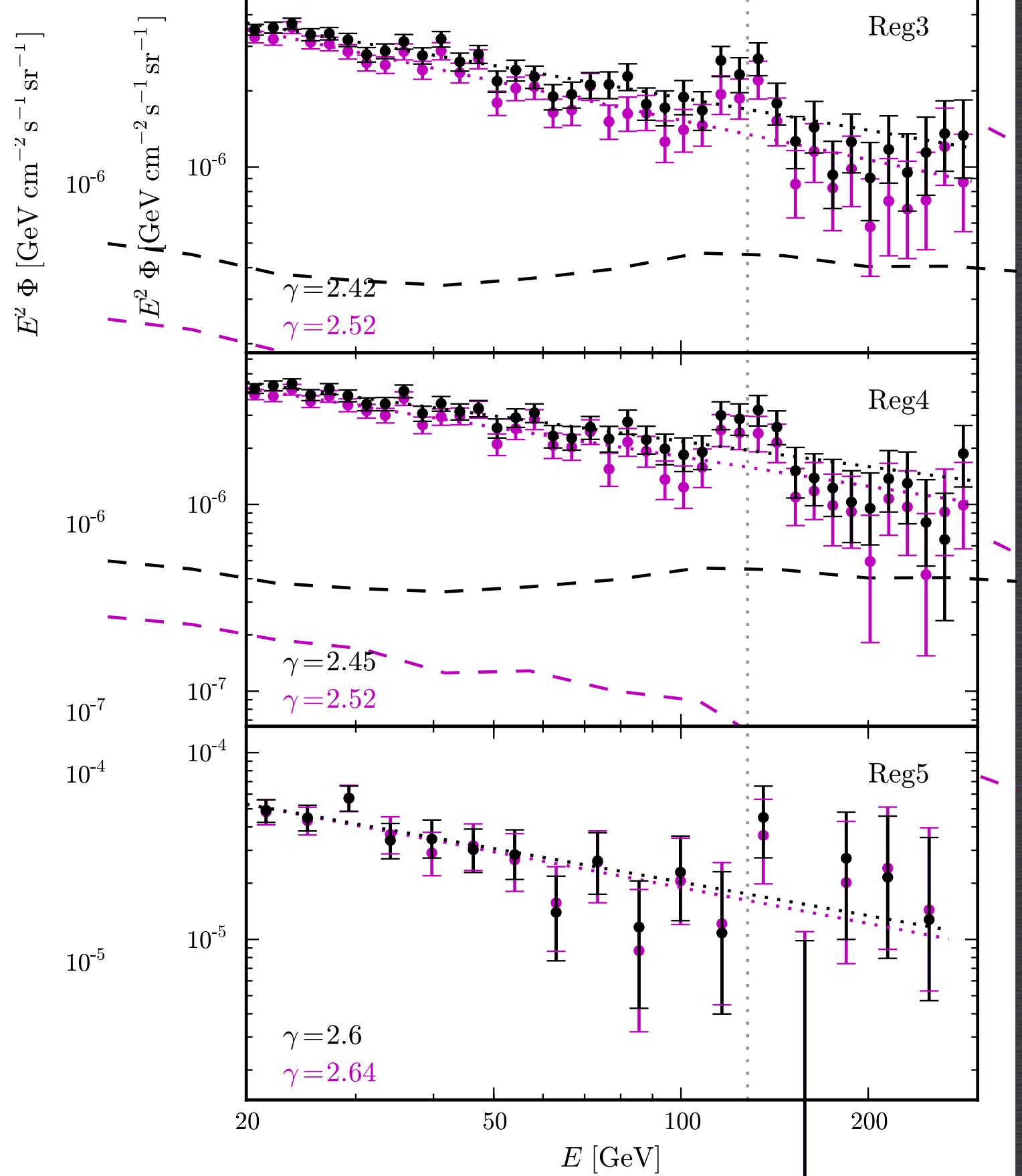
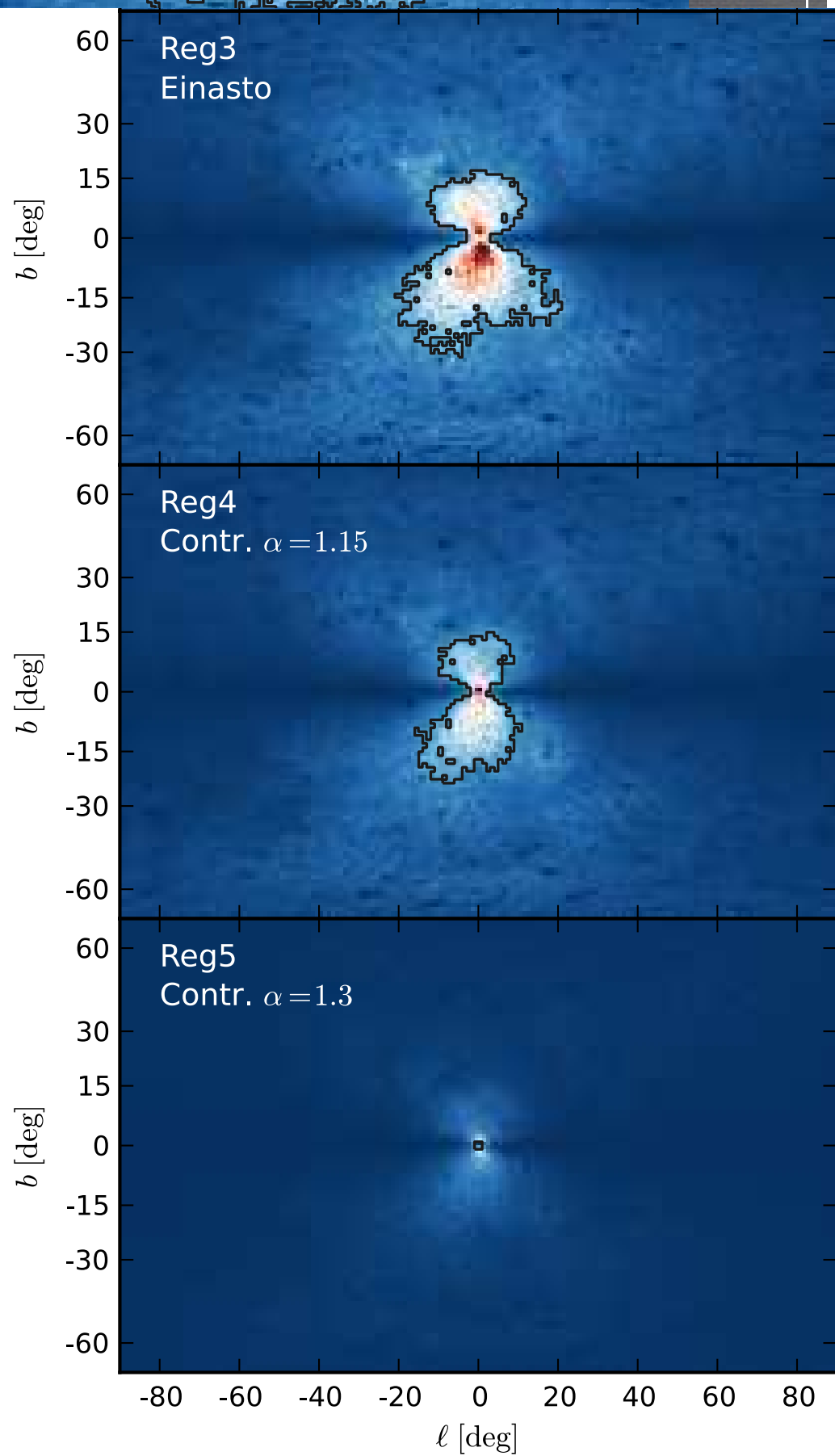


GALACTIC COORDINATES

l : LONGITUDE

b : LATITUDE

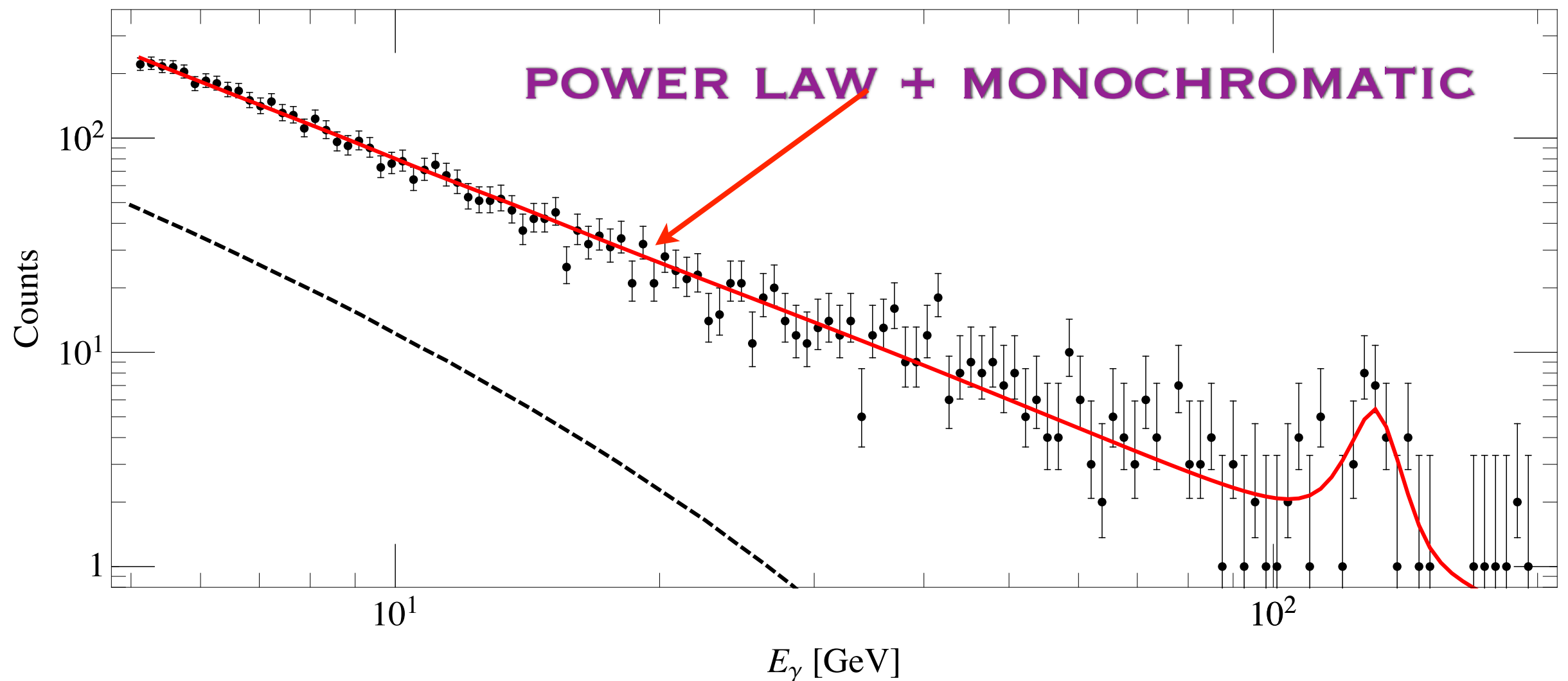
THE CENTER: $l=b=0$



$$v\sigma_{\gamma\gamma} \simeq 10^{-27} \text{ cm}^3 \text{ s}^{-1}$$

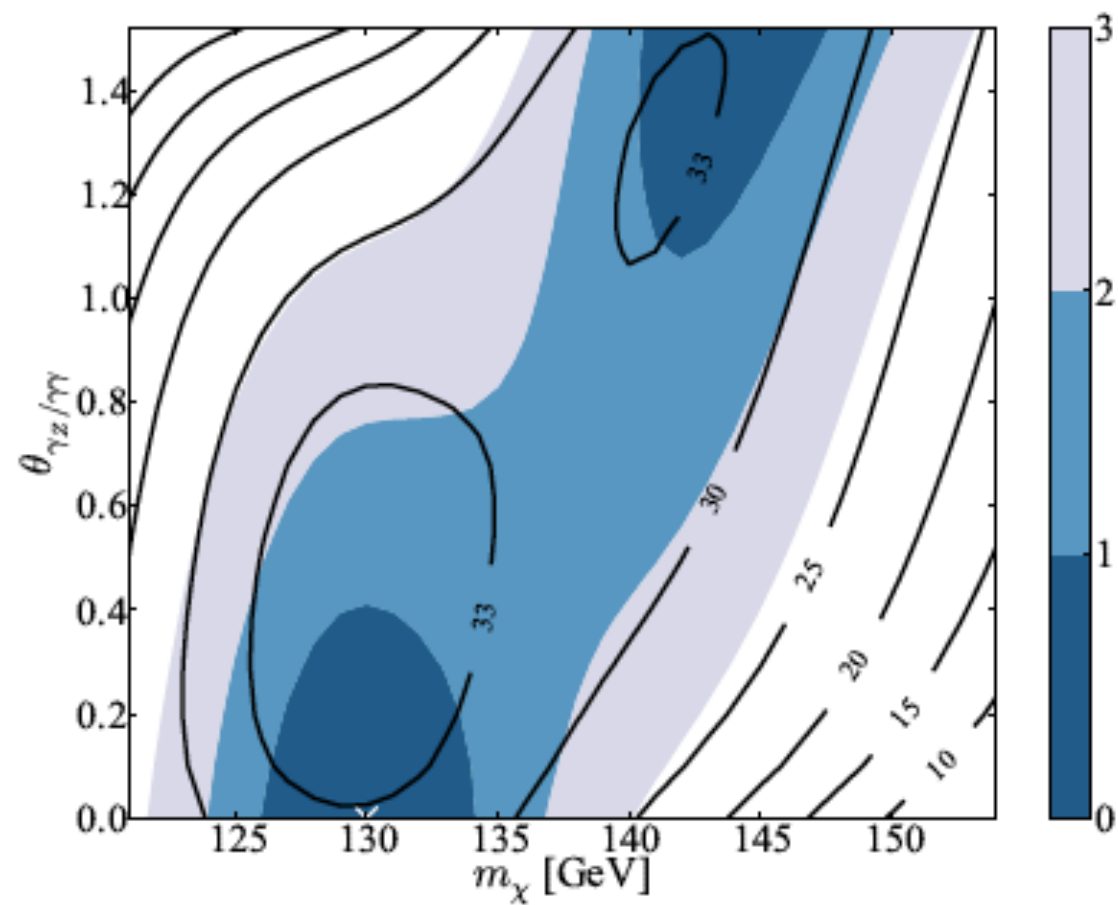
LOOKING AT THE CENTER MORE IN DETAIL.
(INNER 3 DEGREE RADIUS AROUND THE CENTER)

COHEN ET AL, ARXIV:1207.0800

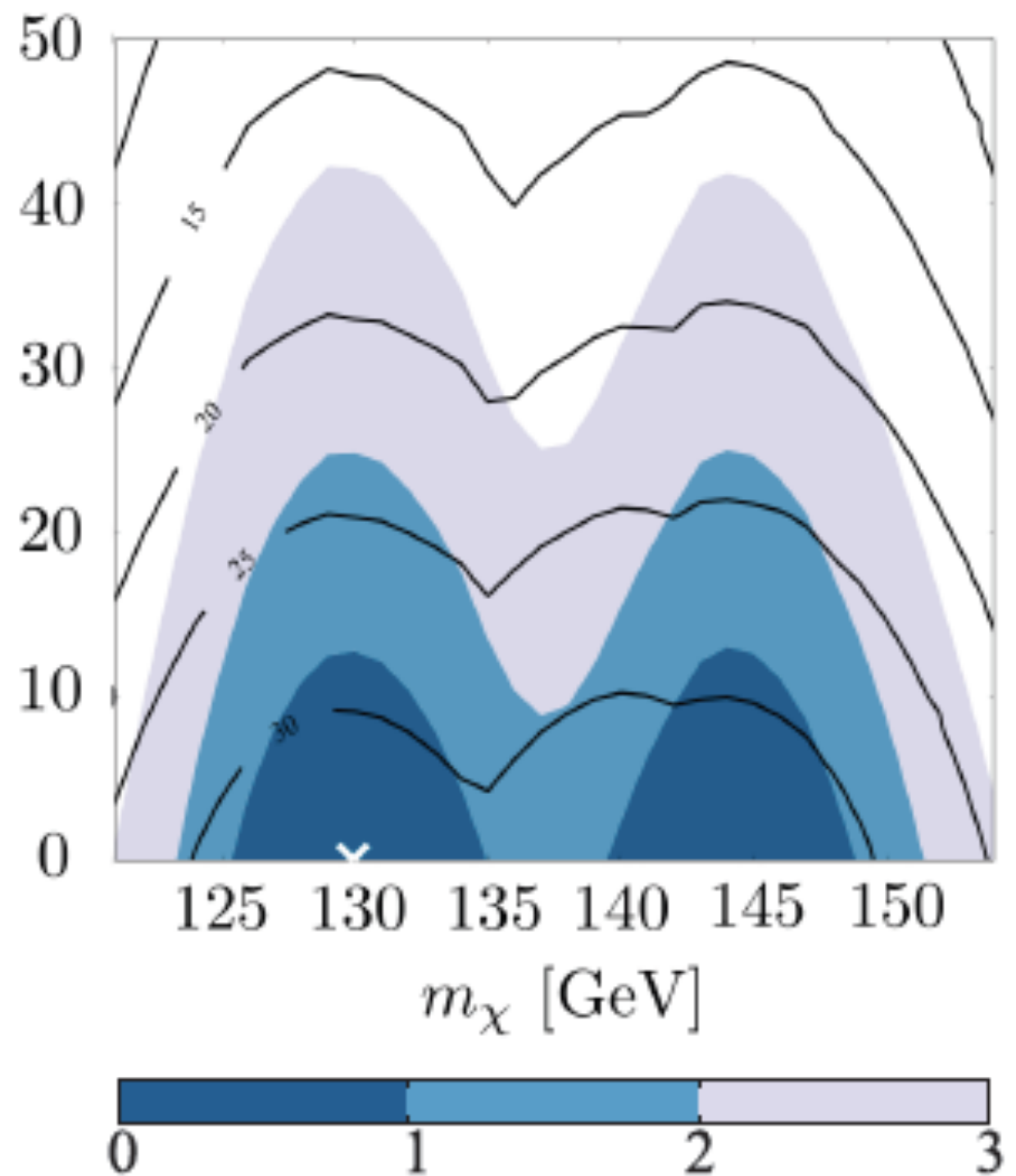


THE BEST FIT: NO ANNIHILATIONS INTO W^+W^- , ZZ , $Z\gamma$ ETC.

$$\sigma_{\gamma Z} / \sigma_{\gamma\gamma}$$



$$\sigma_{W^+W^-} / (2\sigma_{\gamma\gamma} + \sigma_{\gamma Z})$$



COHEN ET AL,
ARXIV:1207.0800

DILEMMA

* MONOCHROMATIC GAMMA

$$v\sigma_{\gamma\gamma} \simeq 10^{-27} \text{ cm}^3\text{s}^{-1} \text{ (LOOP EFFECT)}$$

* CONTINUUM GAMMA

$$v\sigma(\text{DM DM} \rightarrow \text{SM}), i.e. v\sigma(\text{DM DM} \rightarrow W^+W^-) \text{ etc}$$

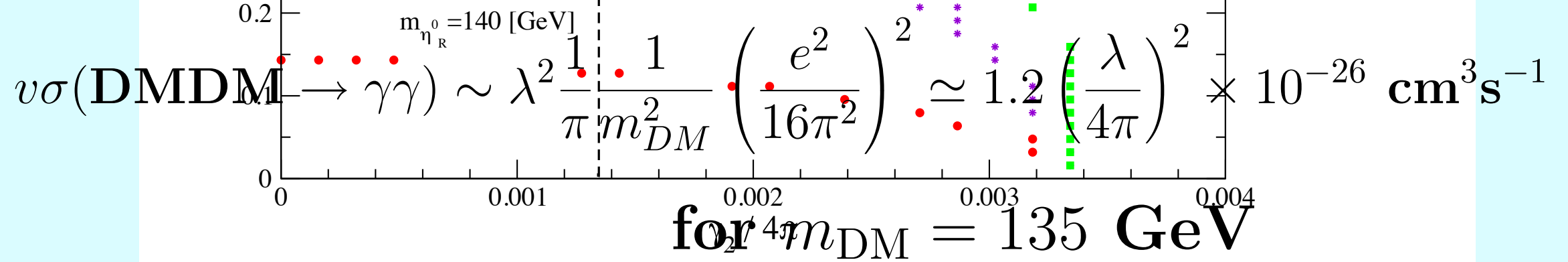
$$\lesssim 10 \times v\sigma(\text{DM DM} \rightarrow \gamma\gamma) \sim 10^{-26} \text{ cm}^3\text{s}^{-1} \text{ (TREE LEVEL)}$$

* OBSERVED RELIC DENSITY

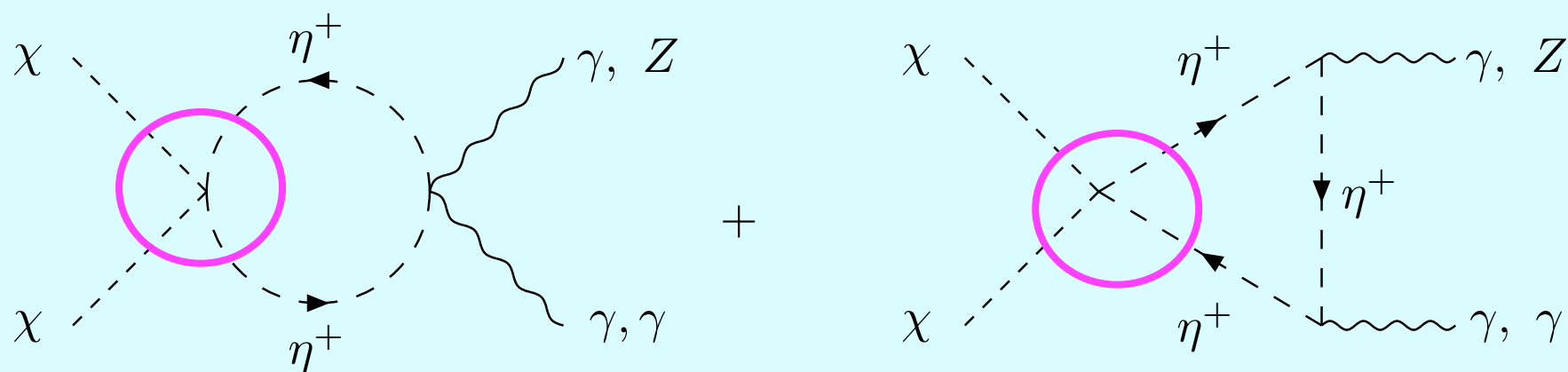
$$\Omega_T h^2 = 0.116$$

$$v\sigma(\text{DM DM}) \simeq 10^{-26} \text{ cm}^3\text{s}^{-1}$$

ONE-LOOP ->

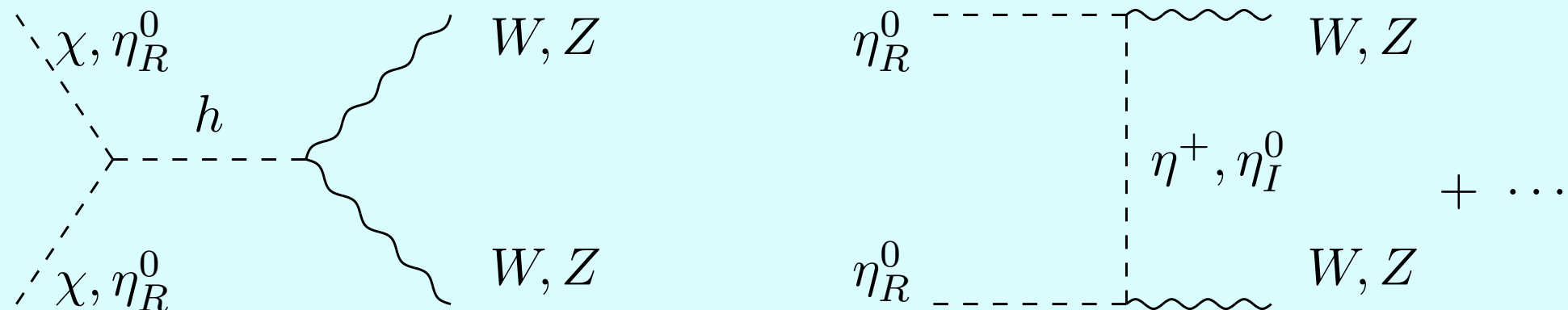


IN OUR MODEL:



HOW TO REALIZE A LARGE σ
 AT THE FREEZE OUT, WHILE SUPPRESSING σ
 IN THE GALAXY, i.e. AT LOW TEMPERATURE.

IN OUR MODEL:



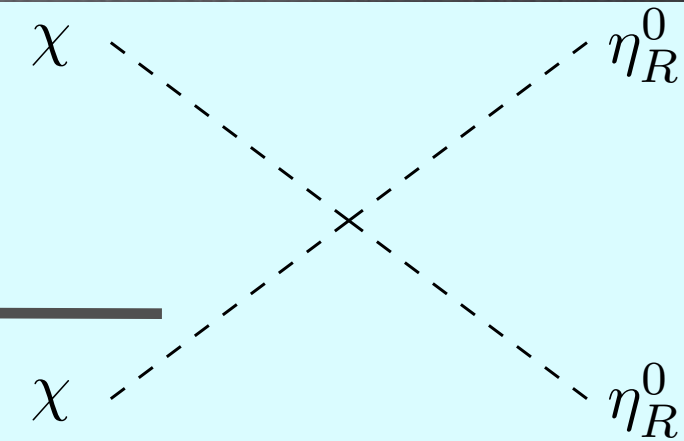
$$(2\gamma_2, \lambda_3)v_h$$

TO MAKE THE CONTINUUM GAMMAS SMALL:

FOR CHI DM: SMALL γ_2

FOR ETA DM: SMALL Ω_η

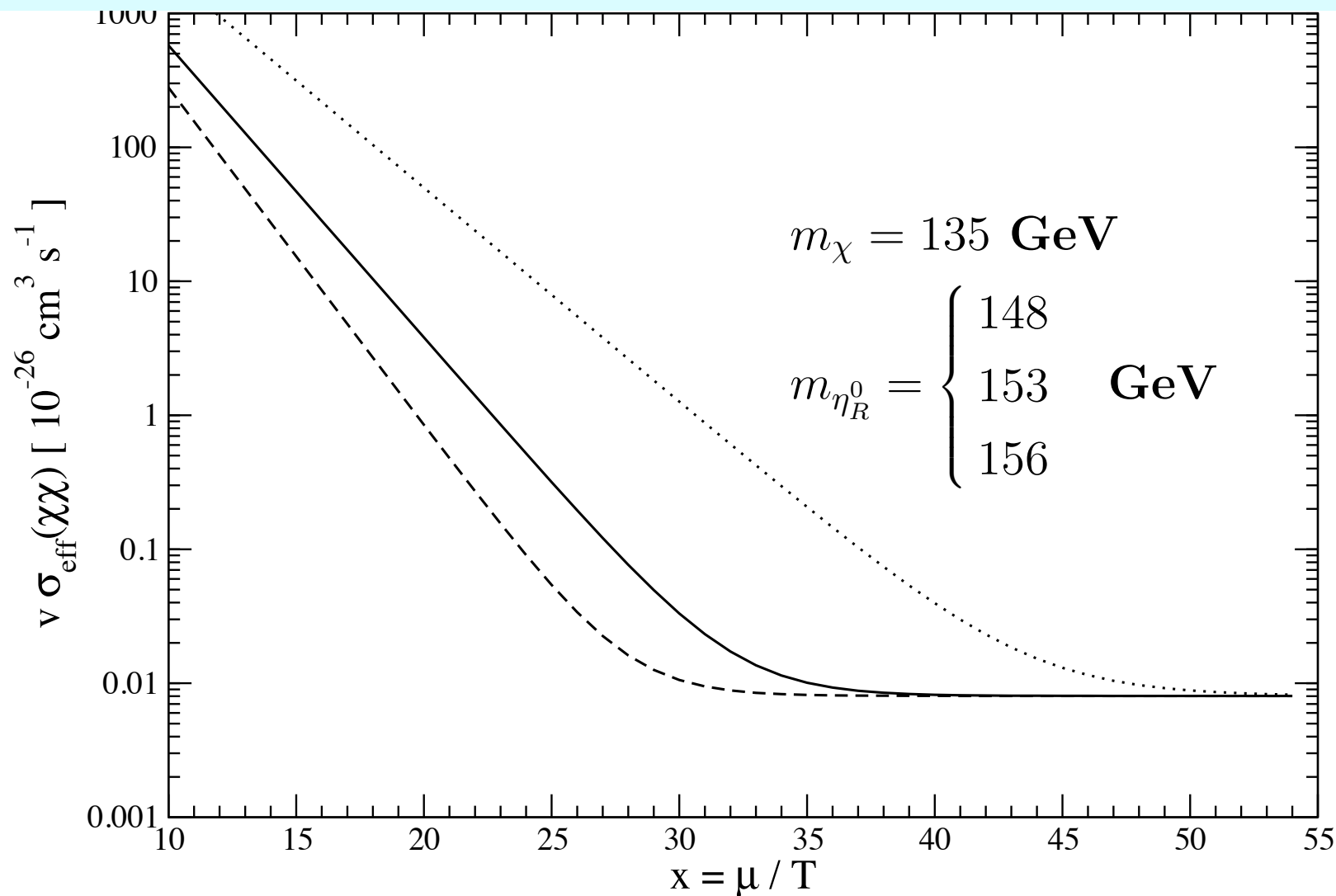
TEMPERATURE DEPENDENT CROSS SECTION:

$$\frac{dY_\chi}{dx} = -0.264 g_*^{1/2} \left[\frac{\mu M_{\text{PL}}}{x^2} \right] \left\{ \langle \sigma(\chi\chi; \text{SM})v \rangle (Y_\chi Y_\chi - \bar{Y}_\chi \bar{Y}_\chi) \right. \\ \left. - \langle \sigma(\eta_R^0 \eta_R^0; \chi\chi)v \rangle \left(Y_{\eta_R^0} Y_{\eta_R^0} - \frac{Y_\chi Y_\chi}{\bar{Y}_\chi \bar{Y}_\chi} \bar{Y}_{\eta_R^0} \bar{Y}_{\eta_R^0} \right) \right\},$$


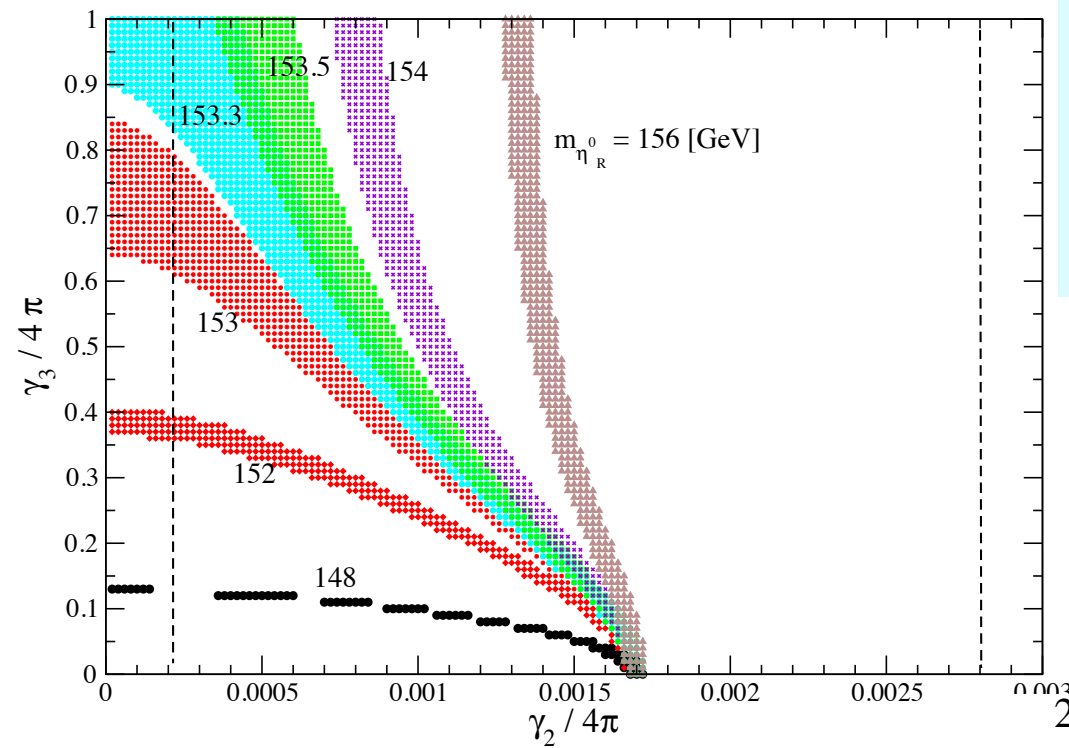
σ_{eff}

DM CONVERSION

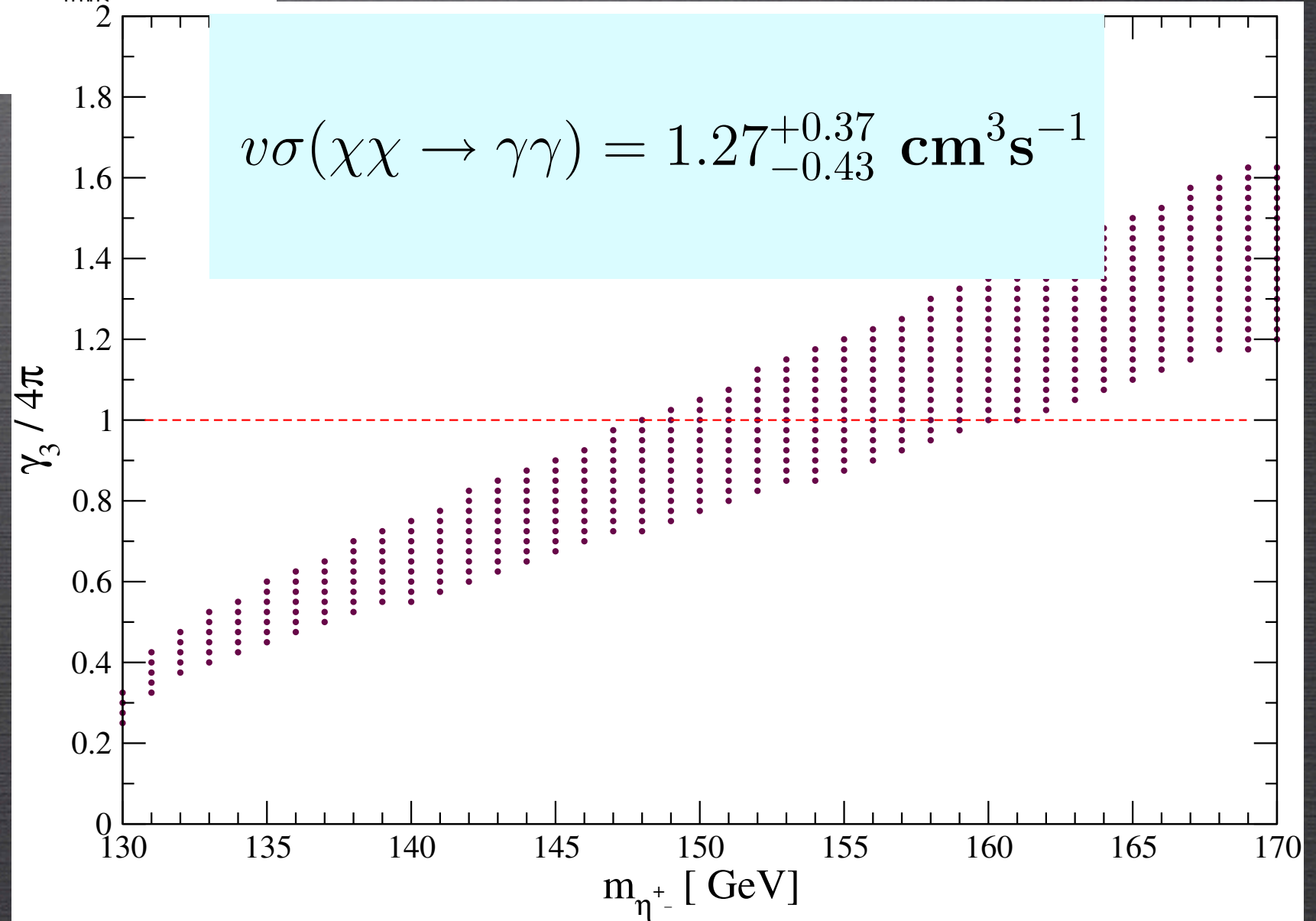
$$\left[\langle \sigma(\chi\chi; \text{SM})v \rangle + \langle \sigma(\eta_R^0 \eta_R^0; \chi\chi)v \rangle \frac{m_{\eta_R^0}^3}{m_\chi^3} \exp 2\left(\frac{m_\chi^2 - m_{\eta_R^0}^2}{m_\chi m_{\eta_R^0}}\right)x \right] (Y_\chi Y_\chi - \bar{Y}_\chi \bar{Y}_\chi)$$



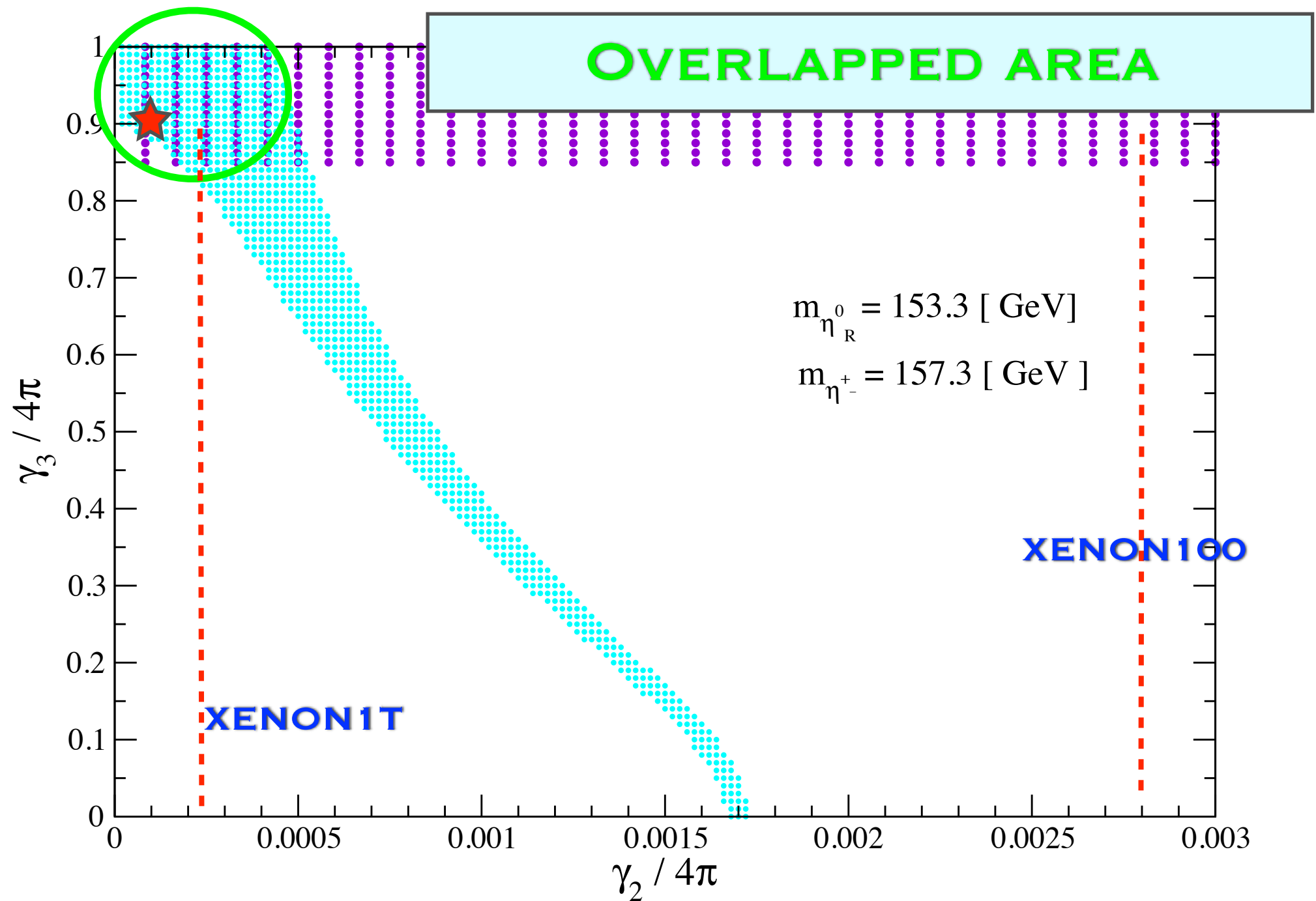
SEE ALSO:
BAEK, KO+SENAHA,
ARXIV:1209.1685



$$\Omega_T = 0.1157 \pm 0.0046(2\sigma)$$



$$v\sigma(\chi\chi \rightarrow \gamma\gamma) = 1.27^{+0.37}_{-0.43} \text{ cm}^3\text{s}^{-1}$$



$$v\sigma(\chi\chi \rightarrow \text{SM}) \simeq 8.0 \times 10^{-29} \text{ cm}^3\text{s}^{-1}$$

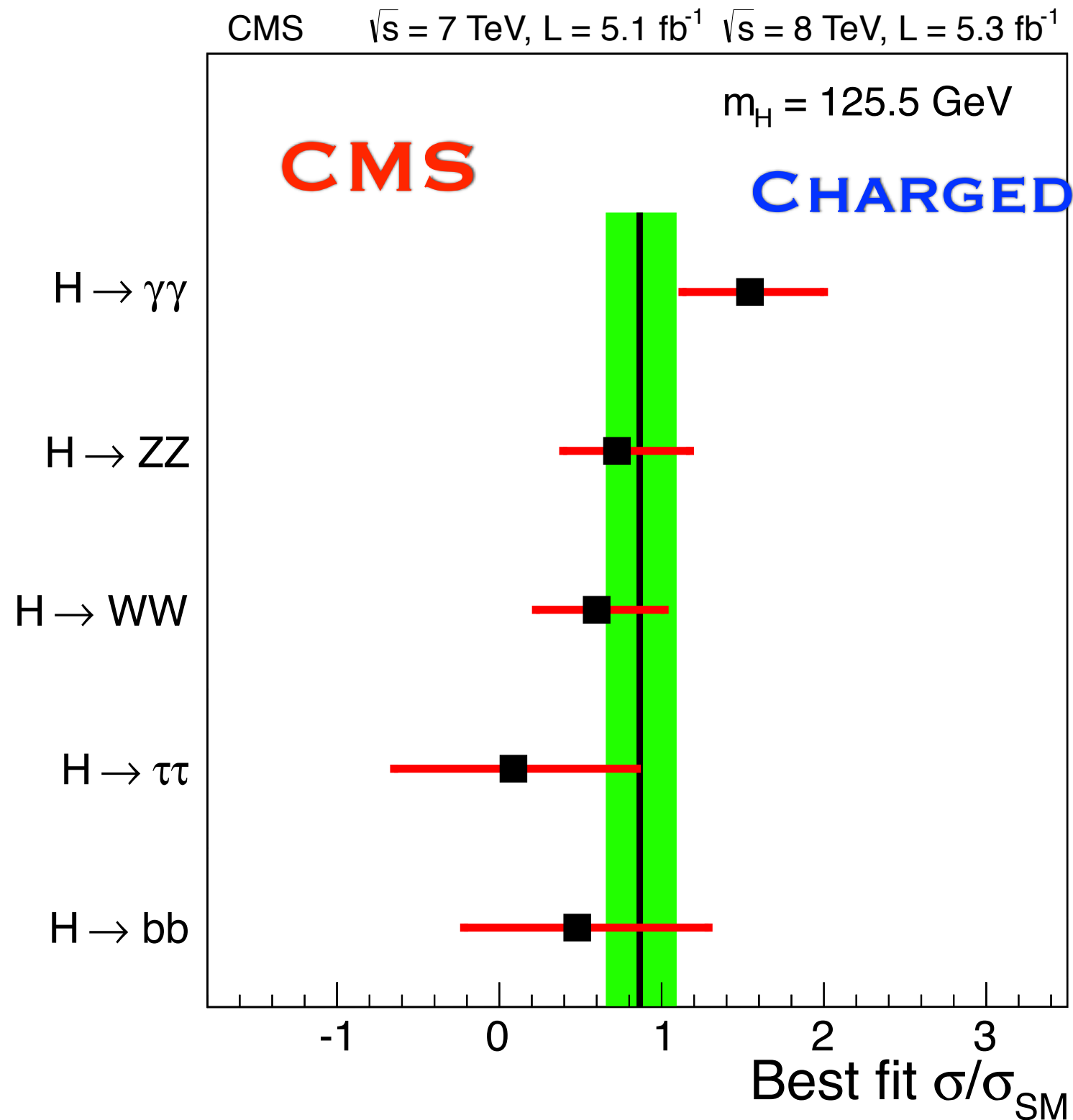
$$v\sigma(\chi\chi \rightarrow W^+W^-) \simeq 3.9 \times 10^{-29} \text{ cm}^3\text{s}^{-1}$$

$$\star v\sigma(\chi\chi \rightarrow ZZ) \simeq 1.7 \times 10^{-29} \text{ cm}^3\text{s}^{-1}$$

$$v\sigma(\chi\chi \rightarrow hh) \simeq 2.5 \times 10^{-29} \text{ cm}^3\text{s}^{-1}$$

$$v\sigma(\chi\chi \rightarrow f\bar{f}) \simeq 1.1 \times 10^{-31} \text{ cm}^3\text{s}^{-1}$$

* YET ANOTHER BONUS:



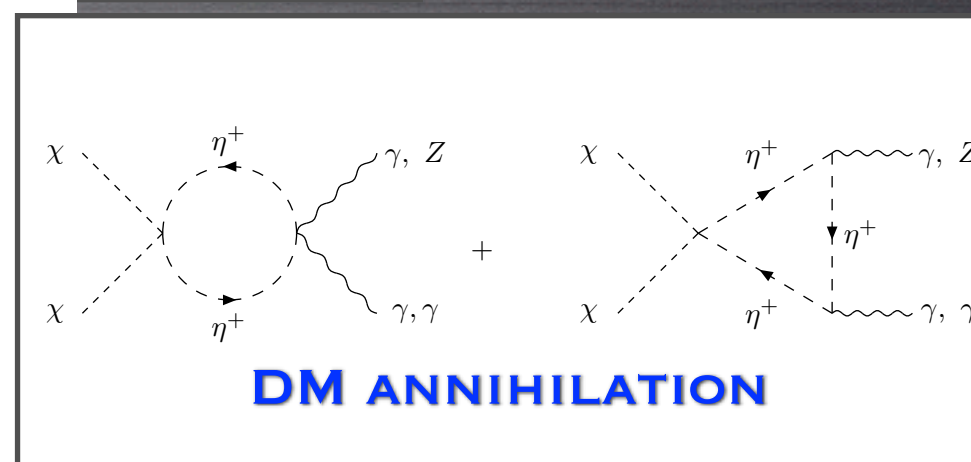
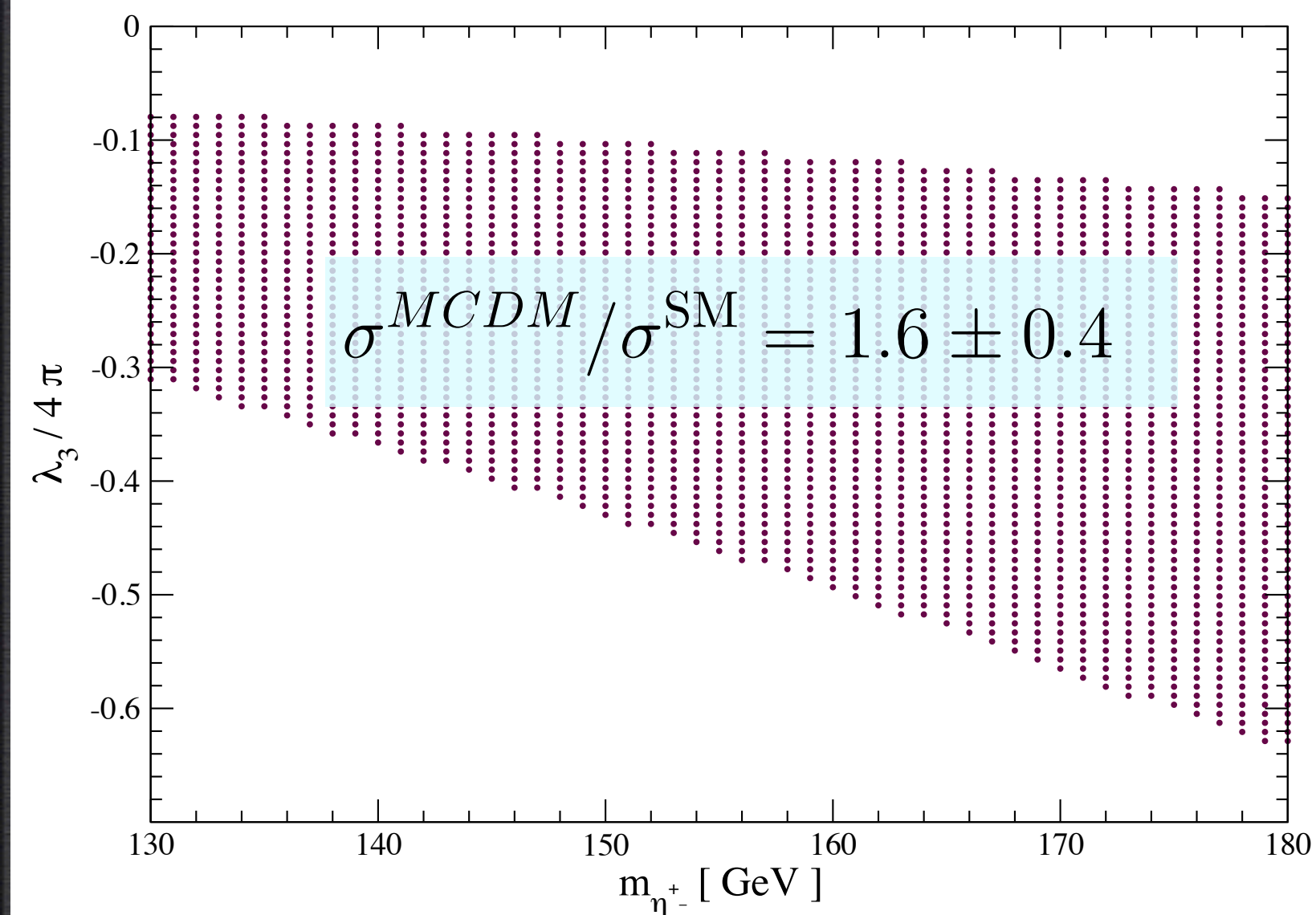
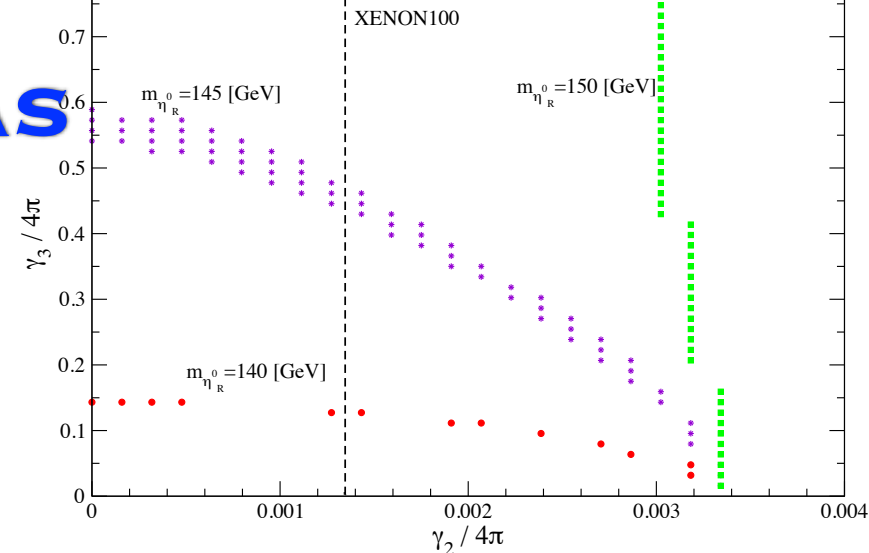
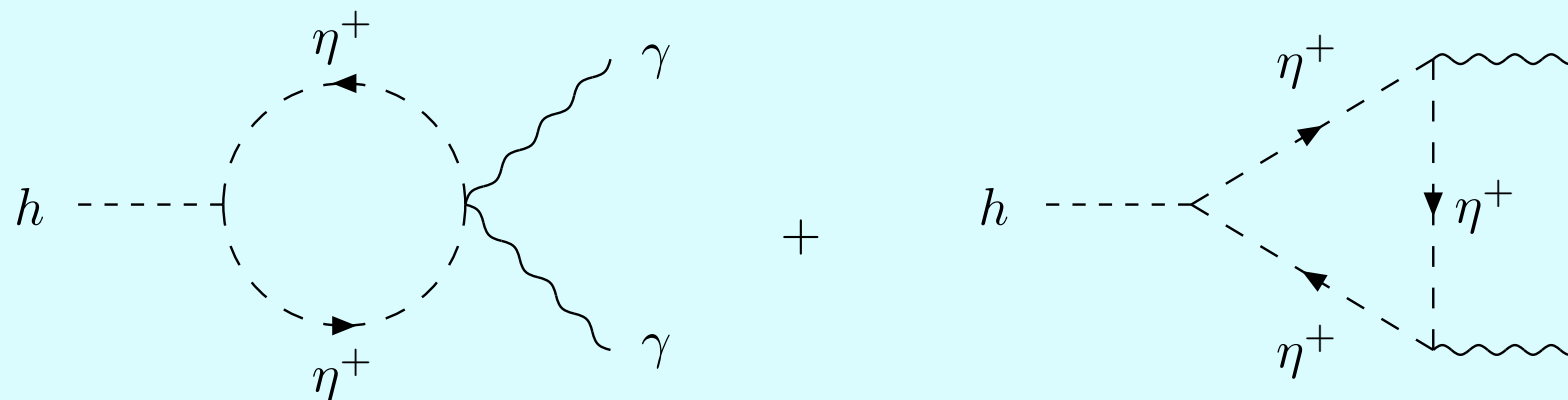
$$\sigma^{MCDM} / \sigma^{\text{SM}} = 1.6 \pm 0.4$$

CMS

CMS COLLABORATION,
ARXIV:1207.7235

ATLAS COLLABORATION
ARXIV:1207.7214

HIGGS DECAY INTO TWO GAMMAS



GUNION, HABER, KANE + DAWSON;
AKERDY, DIAZ + RIVER;
ARHRIB, BENBRIK + GAUR;
SWIEZEWSKA + KRAWCZYK; ETC

IV CONCLUSION

- I A TWO-LOOP SEESAW MODEL IS PROPOSED;
 - * L IS SOFTLY VIOLATED BY A DIM. 2 OPERATOR,
 - * $Z_2 \times Z_2$ IS THE UNBROKEN SYMMETRY, AND
 - * THE HIGGS SECTOR IS MINIMAL.
- 2 WITH η_R^0 and χ AS DM THE MODEL HAS A POTENTIAL TO EXPLAIN:
- * 135 GEV GAMMA-RAY LINE OBSERVED AT THE FERMI LAT AND
 - * ENHANCEMENT OF $h \rightarrow \gamma\gamma$ OBSERVED AT LHC.

TANK YOU VERY MUCH!

$$B(\mu \rightarrow e\gamma) = \frac{3\alpha}{64\pi(G_F m_{\eta^\pm}^2)^2} \left| \sum_k Y_{\mu k}^\nu Y_{ek}^\nu F_2 \left(\frac{M_k^2}{m_{\eta^\pm}^2} \right) \right|^2 \lesssim 2.4 \times 10^{-12}$$

$$F_2(x) = \frac{1}{6(1-x)^4} (1 - 6x + 3x^2 + 2x^3 - 6x^2 \ln x) ,$$

$$\delta a_\mu = \frac{m_\mu^2}{16\pi^2 m_{\eta^\pm}^2} \sum_k Y_{\mu k}^\nu Y_{\mu k}^\nu F_2 \left(\frac{M_k^2}{m_{\eta^\pm}^2} \right)$$

$$|\delta a_\mu| \simeq 2.2 \times 10^{-5} B(\mu \rightarrow e\gamma) \lesssim 3.4 \times 10^{-11}$$

$$\Delta T \simeq 0.54 \left(\frac{m_{\eta^\pm} - m_{\eta_R^0}}{v} \right) \left(\frac{m_{\eta^\pm} - m_{\eta_I^0}}{v} \right) = 0.02_{-0.12}^{+0.11}$$

$$\begin{aligned} (\mathcal{M}_\nu)_{ij} &= \left(\frac{1}{16\pi^2} \right)^2 \frac{\kappa^2 v_h^2}{8} \sum_k Y_{ik}^\nu Y_{jk}^\nu \int_0^\infty dx \{ B_0(-x, m_\chi, m_{\phi_R}) - B_0(-x, m_\chi, m_{\phi_I}) \} \\ &\quad \times \frac{x}{(x + m_\eta^2)^2 (x + M_k^2)} \text{ for } m_\eta = m_{\eta_R^0} \simeq m_{\eta_I^0} \\ &\sim -\lambda_5^{\text{eff}} v_h^2 \sum_k \frac{Y_{ik}^\nu Y_{jk}^\nu}{16\pi^2 M_k} \left(\ln \left(\frac{m_{\eta_R^0}}{M_k} \right)^2 + 1 \right) \text{ for } m_\eta \ll M_k , \end{aligned}$$

Coupled Boltzmann eqs.

D'ERAMO+THALER, 11, BELANGER, KANNIKE, PUKOV+RAIDAL, 12;
AOKI, DUERR, KUBO+TAKANO, 12.



standard

$$\frac{dY_i}{dx} = -0.264 g_*^{1/2} \left[\frac{\mu M_{\text{PL}}}{x^2} \right] \left\{ \langle \sigma(ii; X_i X'_i) v \rangle (Y_i Y_i - \bar{Y}_i \bar{Y}_i) \right.$$

 **conversion**

$$+ \sum_{i>j} \langle \sigma(ii; jj) v \rangle \left(Y_i Y_i - \frac{Y_j Y_j}{\bar{Y}_j \bar{Y}_j} \bar{Y}_i \bar{Y}_i \right) - \sum_{j>i} \langle \sigma(jj; ii) v \rangle \left(Y_j Y_j - \frac{Y_i Y_i}{\bar{Y}_i \bar{Y}_i} \bar{Y}_j \bar{Y}_j \right)$$

 **semi-annihilation**

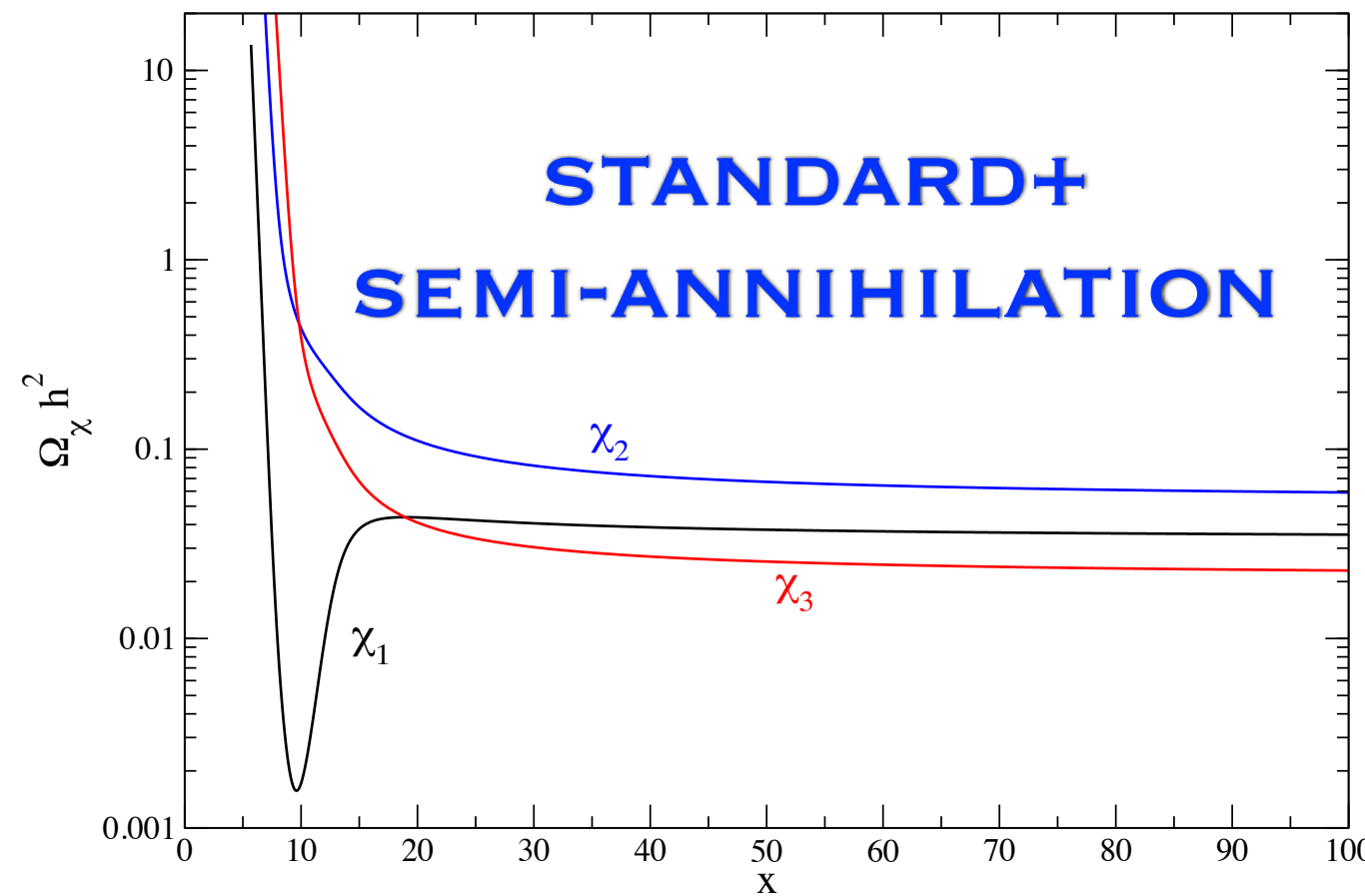
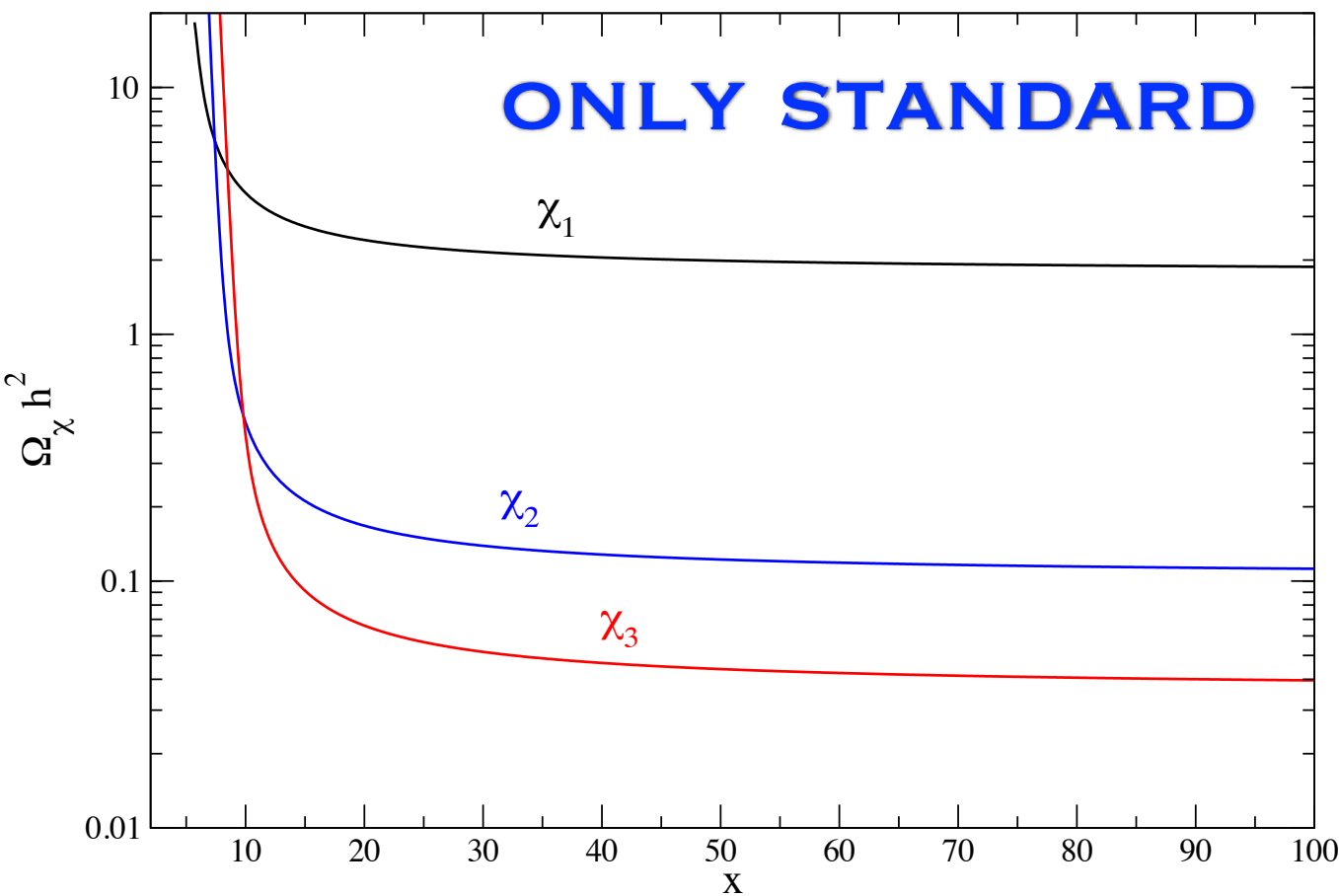
$$+ \sum_{j,k} \langle \sigma(ij; k X_{ijk}) v \rangle \left(Y_i Y_j - \frac{Y_k}{\bar{Y}_k} \bar{Y}_i \bar{Y}_j \right) - \sum_{j,k} \langle \sigma(jk; i X_{jki}) v \rangle \left(Y_j Y_k - \frac{Y_i}{\bar{Y}_i} \bar{Y}_j \bar{Y}_k \right) \}$$

$$Y_i = n_i / s \quad \mu = (\sum_i m_i^{-1})^{-1}$$

$$m_1 = 200 \text{ GeV} , m_2 = 160 \text{ GeV} , m_3 = 140 \text{ GeV}$$

STANDARD: $\sigma_{0,1} = 0.1$, $\sigma_{0,2} = 2$, $\sigma_{0,3} = 6$

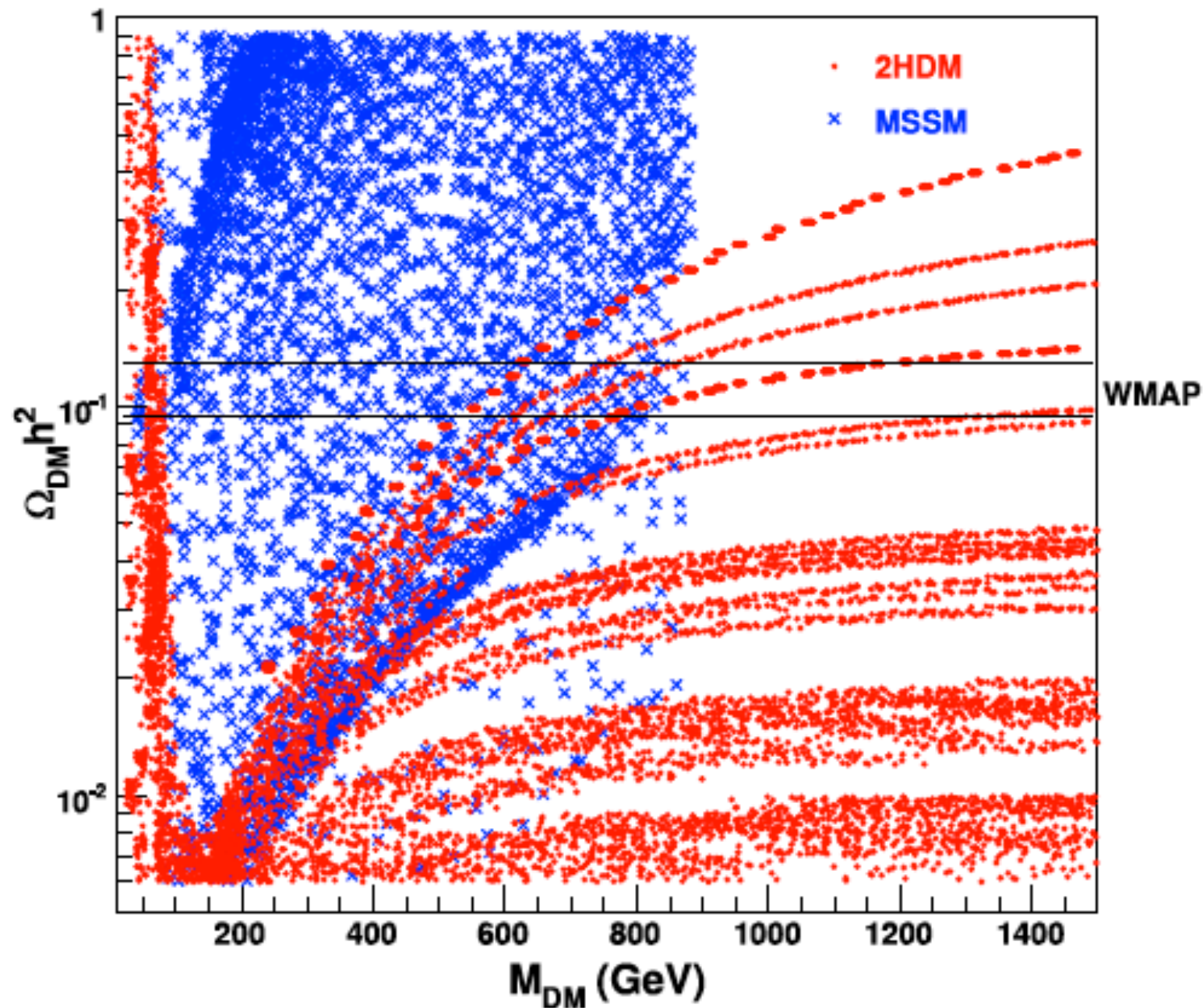
Temperature evolution



AOKI, DUERR, KUBO+TAKANO, 2012

$$\sigma_{0,123} = \sigma_{0,312} = \sigma_{0,231} = 5.1$$

Relic abundance



LOPEZ, OLIVER AND
TYTGAT,'06

Only low mass and higher mass regimes are allowed.

SUSY

Higgsino:

$$\sigma_{\gamma\gamma}v \simeq 1.1 \times 10^{-28} \text{ cm}^3/\text{s}$$

$$\sigma_{\gamma Z}v \simeq 3.7 \times 10^{-28} \text{ cm}^3/\text{s}$$

$$\sigma_{\text{ann}}v \simeq \sigma_{WW}v + \sigma_{ZZ}v \simeq 4.2 \times 10^{-25} \text{ cm}^3/\text{s}$$

Wino:

$$\sigma_{\gamma\gamma}v \simeq 2.5 \times 10^{-27} \text{ cm}^3/\text{s}$$

$$\sigma_{\gamma Z}v \simeq 1.4 \times 10^{-26} \text{ cm}^3/\text{s}$$

$$\sigma_{\text{ann}}v \simeq \sigma_{WW}v \simeq 4.0 \times 10^{-24} \text{ cm}^3/\text{s}$$

Bino:

$$\sigma_{\gamma\gamma}v \simeq \text{few} \times 10^{-30} \text{ cm}^3/\text{s};$$

$$\sigma_{\gamma Z}v \simeq \text{few} \times 10^{-31} \text{ cm}^3/\text{s};$$

$$\sigma_{\text{ann}}v \simeq \sigma_{\ell\bar{\ell}}v \simeq \text{few} \times 10^{-27} \text{ cm}^3/\text{s}.$$

**NEUTRALINO DM CAN NOT EXPLAIN
THE 130 GAMMA RAY LINE.**

