

Higgs phenomenology of the supersymmetric grand unified theory with the Hosotani mechanism

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1, Introduction

T. Appelquist's Decoupling theorem

[T. Appelquist and J. Carazzone, Phys.Rev. D11 (1975) 2856]

→ Testing GUTs is difficult at the low energy.

- Conventional GUTs are tested only by probing relations among coupling constants and mass ratio
- If new TeV scale particles are predicted, we can probe GUTs at low energy experiments

➤ Supersymmetric grand unified theory with the Hosotani mechanism

[K. Kojima, K. Takenaga and T. Yamashita, PhysRevD.84.051701]

[T. Yamashita, Phys. Rev. D84 (2011) 115016]

Doublet-Triplet splitting is naturally realized

New TeV-mass adjoint chiral supermultiplets, color octet, SU(2) triplet and singlet, are predicted

$SU(3)$	$SU(2)$	$U(1)$	
\widehat{H}_d	1	2	-1
\widehat{H}_u	1	2	+1
\widehat{S}	1	1	0
$\widehat{\Delta}$	1	3	0

MSSM doublet CP-even: h, H, S_R^0, Δ_R^0
 MSSM doublet 3+1CP-odd : A, S_I^0, Δ_I^0, G
 Singlet 3+1Charged: $H^\pm, \Delta^\pm, \bar{\Delta}^\pm, G^\pm$
 Triplet

Investigating the Higgs sector, we can test this model

2, Model

➤ Superpotential and soft-SUSY term

$$W = \mu \widehat{H}_u \cdot \widehat{H}_d + \eta \widehat{S} + \frac{\mu_s}{2} \widehat{S}^2 + \mu_\Delta \text{Tr}(\widehat{\Delta}^2) + \lambda_s \widehat{S} \widehat{H}_u \cdot \widehat{H}_d + \lambda_\Delta \widehat{H}_u \cdot \widehat{\Delta} \widehat{H}_d$$

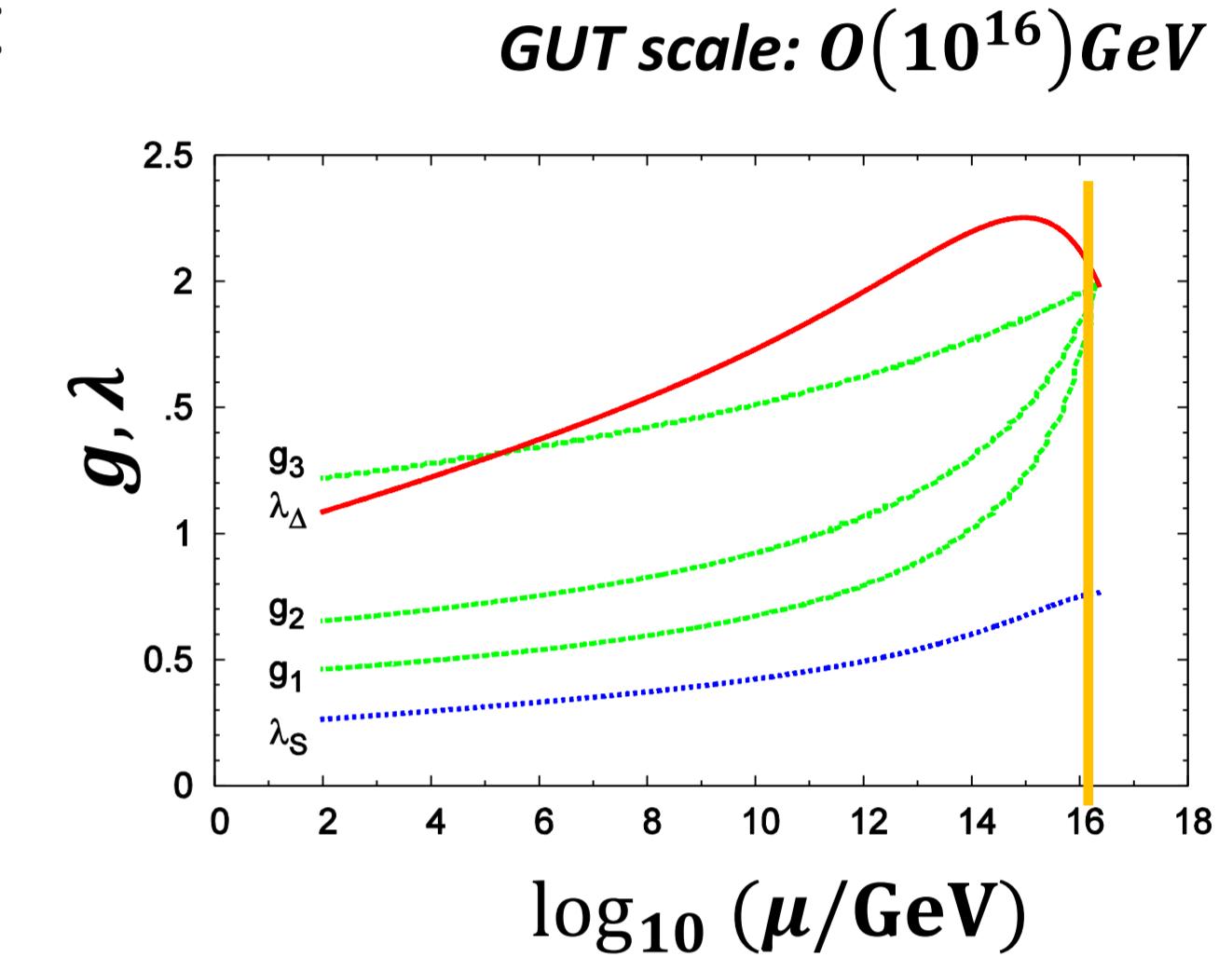
These trilinear couplings contribute to the Higgs physics

$$V_{SOFT} = \tilde{m}_d^2 |H_d|^2 + \tilde{m}_u^2 |H_u|^2 + 2\tilde{m}_\Delta^2 \text{Tr}(\Delta^\dagger \Delta) + \tilde{m}_s^2 |S|^2 + \left[B\mu H_u \cdot H_d + B_\Delta \mu_\Delta \text{Tr}(\Delta^2) + B_s \frac{\mu_s}{2} S^2 + A_\Delta \lambda_\Delta H_u \cdot \Delta H_d + A_S \lambda_s S H_u \cdot H_d + \bar{\eta} S + H.C. \right]$$

➤ Features

- Trilinear coupling $S^3, S\Delta\Delta$ are absent
- $\lambda_s, \lambda_\Delta$ are related at the GUT scale

$$\lambda_\Delta = 2\sqrt{5/3} \lambda_s = g_{GUT} \text{ @GUT scale}$$



$$\lambda_\Delta = 1.1, \lambda_s = 0.26 \text{ @EW scale}$$

- High predictability

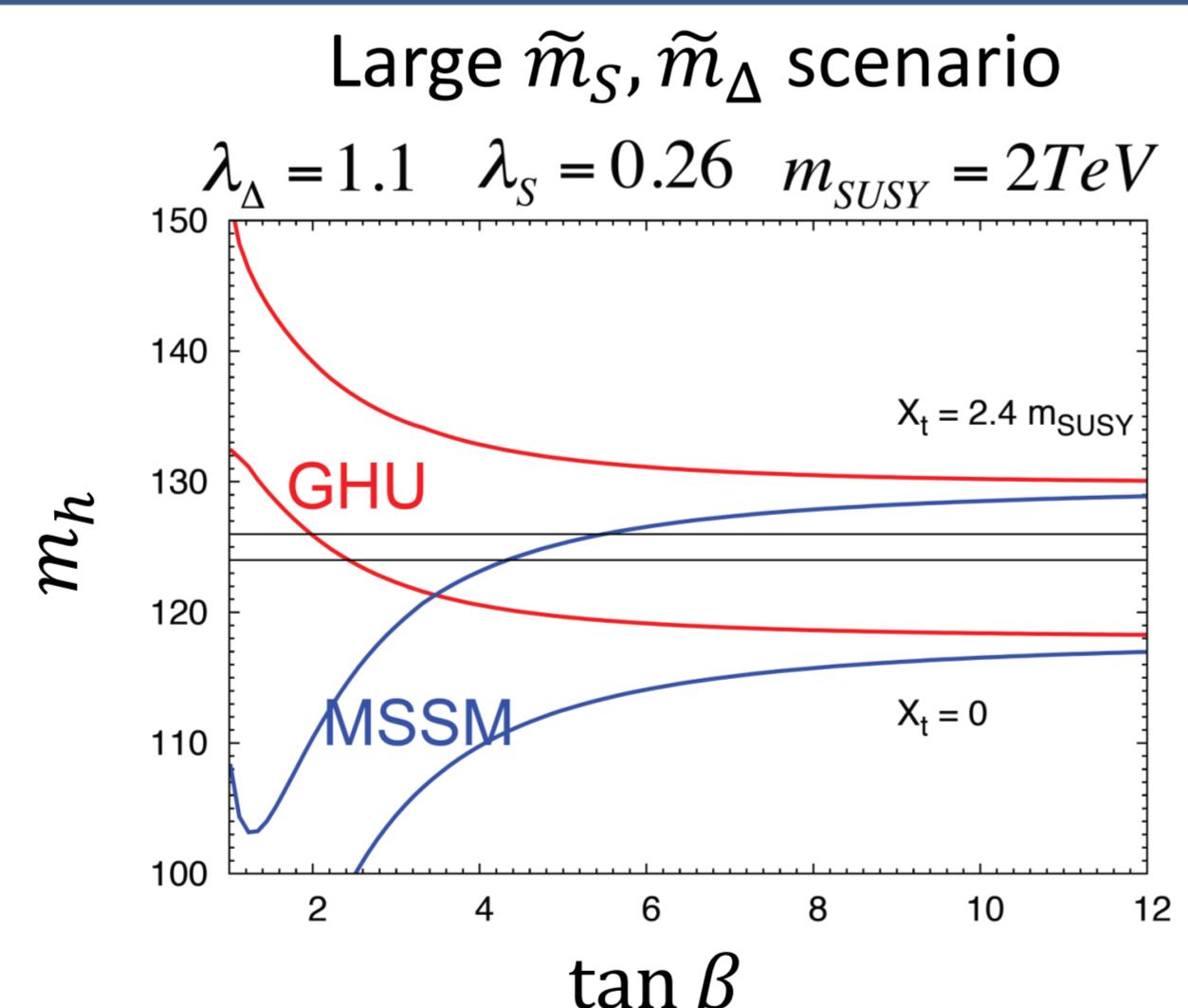
3, Phenomenology of the Higgs sector of Our model

a) Light CP-even Higgs boson mass

$$X_t = A_t - \mu \cot \beta$$

$$m_h^2 \sim m_Z^2 c_{2\beta}^2 + \frac{3m_t^4}{2\pi^2 v^2} \left(\log \frac{m_t^2}{m_b^2} + \frac{X_t^2}{m_t^2} \left(1 - \frac{X_t^2}{12m_t^2} \right) \right) + \frac{1}{2} \lambda_s^2 v^2 S_{2\beta}^2 + \frac{1}{8} \lambda_\Delta^2 v^2 S_{2\beta}^2$$

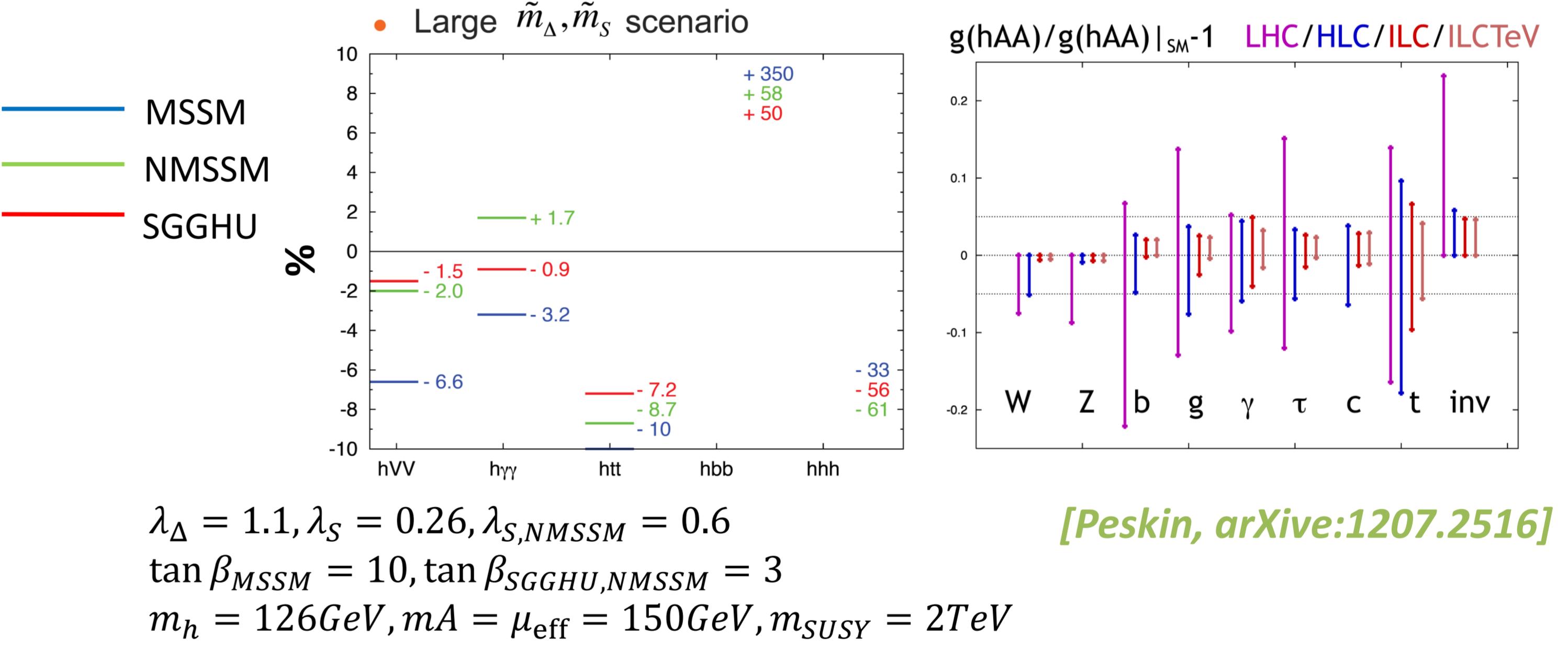
contribution from MSSM
contribution from singlet and triplet



Light CP-even mass reaches a 126GeV

c) Coupling of light Higgs boson

$$\frac{g(hAA)}{g(hAA)_{SM}} - 1$$



We can distinguish models using precision measurement at the ILC

b) Deviation from the MSSM

● Charged Higgs boson

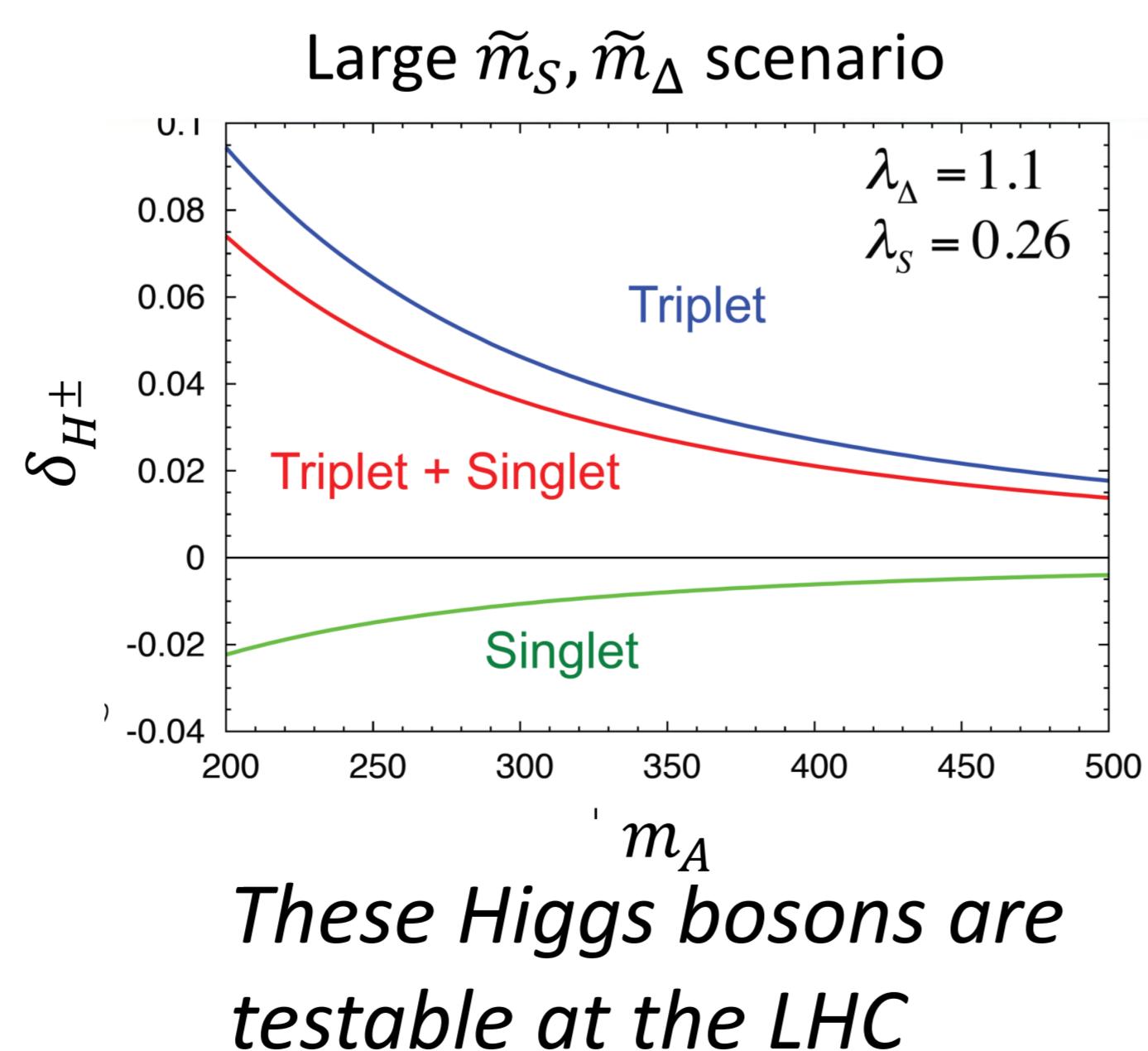
$$m_{H^\pm}^2 = m_{H^\pm}^2 \Big|_{MSSM} (1 + \delta_{H^\pm})^2 \sim m_A^2 + m_W^2 - \frac{1}{2} \lambda_s^2 v^2 + \frac{1}{8} \lambda_\Delta^2 v^2$$

Mass shifts: $O(1)\% - O(10)\%$

● Heavy CP-even Higgs boson

$$m_H \approx m_H \Big|_{MSSM} \approx m_H \Big|_{NMSSM}$$

Mass shifts is very small $< O(1)\%$

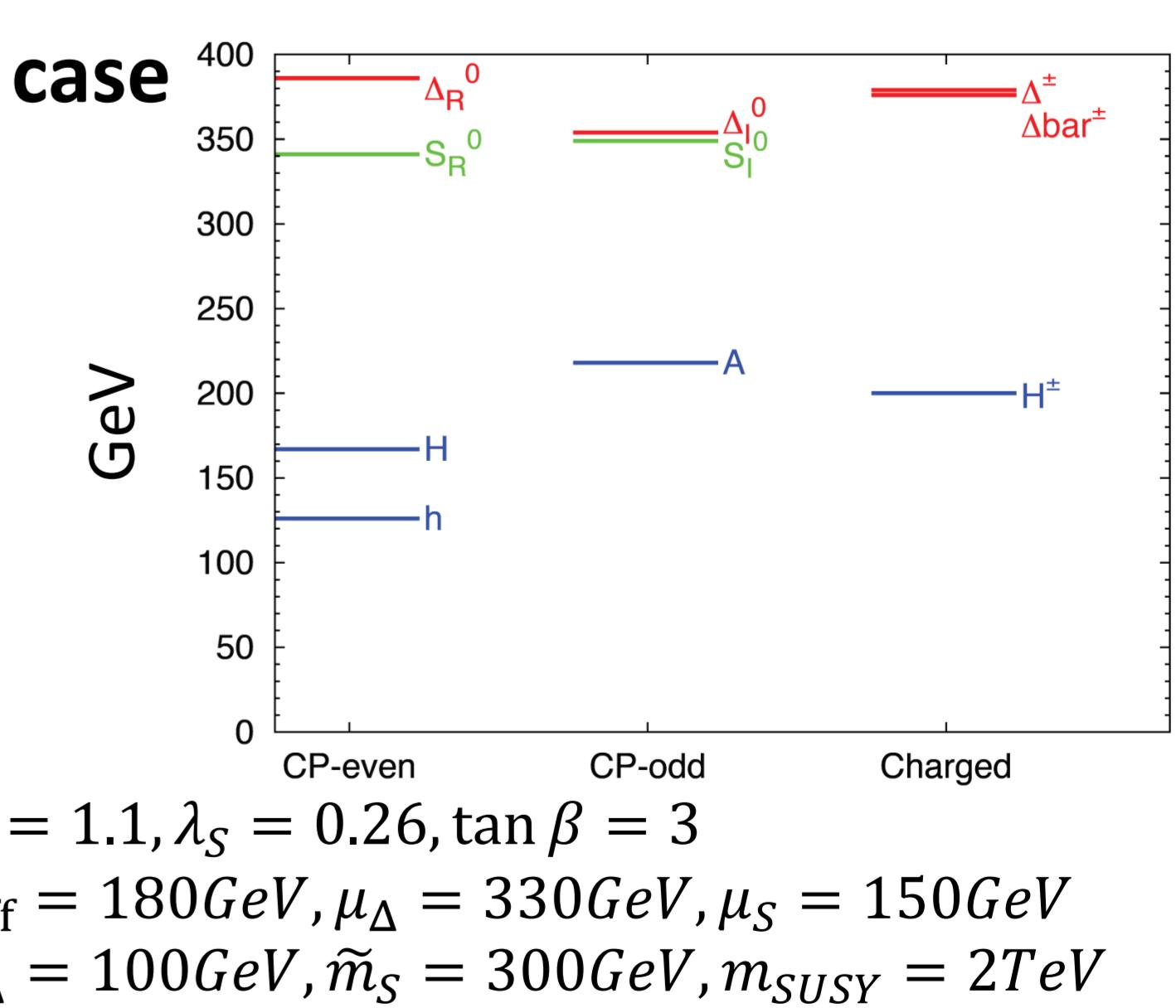


d) Masses of the Higgs bosons for light S, Δ scenario

● Small soft triplet and singlet mass case

In this case, we can directly produce additional Higgs bosons at the ILC

$$\text{ex)} e^+ e^- \rightarrow \Delta^+ \Delta^- \rightarrow tb tb \text{ via large mixing } H^\pm - \Delta^\pm$$



4, Conclusion

The predicted values of the Higgs sector deviates from the MSSM and the SM by $O(1) - O(10)\%$

We can test the SGGHU with two steps at the collider experiment. First, we will test charged Higgs at the LHC. Then, we will test the coupling constants of light CP-even Higgs at the ILC

Supersymmetric grand unified theory with the Hosotani mechanism is a good example of a GUT verifiable at colliders