

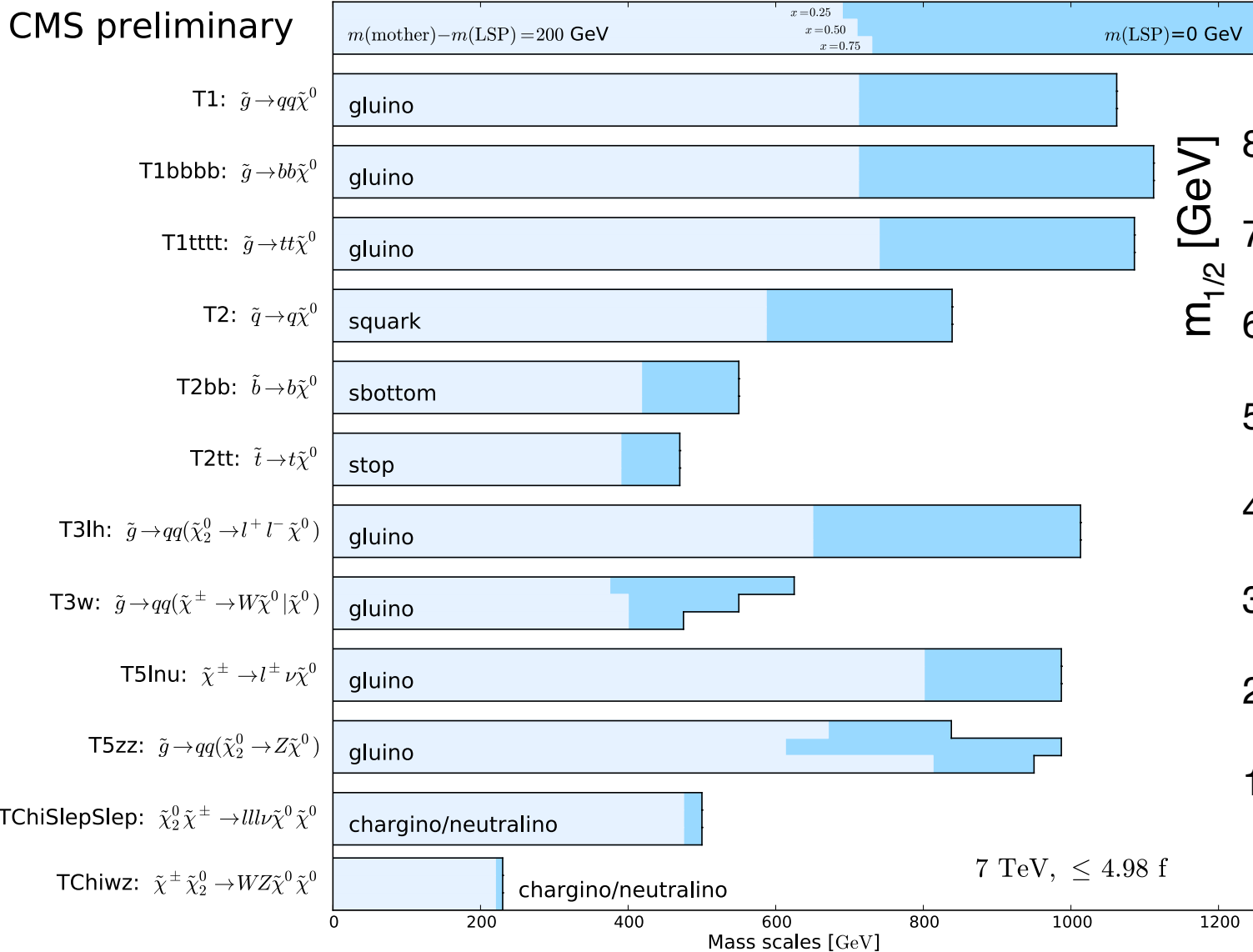
Is there something missing in the MSSM?

Jae Yong Lee (Korea University)

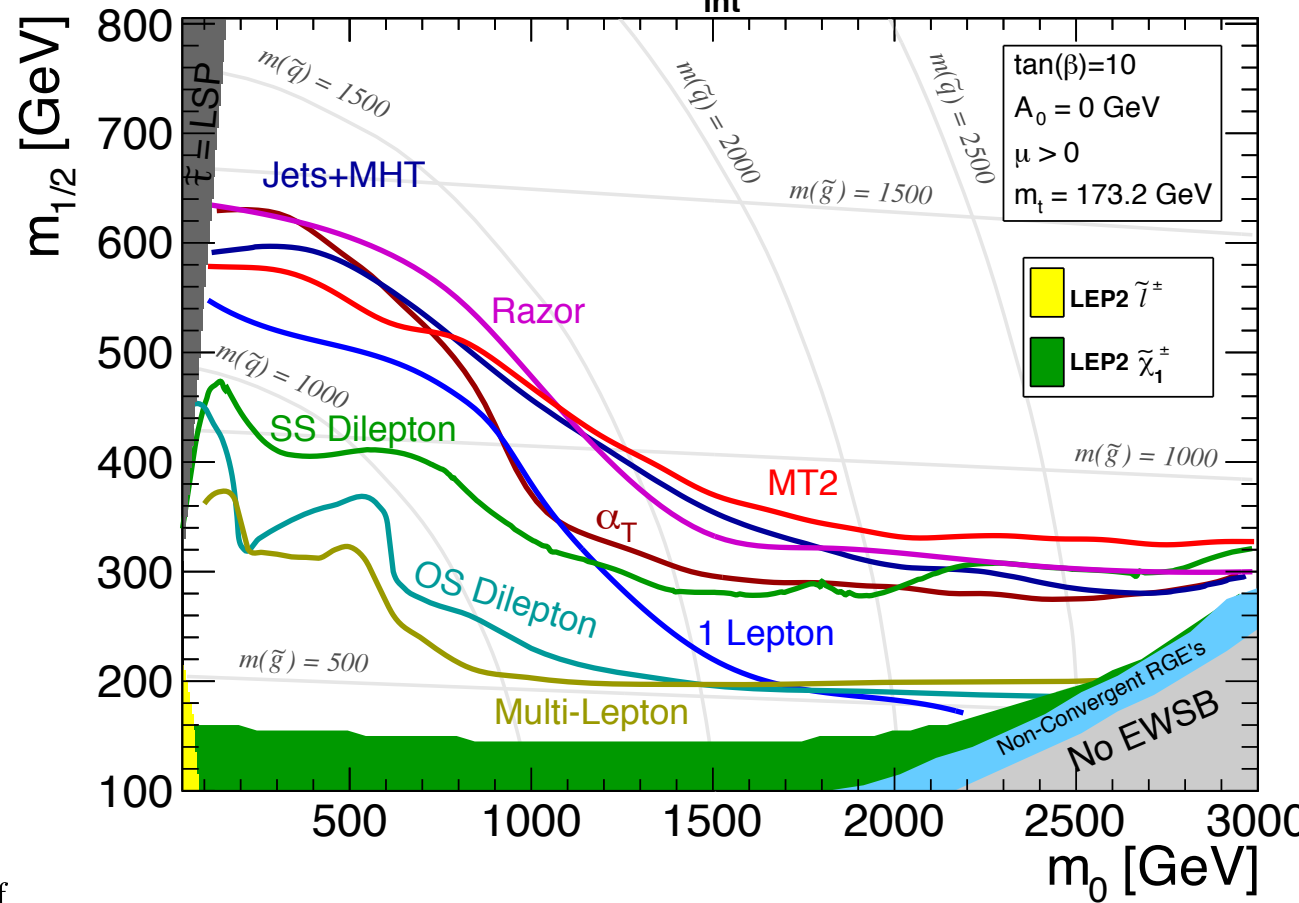
HPNP2013

Feb. 13-16, University of Toyama

CMS preliminary



CMS Preliminary $L_{int} = 4.98 \text{ fb}^{-1}, \sqrt{s} = 7 \text{ TeV}$



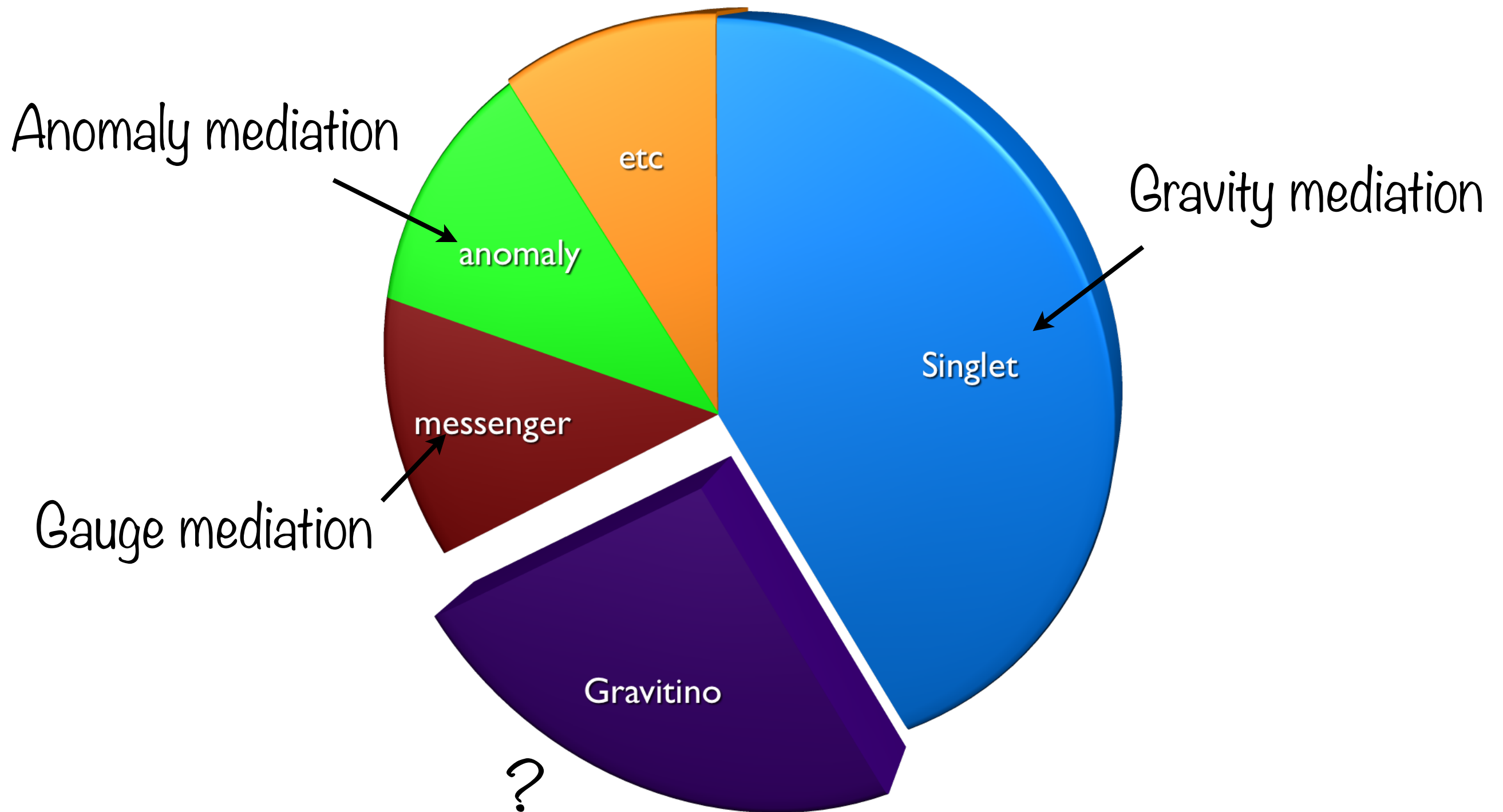
“Where is susy?”

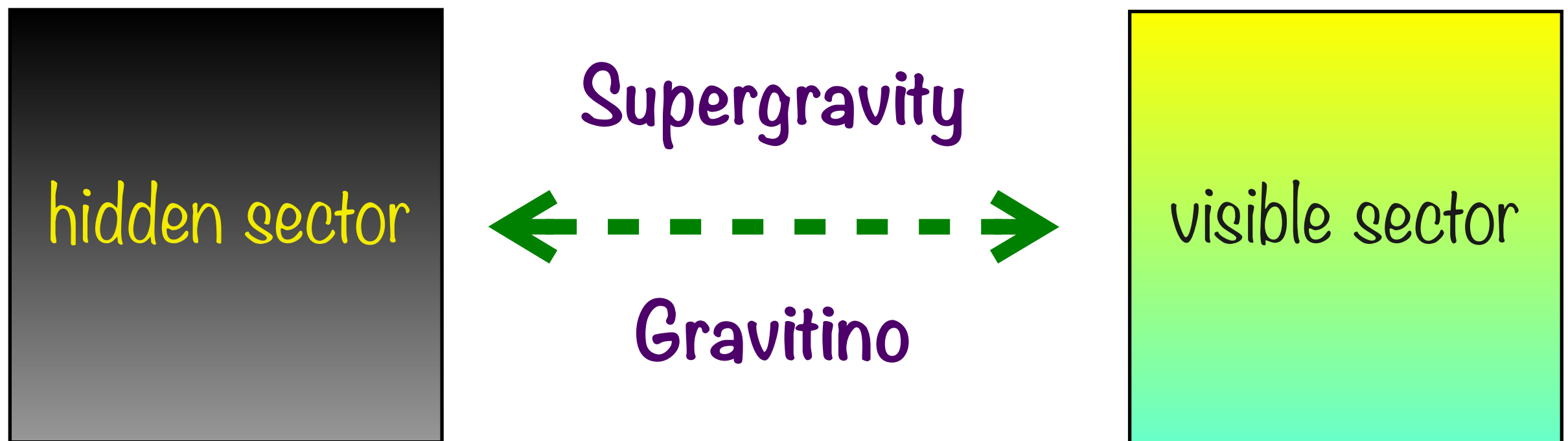




Current Situation of Low-energy Supersymmetry

Contributions to soft mass





- gauge field of local supersymmetry transformations.
- massive  “supersymmetric Higgs mechanism”.
- no gauge coupling for local supersymmetry.

gravitino kinetic term :

$$\mathcal{L}_{3/2} = \frac{1}{\kappa^2} \left[\frac{i}{2} \varepsilon^{\mu\nu\rho\sigma} {}^* \Psi_\mu \gamma_5 \gamma_\nu \partial_\rho \Psi_\sigma + im_{3/2} {}^* \Psi_\mu \Sigma^{\mu\nu} \Psi_\nu + \frac{i}{2} \alpha_s {}^* \Psi \cdot \gamma \not{\partial} \gamma \cdot \Psi \right].$$

massive gravitino propagator :

$$\frac{i}{\kappa^2} \langle \Psi_\mu {}^* \Psi_\nu \rangle = P_{\mu\nu}^{3/2} \frac{1}{m_{3/2} + \not{\partial}} + \sum_{i,j} (P_{ij}^{1/2})_{\mu\nu} [\alpha_{ij}(\square) \not{\partial} + \beta_{ij}(\square)],$$

with

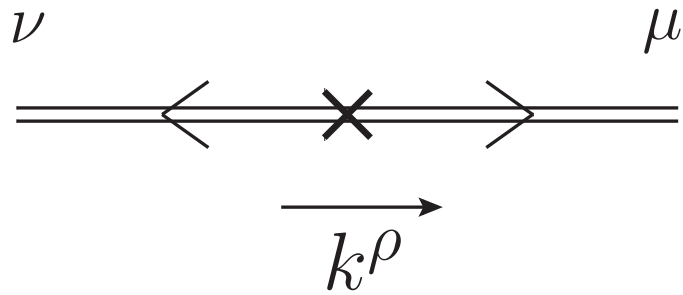
$$P_{\mu\nu}^{3/2} = \theta_{\mu\nu} - \frac{1}{3} \hat{\gamma}_\mu \hat{\gamma}_\nu,$$

$$(P_{11}^{1/2})_{\mu\nu} = \frac{1}{3} \hat{\gamma}_\mu \hat{\gamma}_\nu, \quad (P_{12}^{1/2})_{\mu\nu} = \sqrt{\frac{1}{3}} \hat{\gamma}_\mu \omega_\nu,$$

$$(P_{21}^{1/2})_{\mu\nu} = \sqrt{\frac{1}{3}} \omega_\mu \hat{\gamma}_\nu, \quad (P_{22}^{1/2})_{\mu\nu} = \frac{1}{3} \omega_\mu \omega_\nu,$$

$$\omega_\mu \equiv \frac{\partial_\mu \not{\partial}}{\square}, \quad \hat{\gamma}_\mu \equiv \gamma_\mu - \omega_\mu, \quad \theta_{\mu\nu} \equiv \eta_{\mu\nu} - \omega_\mu \omega_\nu.$$

chirality-flipping propagator (in two component notation)



$$\langle \psi_\mu \bar{\psi}_\nu \rangle_{\frac{3}{2}} = (-i)\kappa^2 \left[\eta_{\mu\nu} - \frac{1}{3} \left(\sigma_\mu \bar{\sigma}_\nu - \frac{k_\nu \sigma_\mu k \cdot \bar{\sigma}}{k^2} - \frac{k_\mu k \cdot \sigma \bar{\sigma}_\nu}{k^2} + 4 \frac{k_\mu k_\nu}{k^2} \right) \right] \frac{m_{3/2}}{k^2 + m_{3/2}^2}$$

$$\begin{aligned} \langle \psi_\mu \bar{\psi}_\nu \rangle_{\frac{1}{2}} = & (-i)\kappa^2 \left[\frac{1}{3} \left(\sigma_\mu \bar{\sigma}_\nu - \frac{k_\nu \sigma_\mu k \cdot \bar{\sigma}}{k^2} - \frac{k_\mu k \cdot \sigma \bar{\sigma}_\nu}{k^2} + \frac{k_\mu k_\nu}{k^2} \right) \beta_{11} + \frac{k_\mu k_\nu}{k^2} \beta_{22} \right. \\ & \left. + \frac{1}{\sqrt{3}} \left(\frac{k_\nu \sigma_\mu k \cdot \bar{\sigma}}{k^2} - \frac{k_\mu k_\nu}{k^2} \right) \beta_{12} + \frac{1}{\sqrt{3}} \left(\frac{k_\mu k \cdot \sigma \bar{\sigma}_\nu}{k^2} - \frac{k_\mu k_\nu}{k^2} \right) \beta_{21} \right] \end{aligned}$$

interaction between gravitino and gaugino

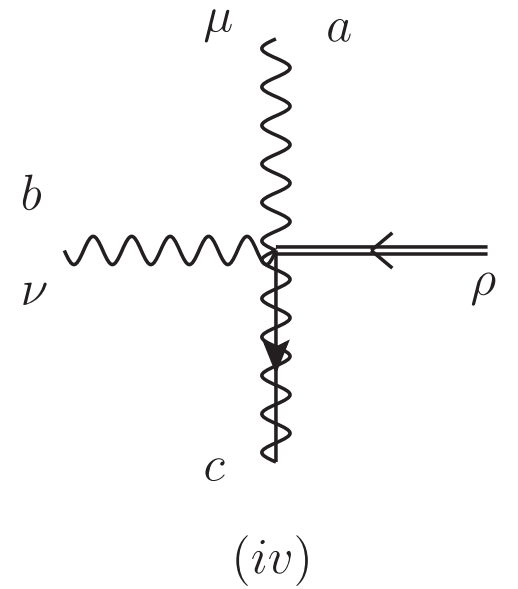
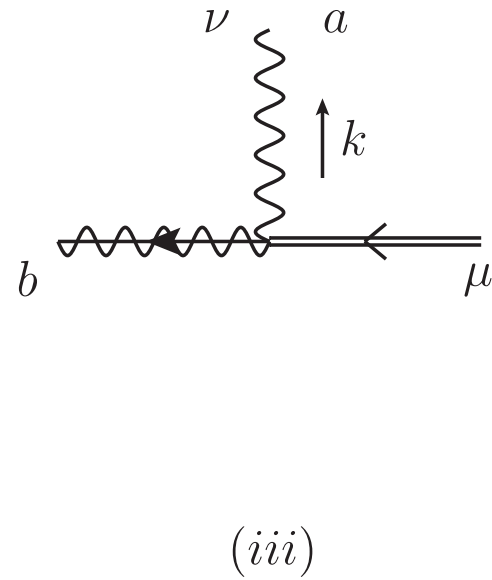
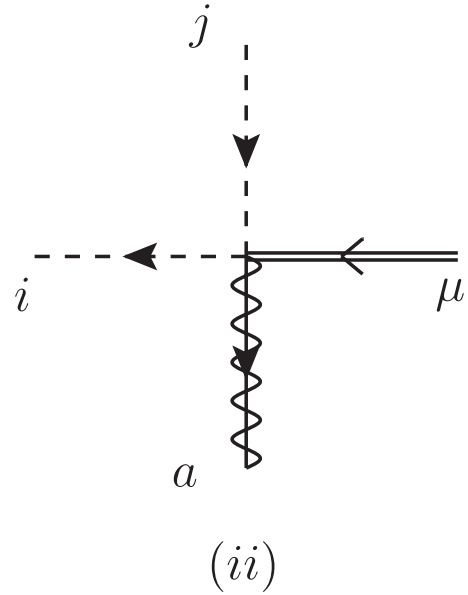
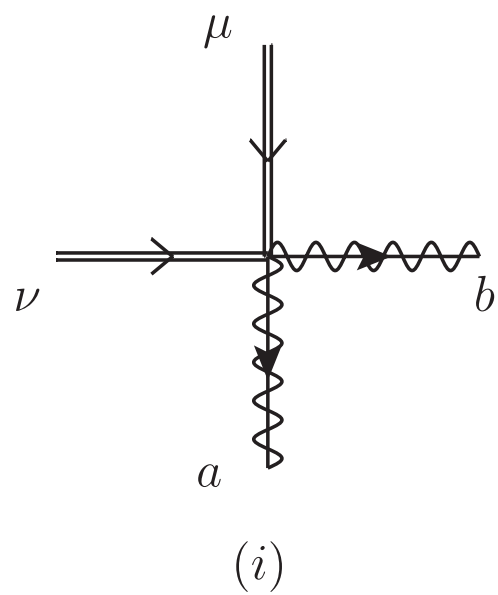
$$\mathcal{L}_{\text{sugra}} \supset \mathcal{L}_i$$

$$\mathcal{L}_1 = \frac{1}{16} [\lambda^a \lambda^a \bar{\psi}_\mu (3g^{\mu\nu} - 2\bar{\sigma}^{\mu\nu}) \bar{\psi}_\nu + \bar{\lambda}^a \bar{\lambda}^a \psi_\mu (3g^{\mu\nu} - 2\sigma^{\mu\nu}) \psi_\nu],$$

$$\mathcal{L}_2 = \frac{1}{2} [\bar{\psi}_\mu \bar{\sigma}^\mu \lambda^a + \psi_\mu \sigma^\mu \bar{\lambda}^a] \sum_{i,j} A_i^* (T^a)^i_j A_j,$$

$$\mathcal{L}_3 = \frac{i}{2} [\partial^\mu \mathcal{A}^{\nu c} - \partial^\nu \mathcal{A}^{\mu c}] [(\psi_\mu \sigma_\nu \bar{\lambda}^c - \bar{\psi}_\mu \bar{\sigma}_\nu \lambda^c) - \frac{i}{2} \varepsilon_{\mu\nu\rho\kappa} (\psi^\rho \sigma^\kappa \bar{\lambda}^c + \bar{\psi}^\rho \bar{\sigma}^\kappa \lambda^c)],$$

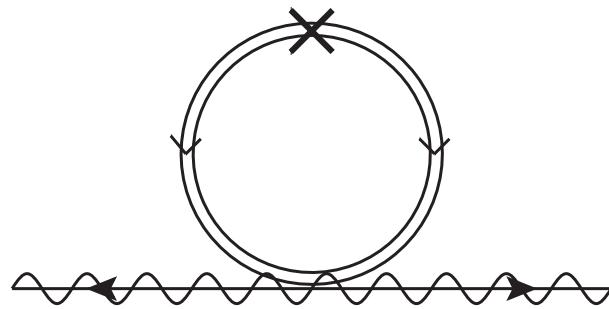
$$\mathcal{L}_4 = \frac{i}{2} g f^{abc} \mathcal{A}^{\mu a} \mathcal{A}^{\nu b} [(\psi_\mu \sigma_\nu \bar{\lambda}^c - \bar{\psi}_\mu \bar{\sigma}_\nu \lambda^c) - \frac{i}{2} \varepsilon_{\mu\nu\rho\kappa} (\psi^\rho \sigma^\kappa \bar{\lambda}^c + \bar{\psi}^\rho \bar{\sigma}^\kappa \lambda^c)],$$



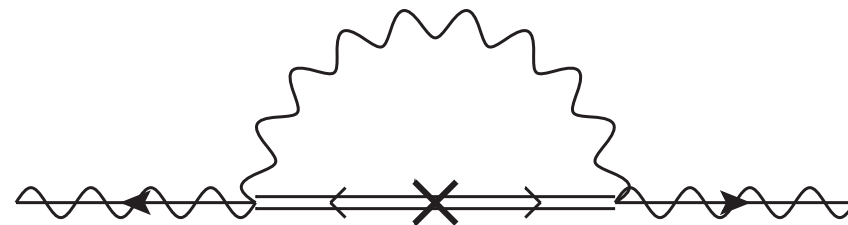
$\mathcal{L}_1$	$\frac{i}{4}[3\eta^{\mu\nu}\delta^{\dot{\alpha}\dot{\beta}} - 2(\bar{\sigma}^{\mu\nu})^{\dot{\alpha}\dot{\beta}}]\delta^{ab}$
$\mathcal{L}_2$	$\frac{i}{2}(T^a)^i_j\sigma^\mu$
$\mathcal{L}_3$	$[\frac{i}{2}(-k_\mu\sigma_\nu + \eta_{\mu\nu}k \cdot \sigma) + \frac{1}{2}k^\rho\varepsilon_{\mu\nu\rho\kappa}\sigma^\kappa]\delta^{ab}$
$\mathcal{L}_4$	$-\frac{g}{2}f^{abc}(\eta^{\mu\rho}\eta^{\nu\kappa} - \eta^{\mu\kappa}\eta^{\nu\rho})\sigma_\kappa + i\frac{g}{2}f^{abc}\varepsilon^{\mu\nu\rho\kappa}\sigma_\kappa$

Table 1: The Feynman rules for the interactions.

Gaugino mass



(a)



(b)

$$m_{1/2} = \frac{m_{3/2}}{16\pi^2} \left[\frac{\Lambda^2}{m_P^2} - \frac{m_{3/2}^2}{m_P^2} \log \left(\frac{\Lambda^2}{m_{3/2}^2} + 1 \right) \right],$$

For $m_{1/2} \ll m_P$,

$$\frac{m_{1/2}}{m_{3/2}} \sim \frac{1}{16\pi^2} \frac{\Lambda^2}{m_P^2}.$$

- No gauge or Yukawa couplings show up in the result.
- Different gaugino acquire the same mass from gravitino at one loop.
- Scale dependence stems only from the gaugino mass.

Outlook

- Evaluate sfermion mass from gravitino loops.
- Reproduce anomaly mediation from explicit one-loop effective supergravity lagrangian.