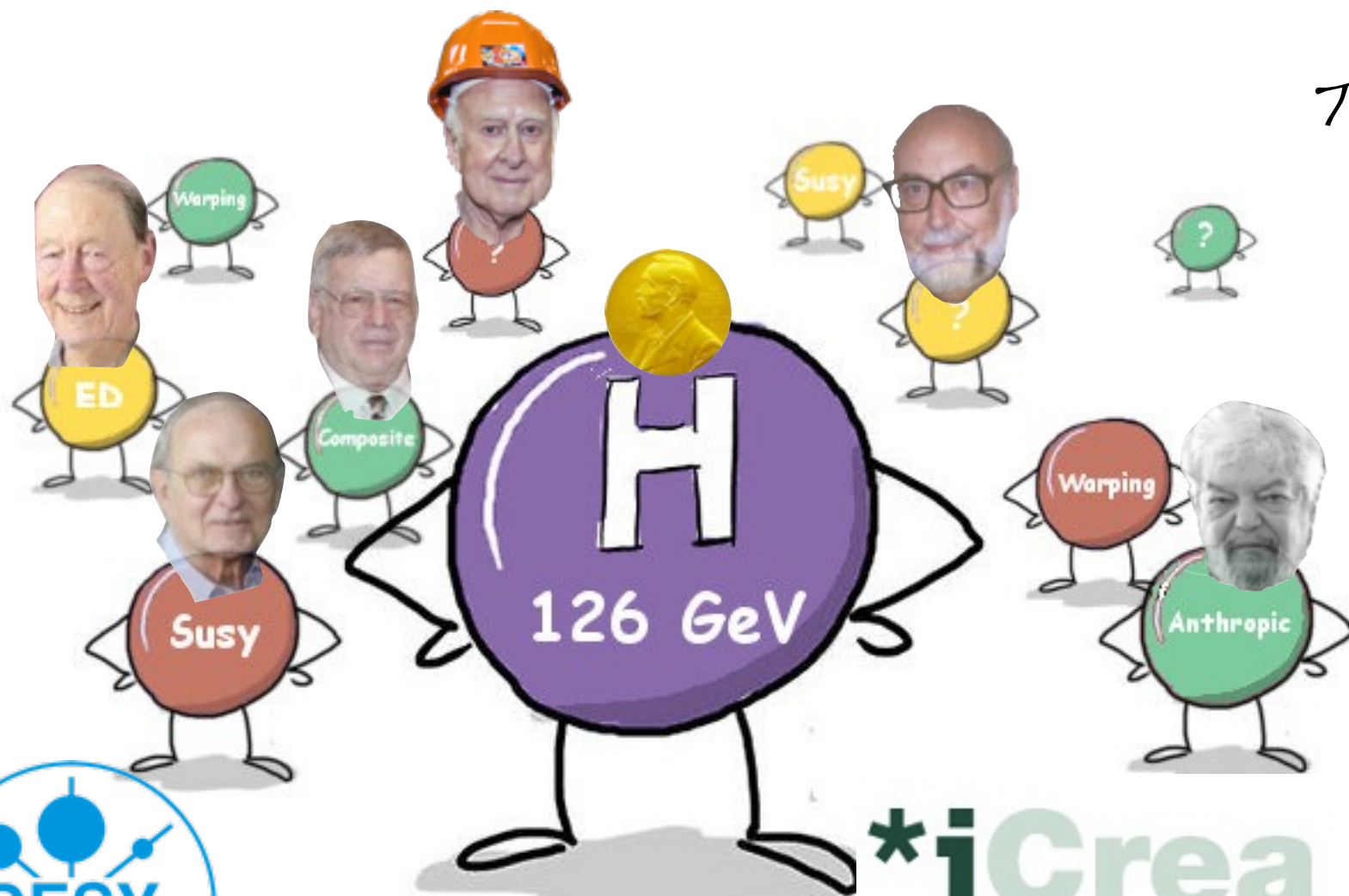


Could the Higgs be composite?

HPNP 2015

Toyama Univ. Feb. 12, 2015



P. Cámara/C. Grojean

***iCrea**

INSTITUCIÓ CATALANA DE
RECERCA I ESTUDIS AVANÇATS

Christophe Grojean

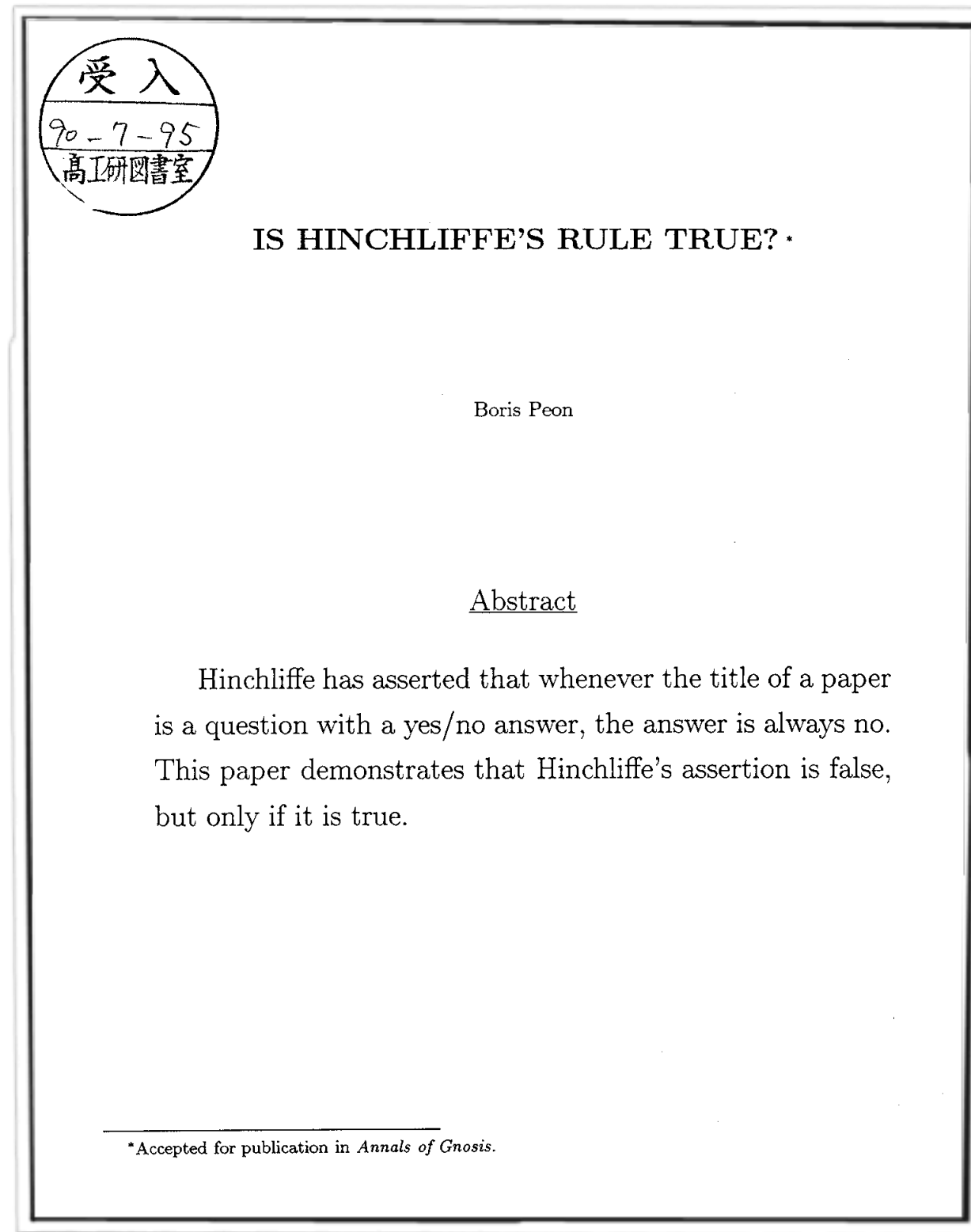
DESY (Hamburg)

ICREA@IFAE (Barcelona)

(christophe.grojean@cern.ch)



Remember Hinchliffe's rule...



However, as physicists, we want to base our answers on experimental data and to keep challenging theoretical prejudices...

We all have a PhD

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For the first time in the history of physics,
we have a *consistent* description of the fundamental constituents of matter and their interactions and this description can be extrapolated to very high energy (up M_{Planck} ?)

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My key message MLM@Aspen'14

- The days of “guaranteed” discoveries or of no-lose theorems in particle physics are over, at least for the time being
- but the big questions of our field remain wild open (hierarchy problem, flavour, neutrinos, DM, BAU,)
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Where and how does the SM break down?
Which machine(s) will reveal this breakdown?

Why should you care about compositeness?

Higgs compositeness means new fundamental interactions

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Pospelov's 38 years rule...

38 years rule = new forces of nature are discovered every 38 years for the last 150 yrs

1. 1860s – first papers of Maxwell on EM. Light is EM excitation. E & M unification.
2. 1897 – Becquerel discovers radioactivity – first evidence of weak charged currents (in retrospect).
3. 1935 – Chadwick gets NP for his discovery of neutron with subsequent checks that there exists strong n-p interaction. Strong force is established.
4. 1973 – Gargamelle experiment sees the evidence for weak neutral currents in ν -N scattering
5. 2011/2012 Discovery of the Higgs, i.e. new Yukawa force.
6. *Prediction: Discovery of a new dark force – 2050?*

(+/- 2 years or so).

M. Pospelov, SHiP collab. meeting, Naples '15

Why should you care about compositeness?

All SM shortcomings are intimately linked to the Higgs elementary nature

$$\mathcal{L}_{\text{Higgs}} = V_0 - \mu^2 H^\dagger H + \lambda (H^\dagger H)^2 + (y_{ij} \bar{\psi}_{Li} \psi_{Rj} H + h.c.)$$

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vacuum energy

cosmological constant

$$V_0 \approx (2 \times 10^{-3} \text{ eV})^4 \ll M_{\text{PL}}^4$$

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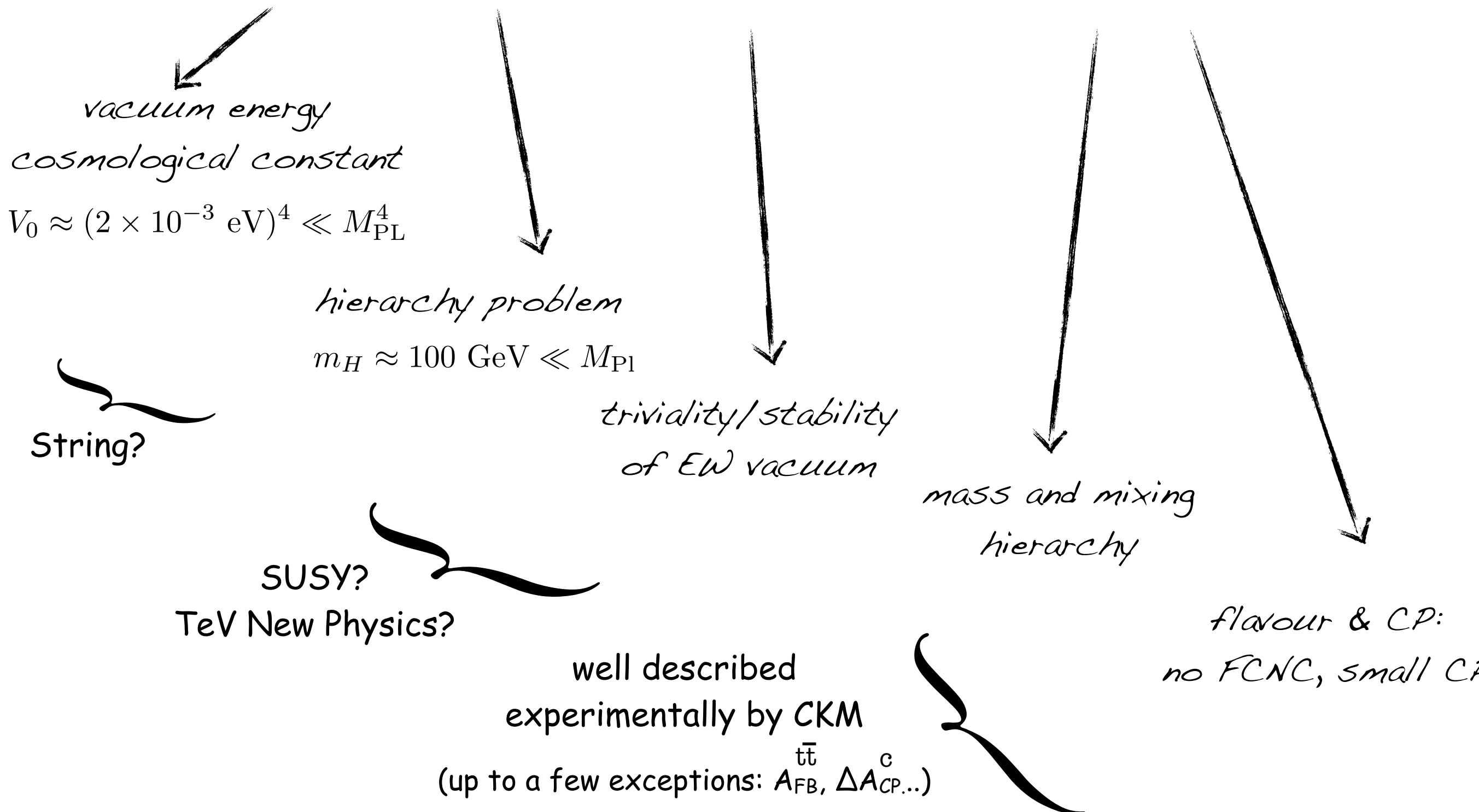


flavour & CP:
no FCNC, small CP

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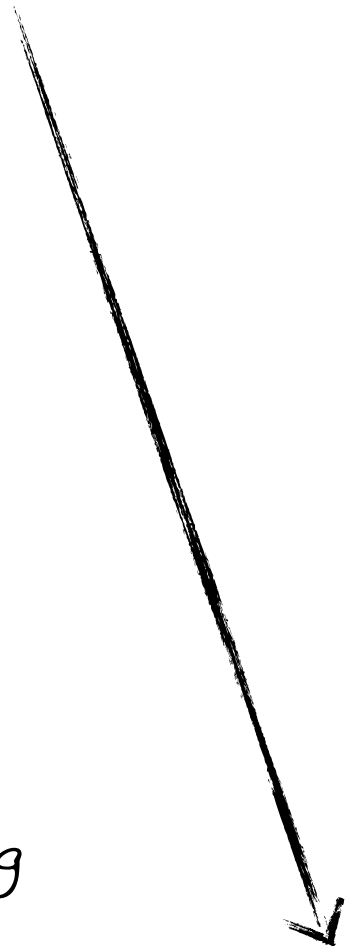
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All these problems because the Higgs boson would be the first elementary particle whose interactions are not endowed with a gauge structure

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Higgs = Elementary or Composite?

Probing the Higgs compositeness

Unlikely we'll ever see the fundamental constituents of the Higgs

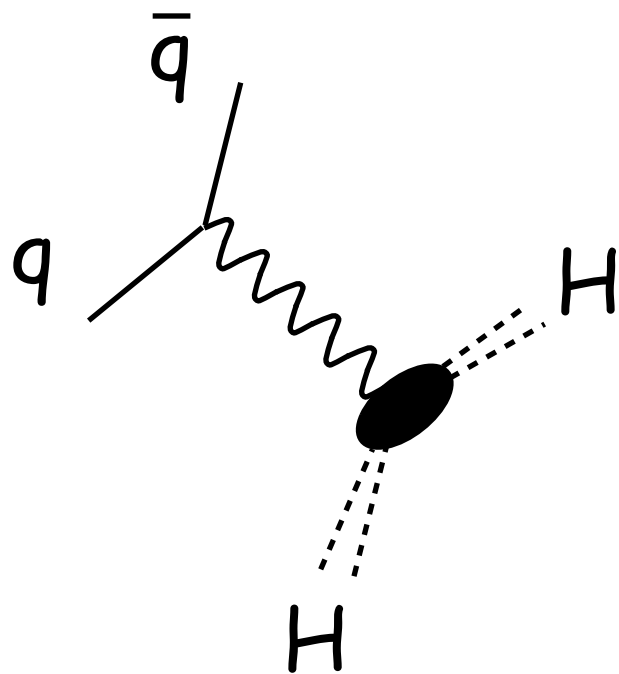
But we can infer that it is not an elementary particle

by measuring its couplings to SM particles

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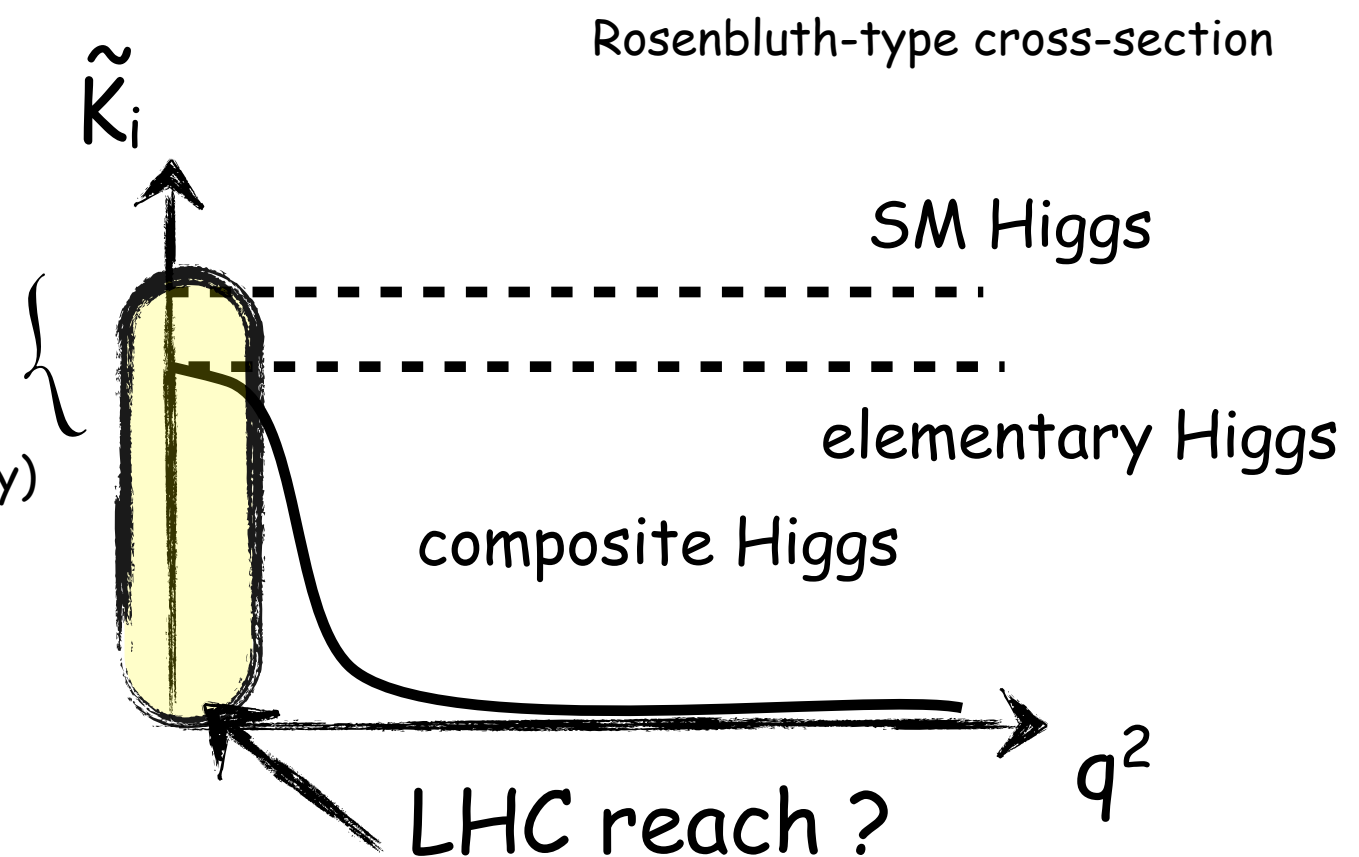
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But we can infer that it is not an elementary particle
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$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{16m_H^2 \sin^4 \theta/2} \frac{E'}{E^3} \left(2\tilde{K}_1 q^2 \sin^2 \theta/2 + \tilde{K}_2 \cos^2 \theta/2 \right)$$

anomalous couplings
(accessible @ LHC with 20-10% accuracy)



Which composite scenario?

Minimal Composite Higgs

ex: $SO(5)/SO(4)$

SILH

$$\xi = \frac{v^2}{f^2} \ll 1$$

Strong sector
 (g^*, f)
PNGB Higgs

$$g_{SM}^2 / g^*$$

SM
 (g, g', y_+)

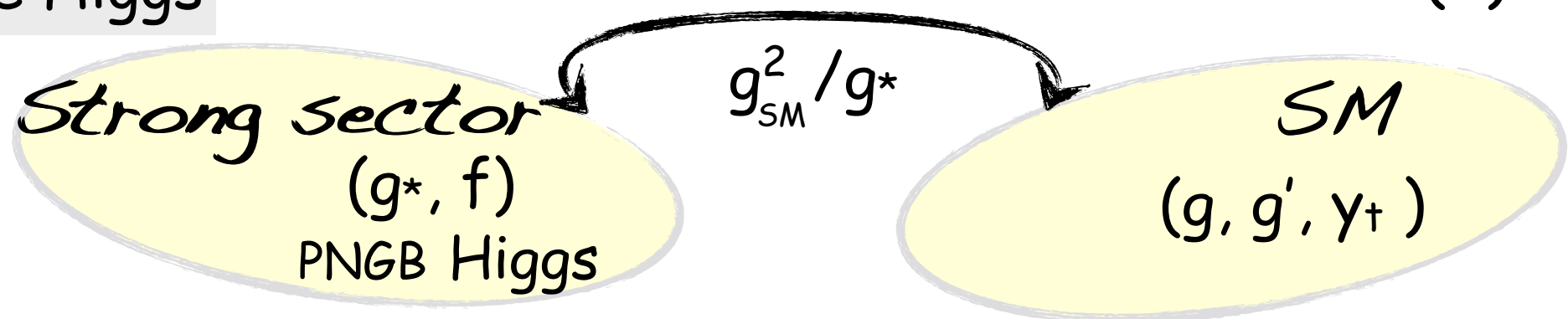
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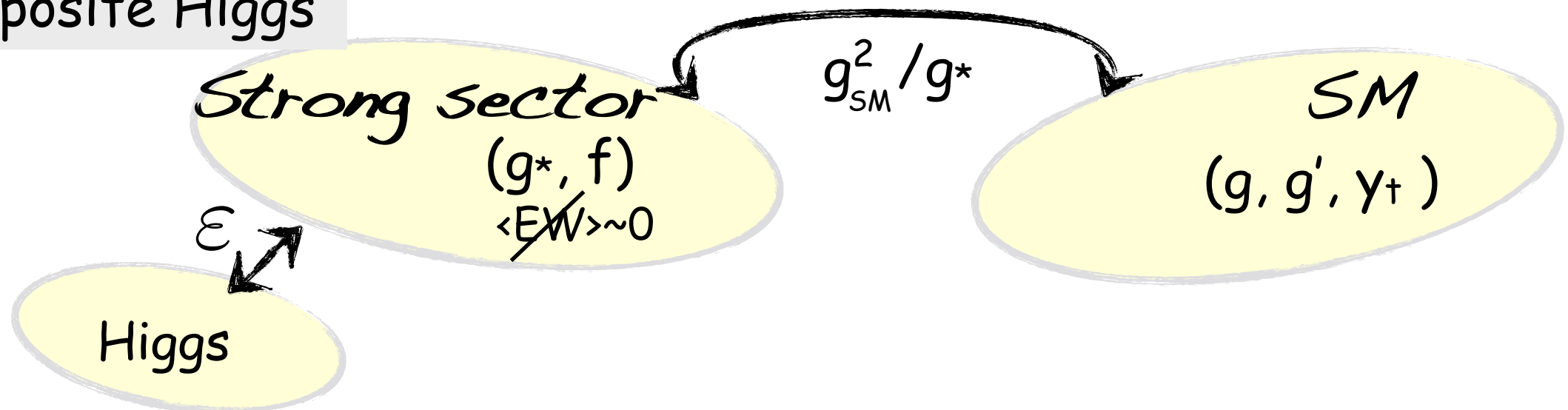
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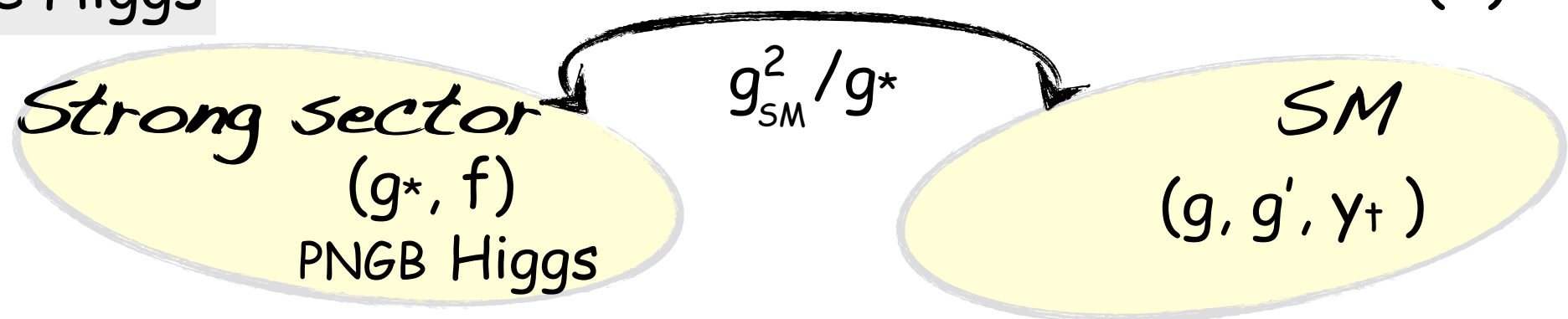
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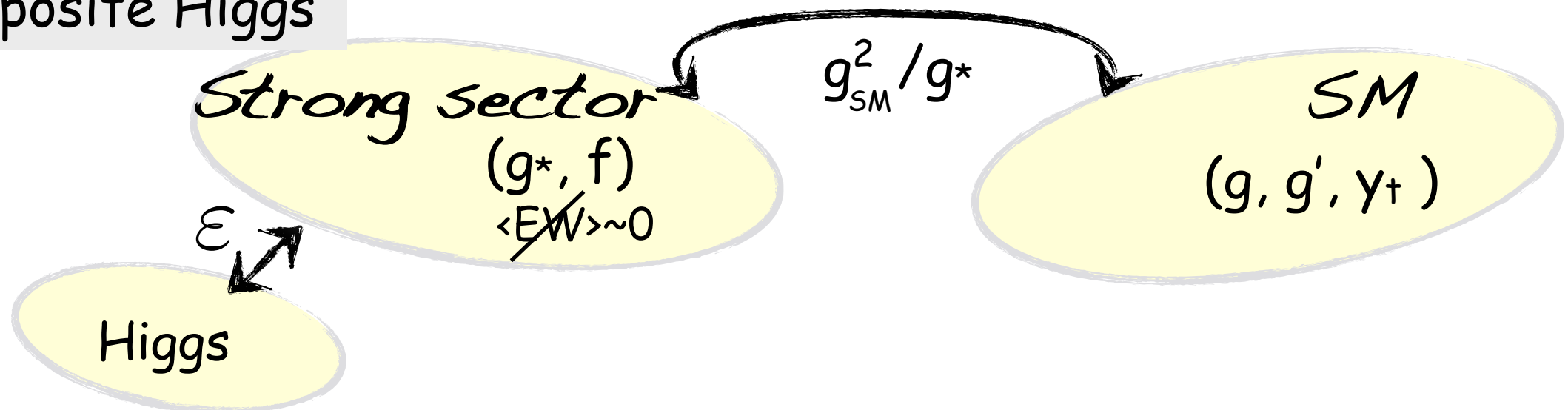
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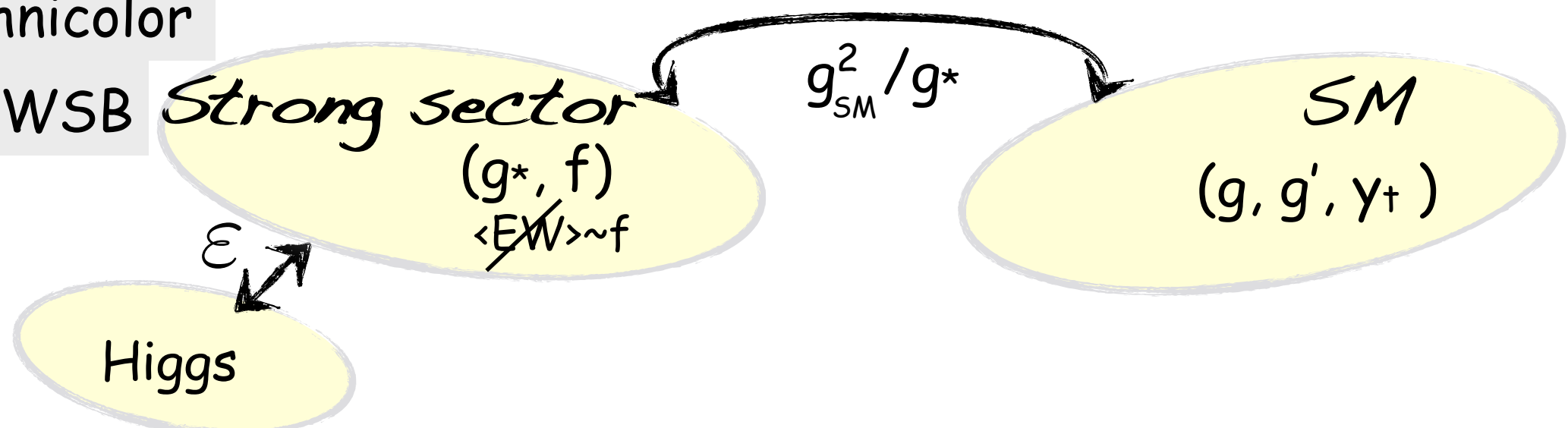
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Bosonic Technicolor

Induced EWSB

$$\epsilon = \frac{f}{v} \ll 1$$



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Minimal Composite Higgs

SILH

$$\xi = \frac{v^2}{f^2} \ll 1$$

$$\frac{1}{f^2} (\partial_\mu |H|^2)^2$$

$$\frac{\lambda_4}{f^2} |H|^6$$

$$\kappa_V \equiv \frac{g_{hVV}}{g_{hVV}^{\text{SM}}} = 1 + \xi$$

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Patterns of Higgs coupling deviations

expected largest relative deviations

	hff	hVV	$h\gamma\gamma$	$h\gamma Z$	hGG	h
MSSM	✓		✓	✓	✓	
NMSSM	✓	✓	✓	✓	✓	
PGB Composite	✓	✓		✓		✓
SUSY Composite	✓	✓	✓	✓	✓	✓
SUSY partly-composite			✓	✓	✓	✓
“Bosonic TC”						✓
Higgs as a dilaton			✓	✓	✓	✓

see also
Muehlleitner's talk

A. Pomarol, Naturalness '15

What is the Higgs the name of?

The SM Higgs couplings are fixed to restore unitarity with mass

$$\Sigma = e^{i\sigma^a \pi^a / v} \quad \text{Goldstone of } SU(2)_L \times SU(2)_R / SU(2)_V \quad D_\mu \Sigma = g V_\mu$$

$$\mathcal{L}_{\text{EWSB}} = \frac{v^2}{4} \text{Tr} (D_\mu \Sigma^\dagger D_\mu \Sigma) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} \right) - \lambda \bar{\psi}_L \Sigma \psi_R \left(1 + c \frac{h}{v} \right)$$

'a', 'b' and 'c' are arbitrary free couplings

For $a=1$: perturbative unitarity in elastic channels $WW \rightarrow WW$

For $b=a^2$: perturbative unitarity in inelastic channels $WW \rightarrow hh$

For $ac=1$: perturbative unitarity in inelastic $WW \rightarrow \psi \psi$

Cornwall, Levin, Tiktopoulos '73

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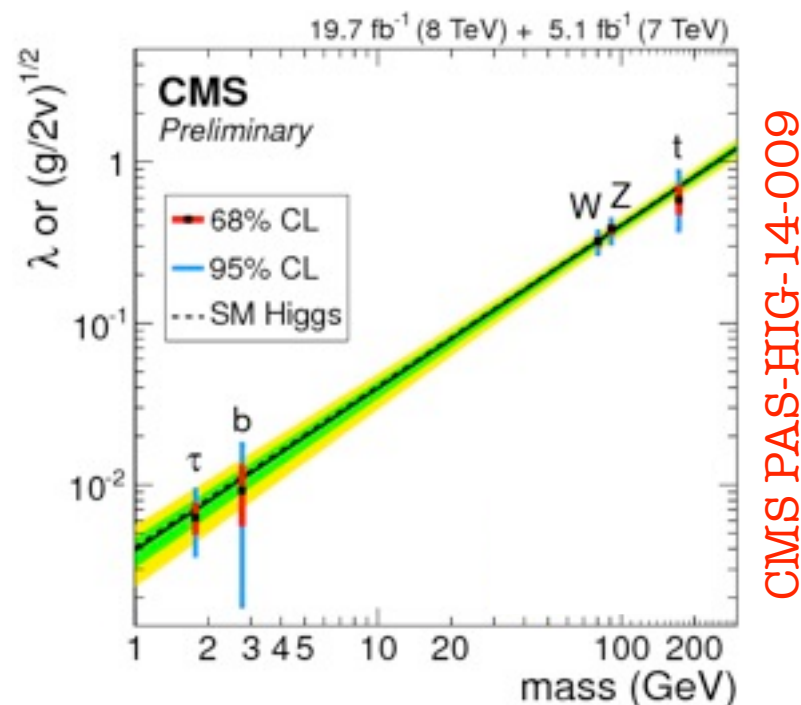
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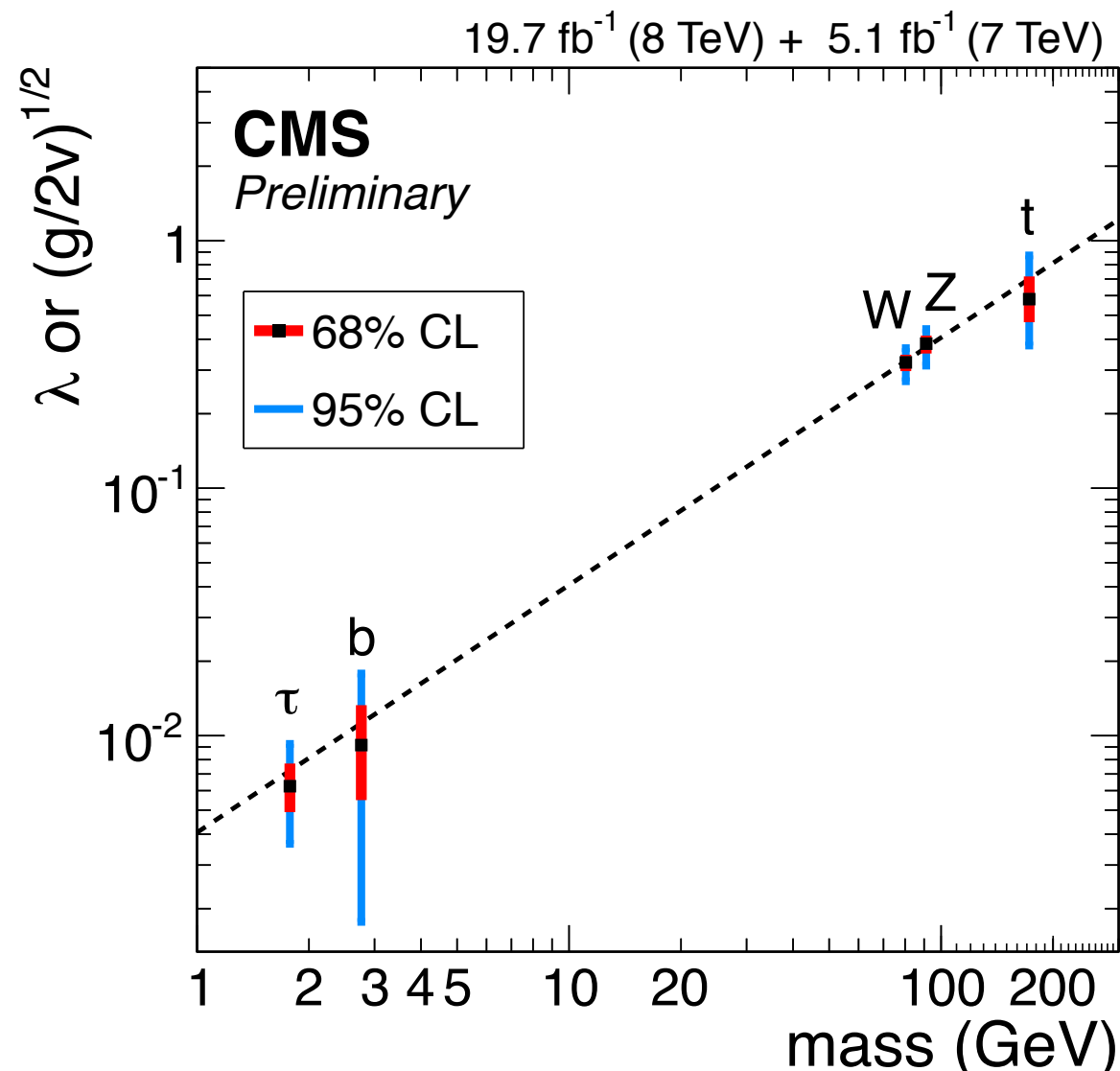


Higgs couplings
are proportional
to the masses of the particles

$$\lambda_\psi \propto \frac{m_\psi}{v}, \quad \lambda_V^2 \equiv \frac{g_{VVh}}{2v} \propto \frac{m_V^2}{v^2}$$

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Higgs group @ Snowmass '13

Facility	LHC	HL-LHC
\sqrt{s} (GeV)	14,000	14,000
$\int \mathcal{L} dt$ (fb ⁻¹)	300/expt	3000/expt
κ_γ	5 – 7%	2 – 5%
κ_g	6 – 8%	3 – 5%
κ_W	4 – 6%	2 – 5%
κ_Z	4 – 6%	2 – 4%
κ_ℓ	6 – 8%	2 – 5%
$\kappa_d = \kappa_b$	10 – 13%	4 – 7%
$\kappa_u = \kappa_t$	14 – 15%	7 – 10%

~ Is this fit theoretically consistent? ~

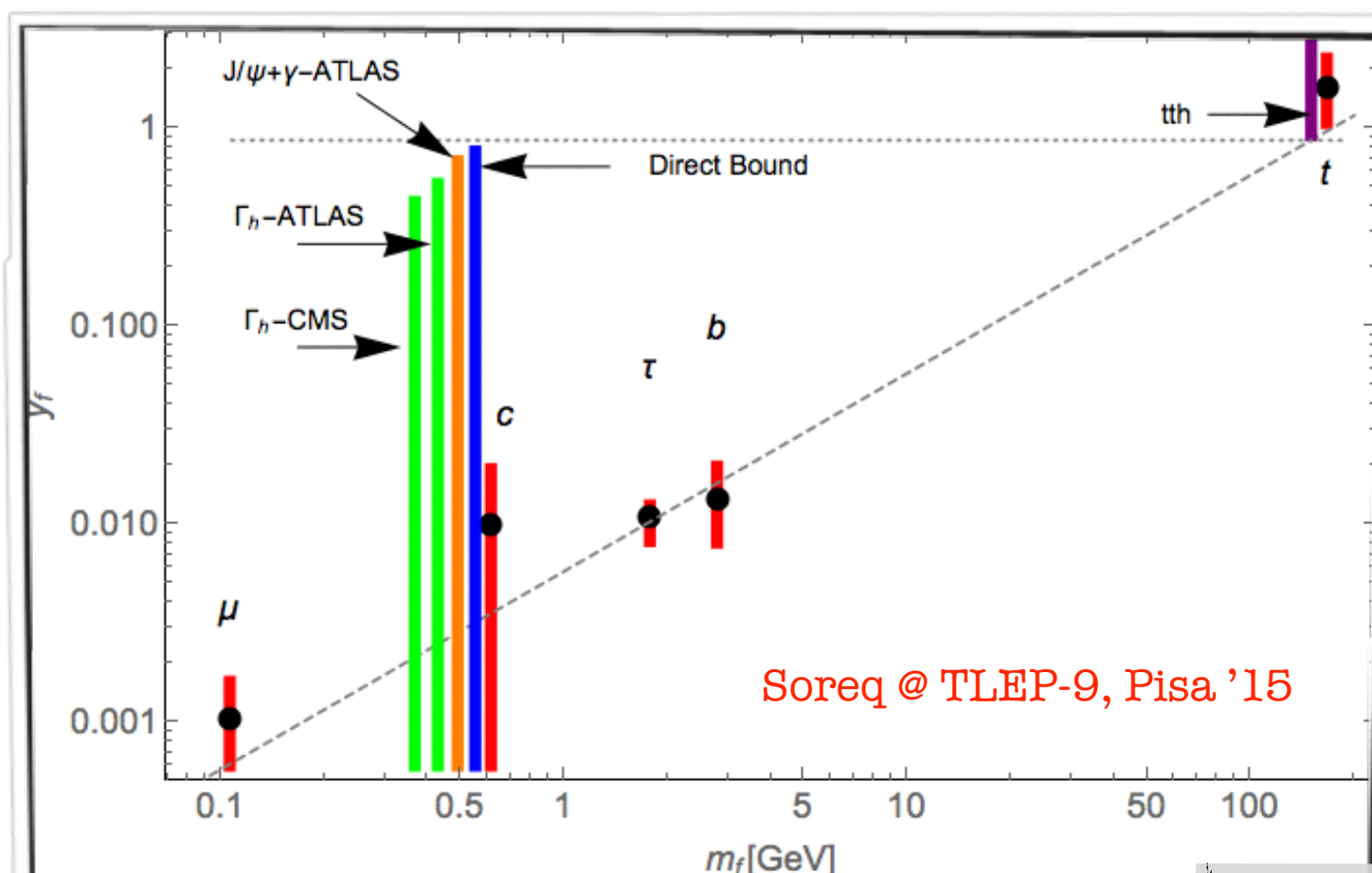
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missing information to complete the picture

◦ width measurement?

◦ couplings to light particles?

inclusive (e.g. c-tagging) or exclusive ($h \rightarrow J/\Psi + \gamma$)

◦ coupling to top?

known indirectly ($gg \rightarrow h$) or via difficult tth channel

~ Is this fit theoretically consistent? ~

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in the bottom coupling without generating other coupling structures not taken into account in the fit?



Precision program in single Higgs processes


(assuming a mass gap between weak scale and new physics scale)

Higgs/BSM Primaries

Several deformations away from the SM are harmless in the vacuum and need a Higgs field to be probed

e.g.
$$\frac{1}{g_s^2} G_{\mu\nu}^2 + \frac{|H|^2}{\Lambda^2} G_{\mu\nu}^2 \rightarrow \left(\frac{1}{g_s^2} + \frac{v^2}{\Lambda^2} \right) G_{\mu\nu}^2$$

operator is not visible in the vacuum
(redefinition of input parameter)



But can affect h physics:



Higgs/BSM Primaries

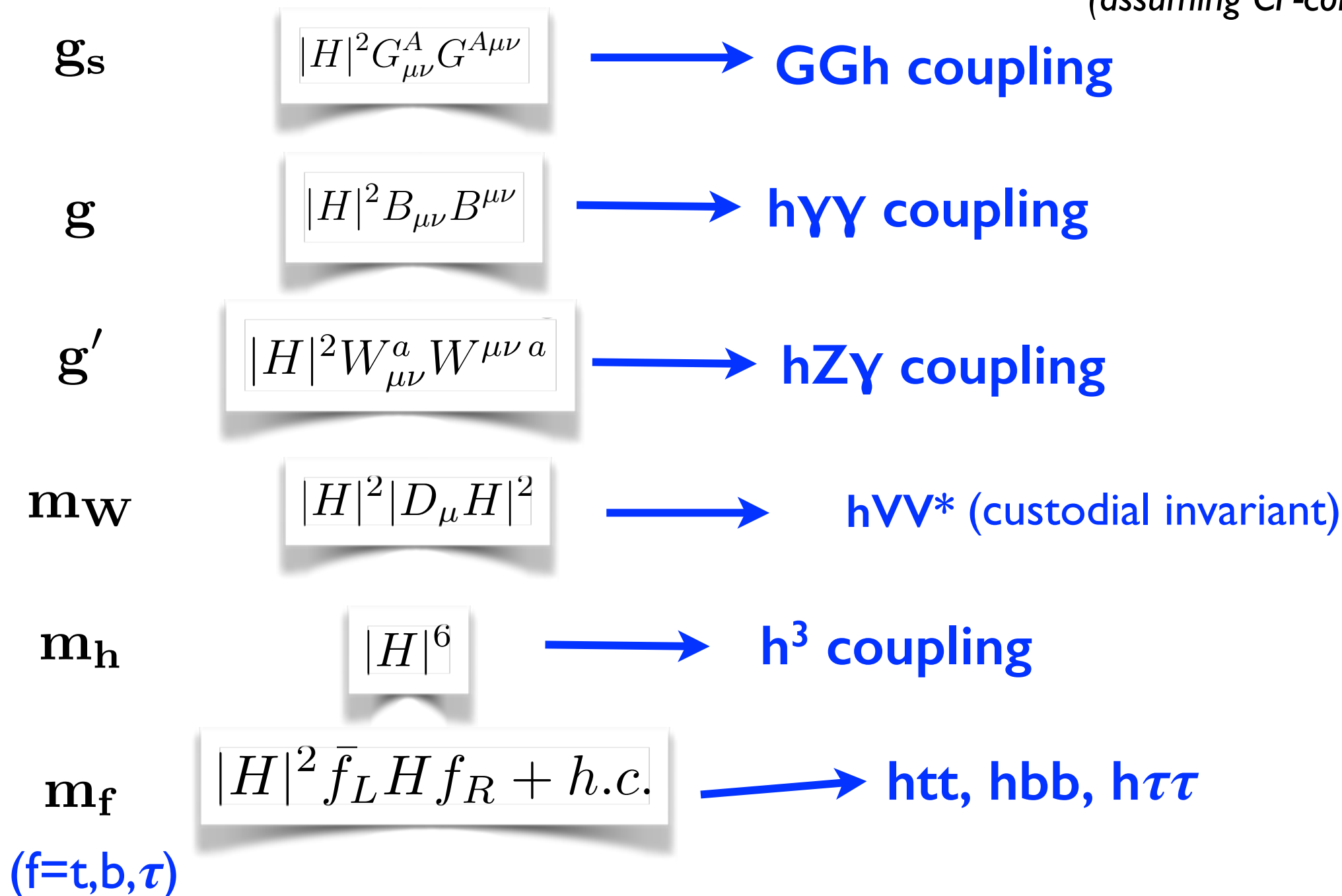
How many of these effects can we have?

Pomarol, Riva '13

Elias-Miro et al '13

Gupta, Pomarol, Riva '14

As many as parameters in the SM: **8** for one family
(assuming CP-conservation)



(courtesy of A. Pomarol@HiggsHunting2014)

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g_s

$$|H|^2 G_{\mu\nu}^A G^{A\mu\nu}$$

→ **GGh coupling**

g

$$|H|^2 B_{\mu\nu} B^{\mu\nu}$$

→ **h $\gamma\gamma$ coupling**

yet to be measured
at the LHC

g'

$$|H|^2 W_{\mu\nu}^a W^{a\mu\nu}$$

→ **hZ γ coupling**

m_W

$$|H|^2 |D_\mu H|^2$$

→ **hVV* (custodial invariant)**

m_h

$$|H|^6$$

→ **h³ coupling**

m_f

$$|H|^2 \bar{f}_L H f_R + h.c.$$

→ **htt, hbb, h $\tau\tau$**

(f=t,b, τ)

the 6 others have been measured (~15%) up to a flat direction
between the top/gluon/photon couplings

(courtesy of A. Pomarol@HiggsHunting2014)

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Almost a 1-to-1 correspondence
with the 8 κ 's in the Higgs fit

Coupling	300 fb ⁻¹ Theory unc.:			3000 fb ⁻¹ Theory unc.:		
	All	Half	None	All	Half	None
κ_Z	8.1%	7.9%	7.9%	4.4%	4.0%	3.8%
κ_W	9.0%	8.7%	8.6%	5.1%	4.5%	4.2%
κ_t	22%	21%	20%	11%	8.5%	7.6%
κ_b	23%	22%	22%	12%	11%	10%
κ_τ	14%	14%	13%	9.7%	9.0%	8.8%
κ_μ	21%	21%	21%	7.5%	7.2%	7.1%
κ_g	14%	12%	11%	9.1%	6.5%	5.3%
κ_γ	9.3%	9.0%	8.9%	4.9%	4.3%	4.1%
$\kappa_{Z\gamma}$	24%	24%	24%	14%	14%	14%

Atlas projection

With some important differences:

- 1) width approximation built-in
- 2) κ_W/κ_Z is not a primary
(constrained by $\Delta\rho$ and TGC)
- 3) $\kappa_g, \kappa_\gamma, \kappa_{Z\gamma}$ do not separate UV and IR contributions

8

for one family

(assuming CP-conservation)

GGh coupling

h $\gamma\gamma$ coupling

yet to be measured
at the LHC

hZ γ coupling

hVV* (custodial invariant)

h³ coupling

htt, hbb, h $\tau\tau$

Don't forget LEP!

The parameter 'a' controls the size of the one-loop IR contribution to the LEP precision observables

$$\mathcal{L} \supset \frac{1}{f^2} |H|^2 |D_\mu H|^2$$

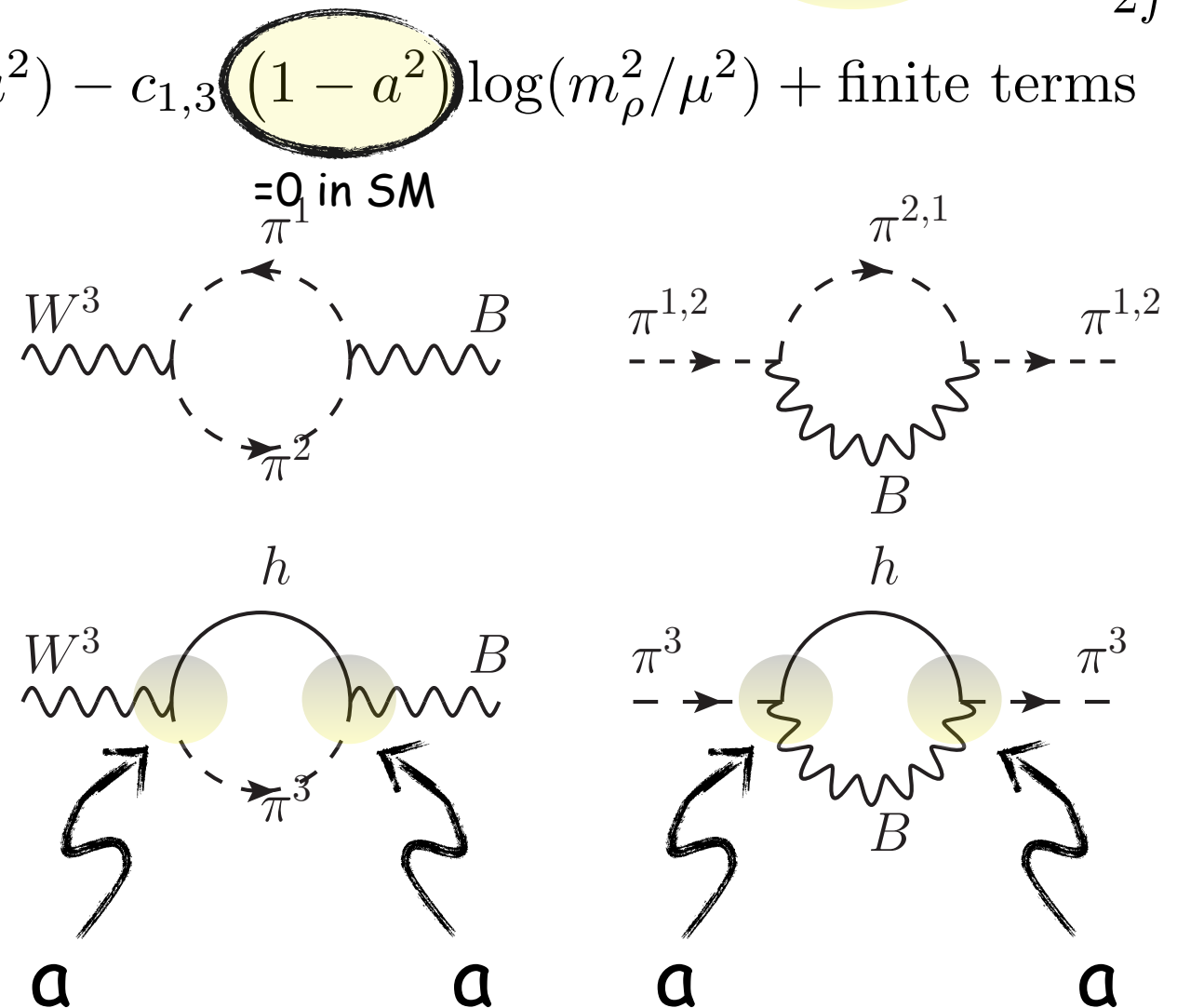
$$\Rightarrow a = \kappa_V = 1 + \frac{v^2}{2f^2}$$

$$\epsilon_{1,3} = c_{1,3} \log(m_Z^2/\mu^2) - c_{1,3} a^2 \log(m_h^2/\mu^2) - c_{1,3} (1 - a^2) \log(m_\rho^2/\mu^2) + \text{finite terms}$$

$$c_1 = + \frac{3}{16\pi^2} \frac{\alpha(m_Z)}{\cos^2 \theta_W} \quad c_3 = - \frac{1}{12\pi} \frac{\alpha(m_Z)}{4 \sin^2 \theta_W}$$

$$\Delta\epsilon_{1,3} = -c_{1,3} (1 - a^2) \log(m_\rho^2/m_h^2)$$

Barbieri, Bellazzini, Rychkov, Varagnolo '07



Log. div. cancel only for $a=1$ (SM)
 $a \neq 1$ log. sensitivity on the scale of new physics

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The parameter 'a' controls the size of the one-loop IR contribution to the LEP precision observables

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Barbieri, Bellazzini, Rychkov, Varagnolo '07

EW fit:

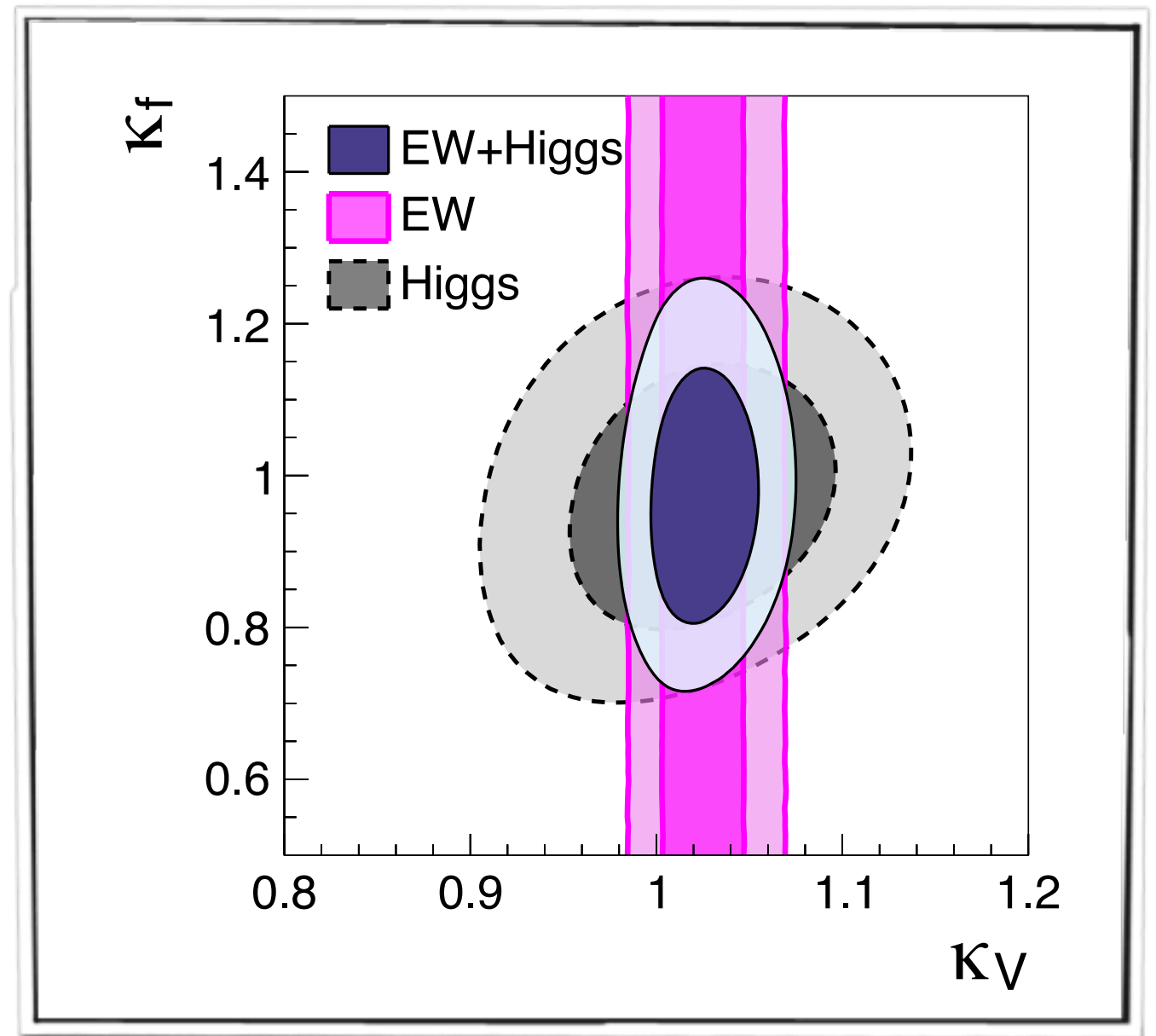
$$0.98 \leq a^2 \leq 1.12$$

Ciuchini et al '13

see also Grojean et al '13

The LEP indirect constraints on the other BSM primaries are not competitive

Elias-Miro et al '13



Ciuchini et al '13

CP violation in Higgs physics?

Is CP a good symmetry of Nature? 2 CP-violating couplings in the SM:

V_{CKM} (large, $O(1)$), but screened by small quark masses) and θ_{QCD} (small, $O(10^{-10})$)

Can the 0^+ SM Higgs boson have CP violating couplings?

Among the 59 irrelevant directions, 6 ~~CP~~ Higgs/BSM primaries

$$\begin{aligned} \Delta\mathcal{L}_{BSM} = & i\delta\tilde{g}_{hff} h\bar{f}_L f_R + h.c. & (f=b, \tau, t) \\ & + \tilde{\kappa}_{GG} \frac{h}{v} G^{\mu\nu} \tilde{G}_{\mu\nu} & (\tilde{F}_{\mu\nu} \equiv \epsilon_{\mu\nu\rho\sigma} F^{\rho\sigma}) \\ & + \tilde{\kappa}_{\gamma\gamma} \frac{h}{v} F^{\gamma\mu\nu} \tilde{F}_{\mu\nu} \\ & + \tilde{\kappa}_{\gamma Z} \frac{h}{v} F^{\gamma\mu\nu} \tilde{F}_{\mu\nu}^Z \end{aligned}$$

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Among the 59 irrelevant directions, 6 ~~CP~~ Higgs/BSM primaries

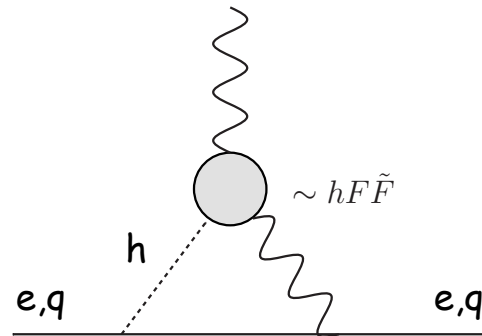
$$\begin{aligned}\Delta\mathcal{L}_{\text{BSM}} = & i\delta\tilde{g}_{hff} h\bar{f}_L f_R + h.c. & (f=b, \tau, t) \\ & + \tilde{\kappa}_{GG} \frac{h}{v} G^{\mu\nu} \tilde{G}_{\mu\nu} & (\tilde{F}_{\mu\nu} \equiv \epsilon_{\mu\nu\rho\sigma} F^{\rho\sigma}) \\ & + \tilde{\kappa}_{\gamma\gamma} \frac{h}{v} F^{\gamma\mu\nu} \tilde{F}_{\mu\nu} \\ & + \tilde{\kappa}_{\gamma Z} \frac{h}{v} F^{\gamma\mu\nu} \tilde{F}_{\mu\nu}^Z\end{aligned}$$

see also
Boudjema's talk

operators with γ :

already severely constrained
by e and q EDMs

McKeen, Pospelov, Ritz '12



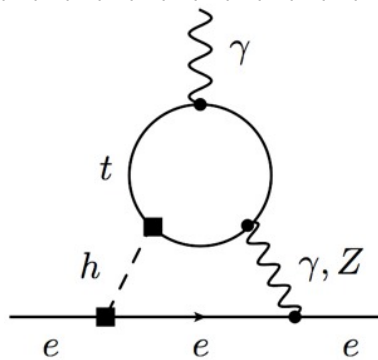
$$\tilde{\kappa}_{\gamma\gamma} \sim \tilde{\kappa}_{\gamma Z} \leq 10^{-4}$$

$$\Lambda_{\text{CP}} > 25 \text{ TeV}$$

operators with top:

already severely constrained
by e and q EDMs

Brod, Haisch, Zupan '13



$$\delta\tilde{g}_{htt} \leq 0.01$$

$$\Lambda_{\text{CP}} > 2.5 \text{ TeV}$$

Caveats: h couplings to light particles can be significantly reduced



Boosted and off-shell Higgs channels

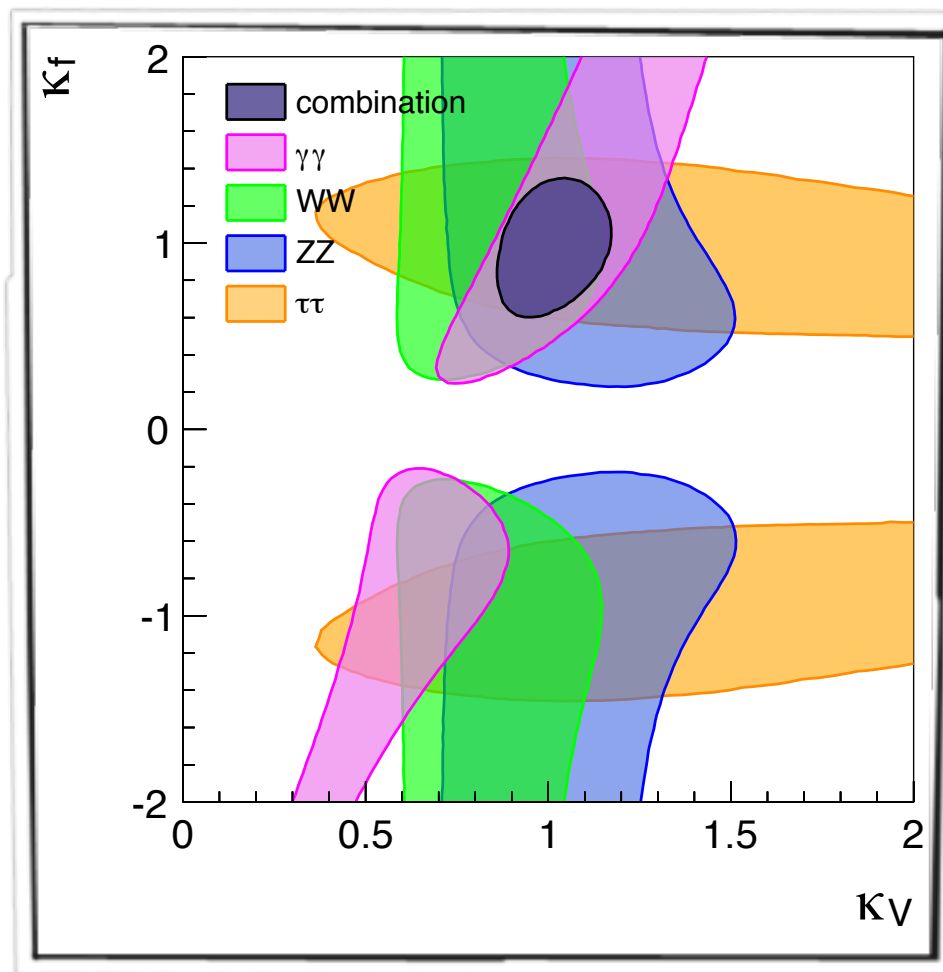
Why going beyond inclusive Higgs processes?

So far the LHC has mostly produced Higgses on-shell
in processes with a characteristic scale $\mu \approx m_H$

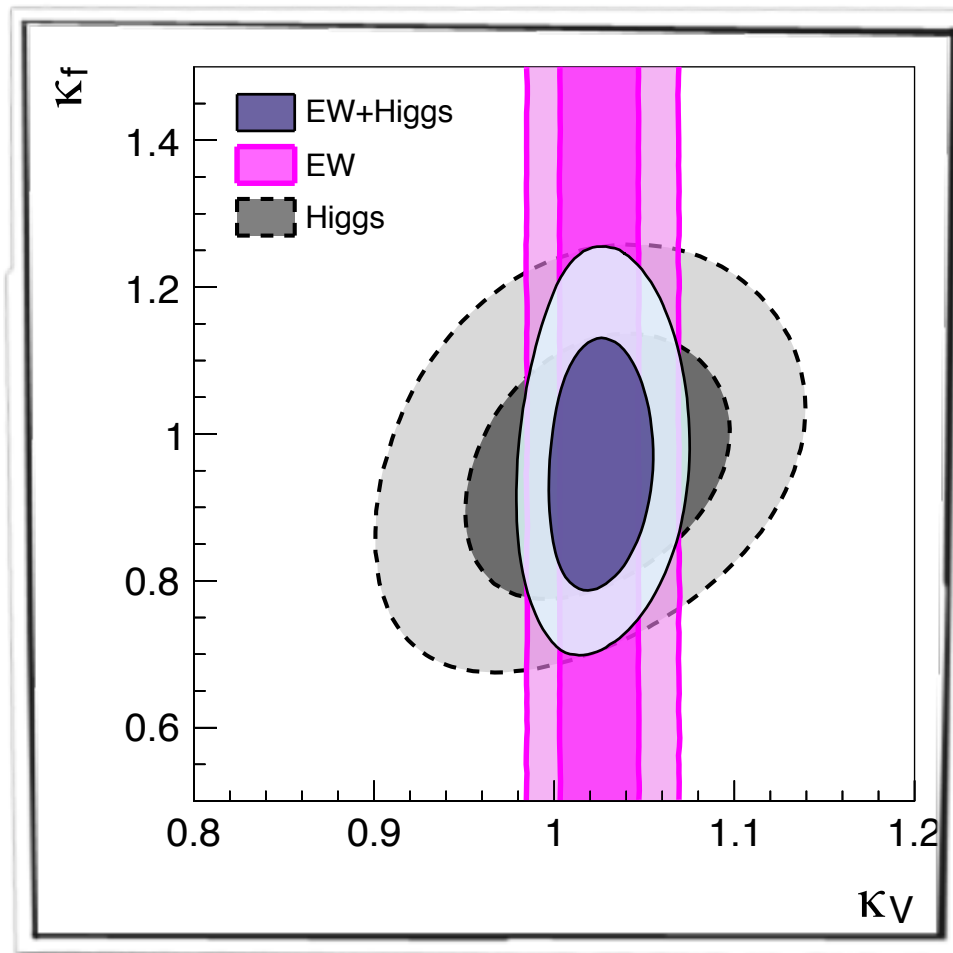
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↓ ↓
access to Higgs couplings @ m_H



Ciuchini et al '13



Ciuchini et al '13

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access to Higgs couplings @ m_H

Producing a Higgs with boosted additional particle(s)
probe the Higgs couplings @ large energy
(important to check that the Higgs boson ensures perturbative unitarity)

Probing new corrections to the SM Lagrangian?

on-shell Z @ LEP1

constraints on
S and T oblique corrections

off-shell Z @ LEP2

constraints on
W and Y oblique corrections
(same order as S and T but cannot be probed @ LEP1)

But... off-shell Higgs data do not probe new corrections
that cannot be constrained by on-shell data

Boosted Higgs

inability to resolve the top loops

- the bearable lightness of the Higgs: rich spectroscopy w/ multiple decays channels
- the unbearable lightness: loops saturate and don't reveal the physics @ energy physics ^(*)

$m_H(\text{GeV})$	$\frac{\sigma_{NLO}(m_t)}{\sigma_{NLO}(m_t \rightarrow \infty)}$	$\frac{\sigma_{NLO}(m_t, m_b)}{\sigma_{NLO}(m_t \rightarrow \infty)}$
125	1.061	0.988
150	1.093	1.028
200	1.185	1.134

e.g. Grazzini, Sargsyan '13



the inclusive rate
doesn't "see" the finite mass of the top

^(*) unless it doesn't decouple
(e.g. 4th generation)

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the inclusive rate
doesn't "see" the finite mass of the top

⇒ cannot disentangle ○ long distance physics (modified top coupling)
○ short distance physics (new particles running in the loop) ⇐

$$\mathcal{L} = \frac{\alpha_s c_g}{12\pi} |H|^2 G_{\mu\nu}^a{}^2 + \frac{\alpha c_\gamma}{2\pi} |H|^2 F_{\mu\nu} + y_t c_t \bar{q}_L \tilde{H} t_R |H|^2$$

$$\frac{\sigma(gg \rightarrow h)}{\text{SM}} = (1 + (c_g - c_t)v^2)^2 \quad \frac{\Gamma(h \rightarrow \gamma\gamma)}{\text{SM}} = (1 + (c_\gamma - 4c_t/9)v^2)^2$$

fermionic top-partners in composite Higgs models exactly lead to $\Delta c_t = \Delta c_g = \frac{9}{4} \Delta c_\gamma$.

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see also
Boudjema's talk

having access to $h\bar{t}t$ final state will resolve this degeneracy
but notoriously difficult channel

14%-4% @ LHC¹⁴₃₀₀-LHC¹⁴₃₀₀₀ vs 10%-4% @ ILC⁵⁰⁰₅₀₀-ILC¹⁰⁰⁰₁₀₀₀

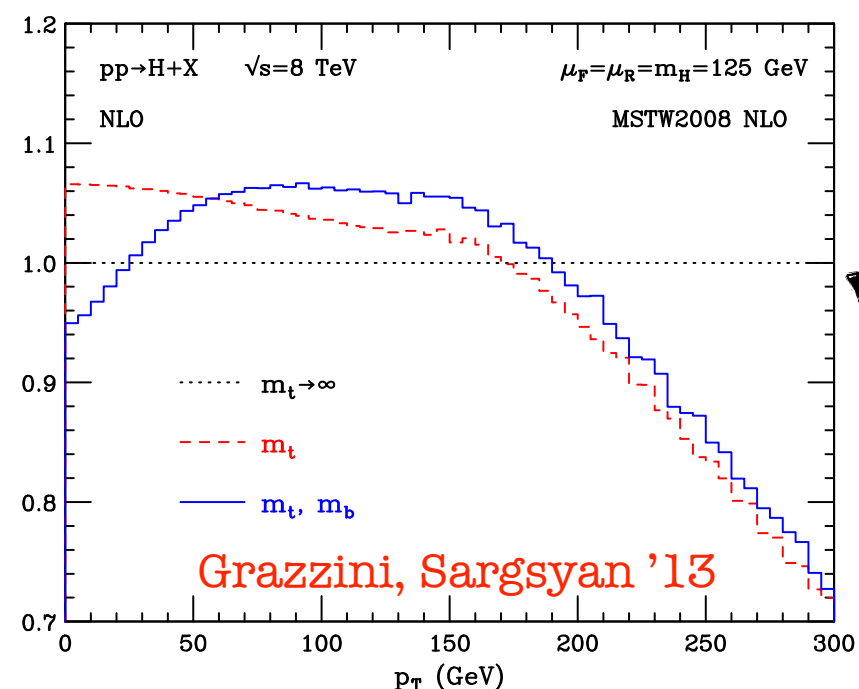
Resolving top loop: Boosted Higgs

cut open the top loops

high $p_T \approx$ Higgs off-shell
we "see" the details of the particles
running inside the loops

Baur, Glover '90

Langenegger, Spira, Starodumov, Trueb '06



Note: LO only

NLO_{m_t} is not known

$1/m_t$ corrections known $O(\alpha_s^4)$

few % up to $p_T \sim 150$ GeV

Harlander et al '12

the high p_T tail
is tens' % sensitive
to the mass of top

Resolving top loop: Boosted Higgs

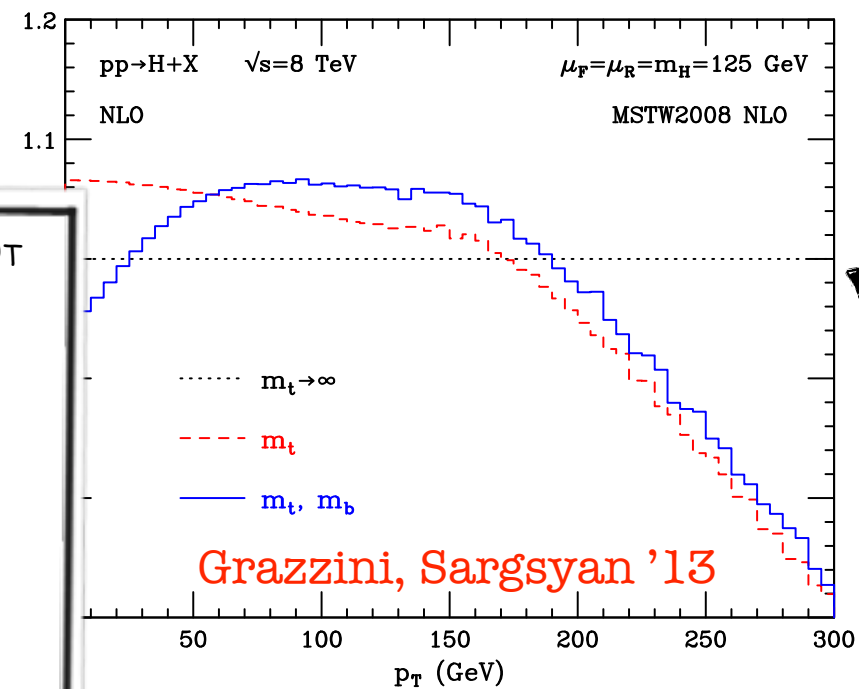
cut open the top loops

Don't think it is easy to produce a Higgs with high p_T

\sqrt{s} [TeV]	p_T^{\min} [GeV]	$\sigma_{p_T^{\min}}^{\text{SM}}$ [fb]	δ	ϵ	gg, qg [%]
14	100	2200	0.016	0.023	67, 31
	150	830	0.069	0.13	66, 32
	200	350	0.20	0.31	65, 34
	250	160	0.39	0.56	63, 36
	300	75	0.61	0.89	61, 38
	350	38	0.86	1.3	58, 41
	400	20	1.1	1.8	56, 43
	450	11	1.4	2.3	54, 45
	500	6.3	1.7	2.9	52, 47
	550	3.7	2.0	3.6	50, 49
	600	2.2	2.3	4.4	48, 51
	650	1.4	2.6	5.2	46, 53
	700	0.87	3.0	6.2	45, 54
	750	0.56	3.3	7.2	43, 56
	800	0.37	3.7	8.4	42, 57

+ 1000
reduction

Grojean, Salvioni, Schlaffer, Weiler '13



Note: LO only

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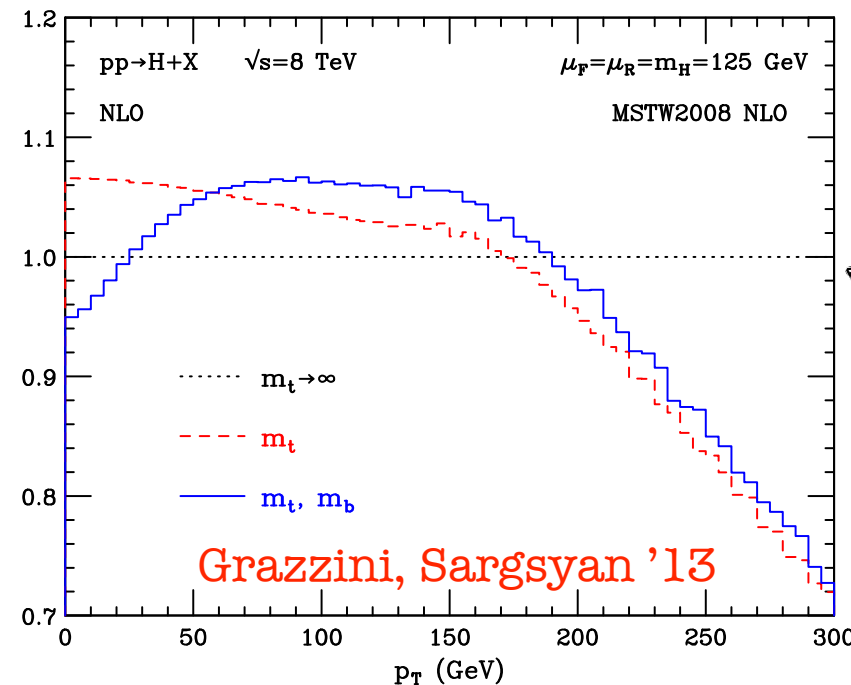
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we "see" the details of the particles
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Baur, Glover '90

Langenegger, Spira, Starodumov, Trueb '06



Note: LO only

NLO_{mt} is not known

1/ m_t corrections known $O(\alpha_s^4)$

few % up to $p_T \sim 150$ GeV

Harlander et al '12

the high p_T tail
is tens' % sensitive
to the mass of top

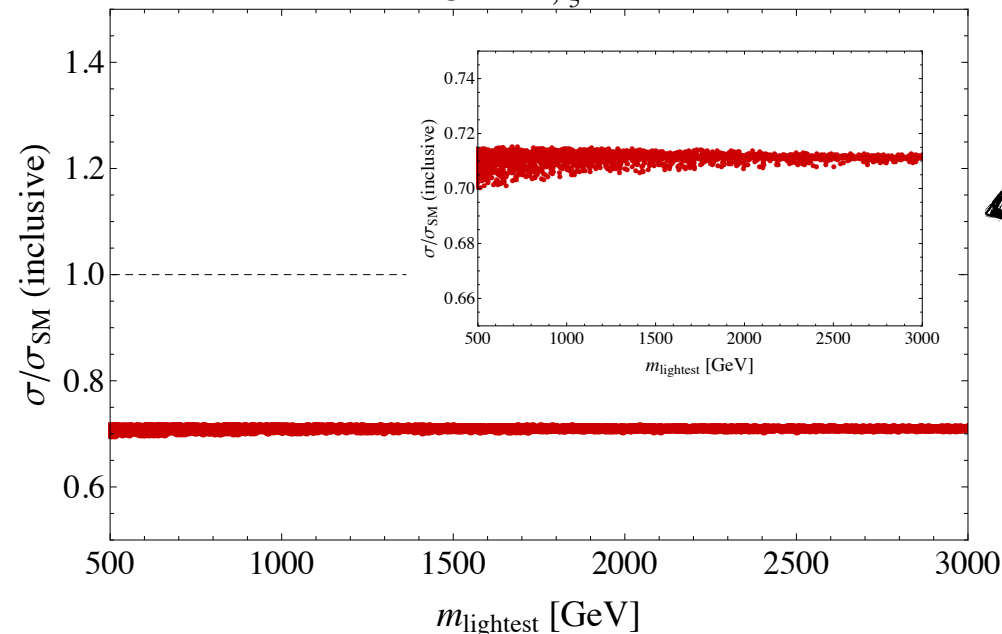
Composite Higgs Model top partners contributions

see also Banfi, Martin, Sanz '13

see also Azatov, Paul '13

Grojean, Salvioni, Schlaffer, Weiler '13

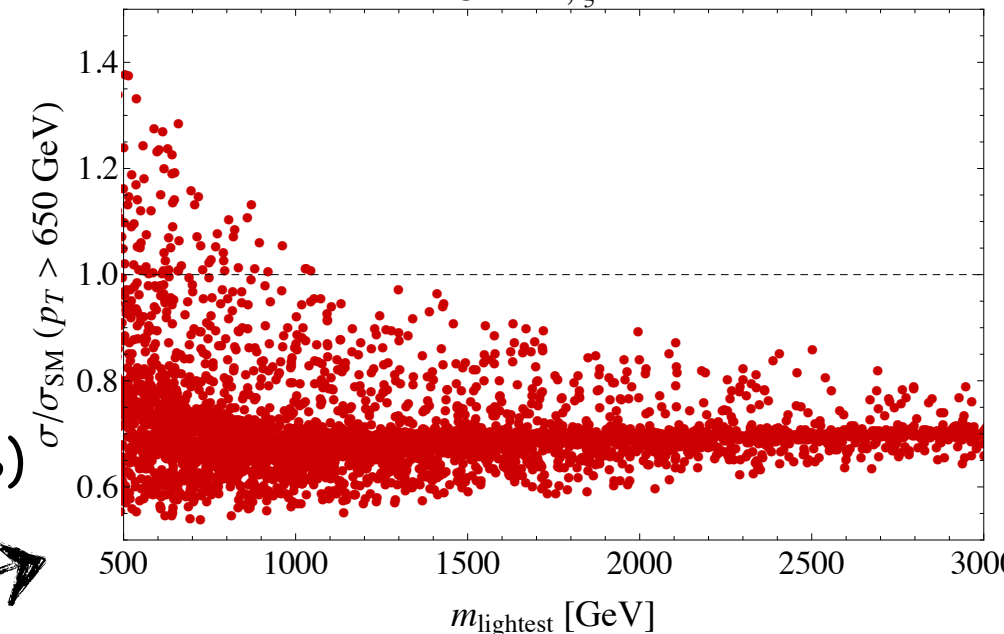
MCHM 5, $\xi = 0.1$



inclusive rate: $O(\%)$

with high- p_T cut: $O(\times 10\%)$

MCHM 5, $\xi = 0.1$



high- p_T tail "sees" the top partners that are missed by the inclusive rate

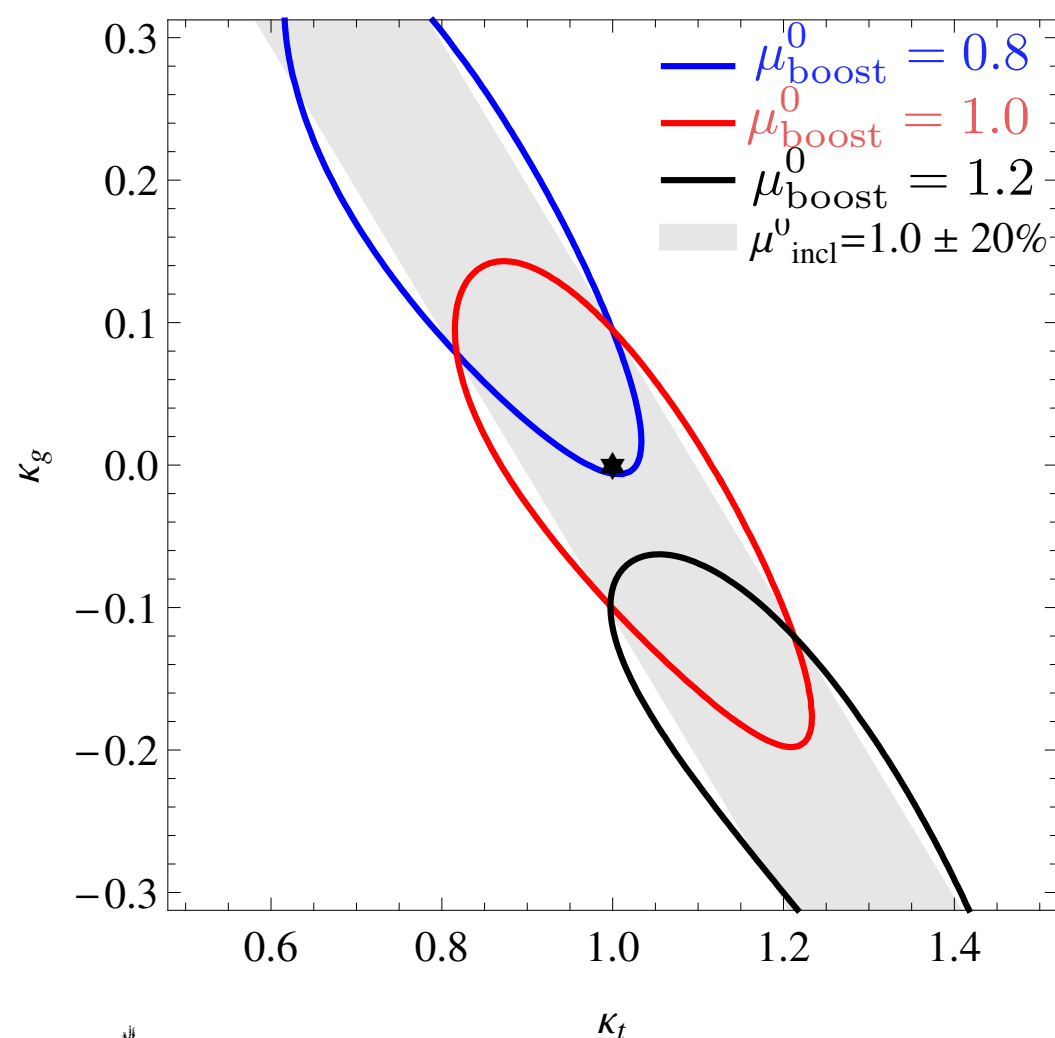
Boosted Higgs

high p_T tail discriminates short and long distance physics contribution to $gg \rightarrow h$

$$\sqrt{s} = 14 \text{ TeV}, \int dt \mathcal{L} = 3 \text{ ab}^{-1}, p_T > 650 \text{ GeV}$$

(partonic analysis in the boosted "ditau-jets" channel)

see Schlaffer et al '14 for a more complete analysis including WW channel



10-20% precision on κ_t



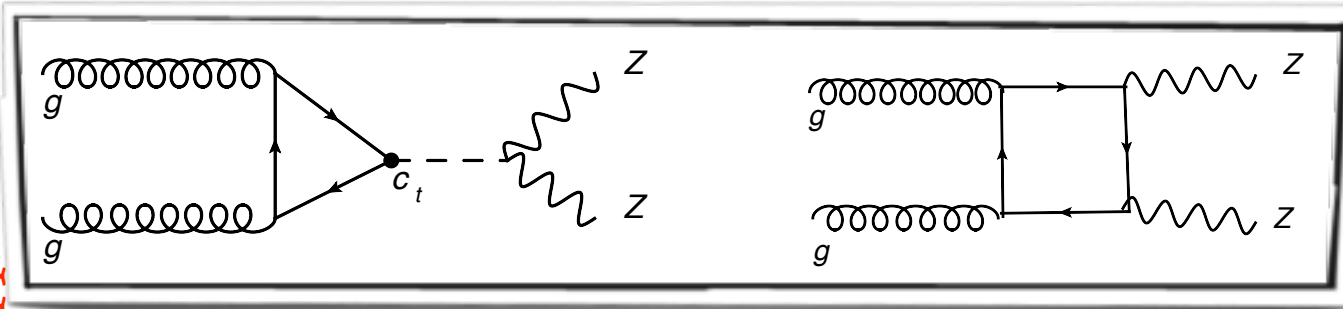
competitive/complementary to htt channel
for the measure the top-Higgs coupling

Are the NLO_m QCD corrections (not known) going to destroy all the sensitivity?
Frontier priority: N³LO_∞ for inclusive xs or NLO_{mt} for p_T spectrum?

Off-shell Higgs: $gg \rightarrow h^* \rightarrow ZZ \rightarrow 4l$

off-shell effects enhanced by the particular couplings of H to V_L

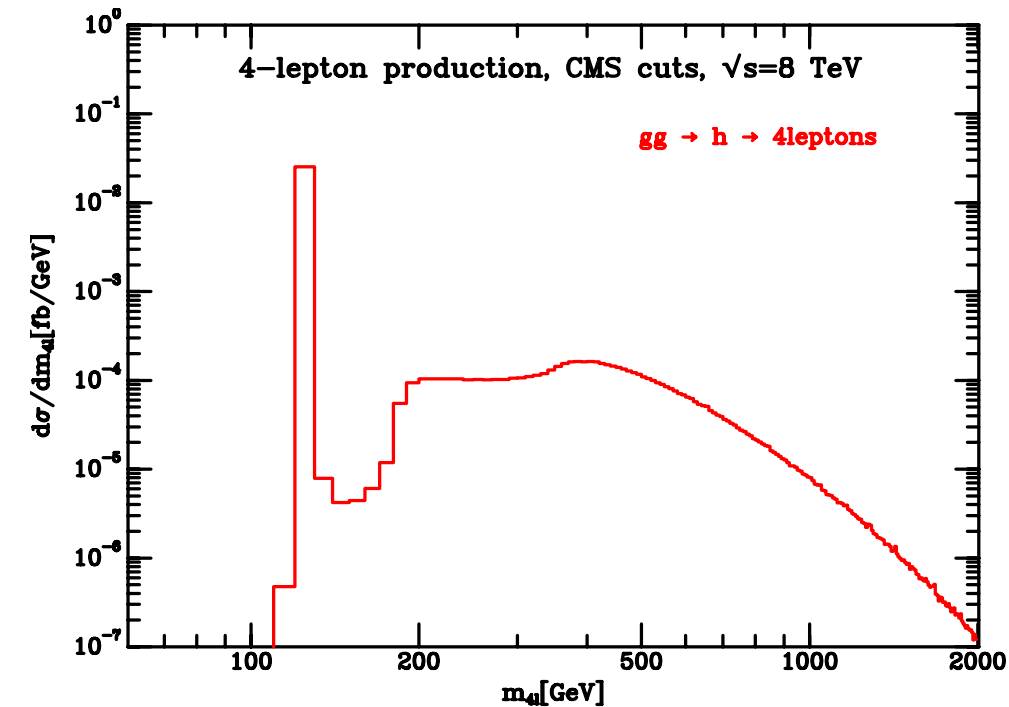
Glover, van der Bij '89



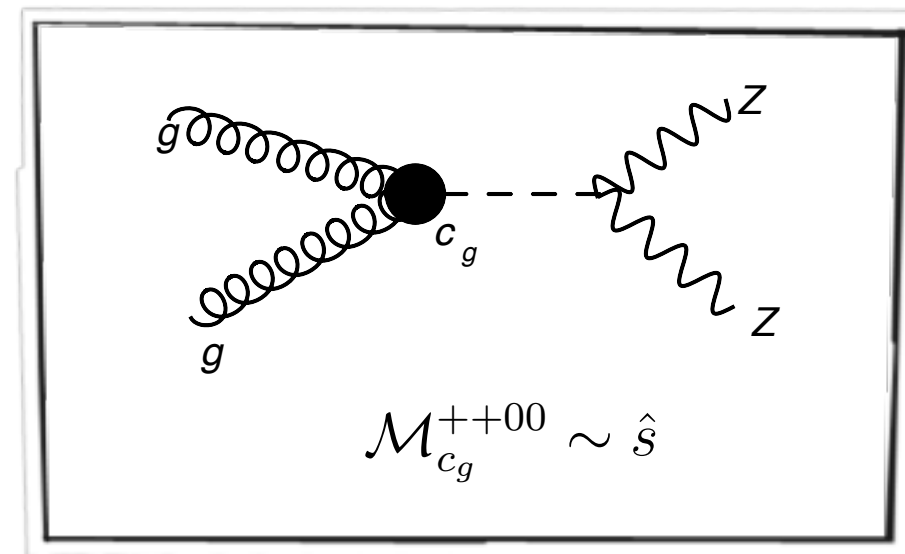
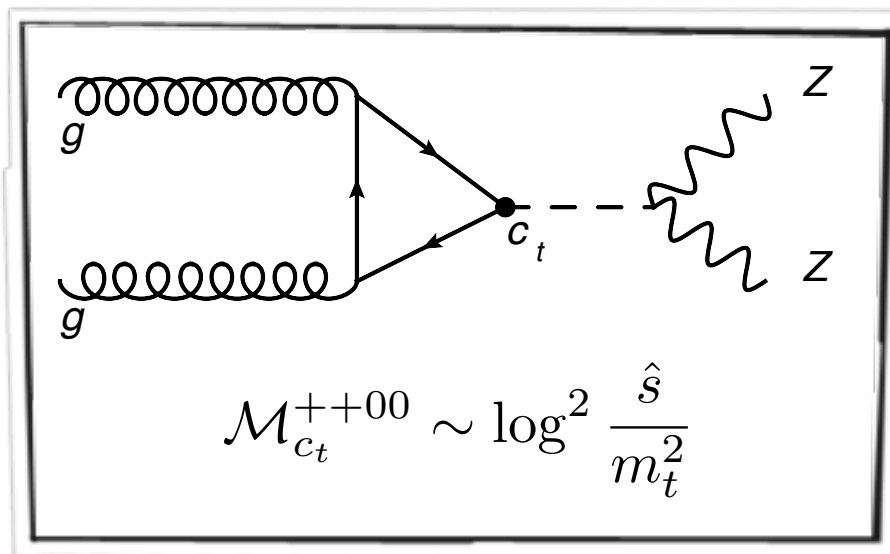
$$\mathcal{M}_{\text{Higgs}}^{++00} \sim \log^2 \frac{\hat{s}}{m_t^2} \quad \mathcal{M}_{\text{box}}^{++00} \sim -\log^2 \frac{\hat{s}}{m_t^2}$$

SM: cancelation forced by unitarity

BSM: deviations of Higgs couplings at large s will be amplified



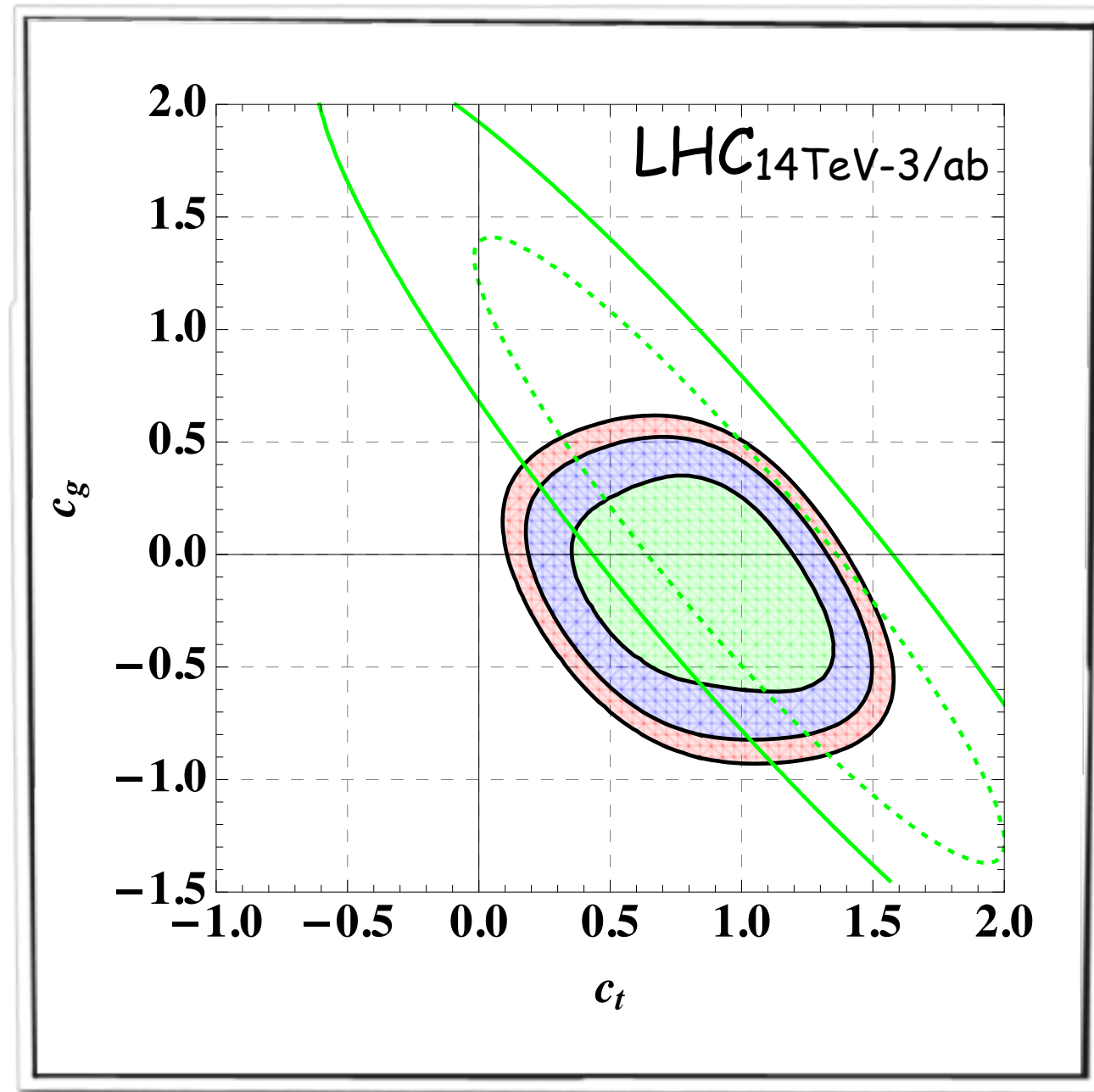
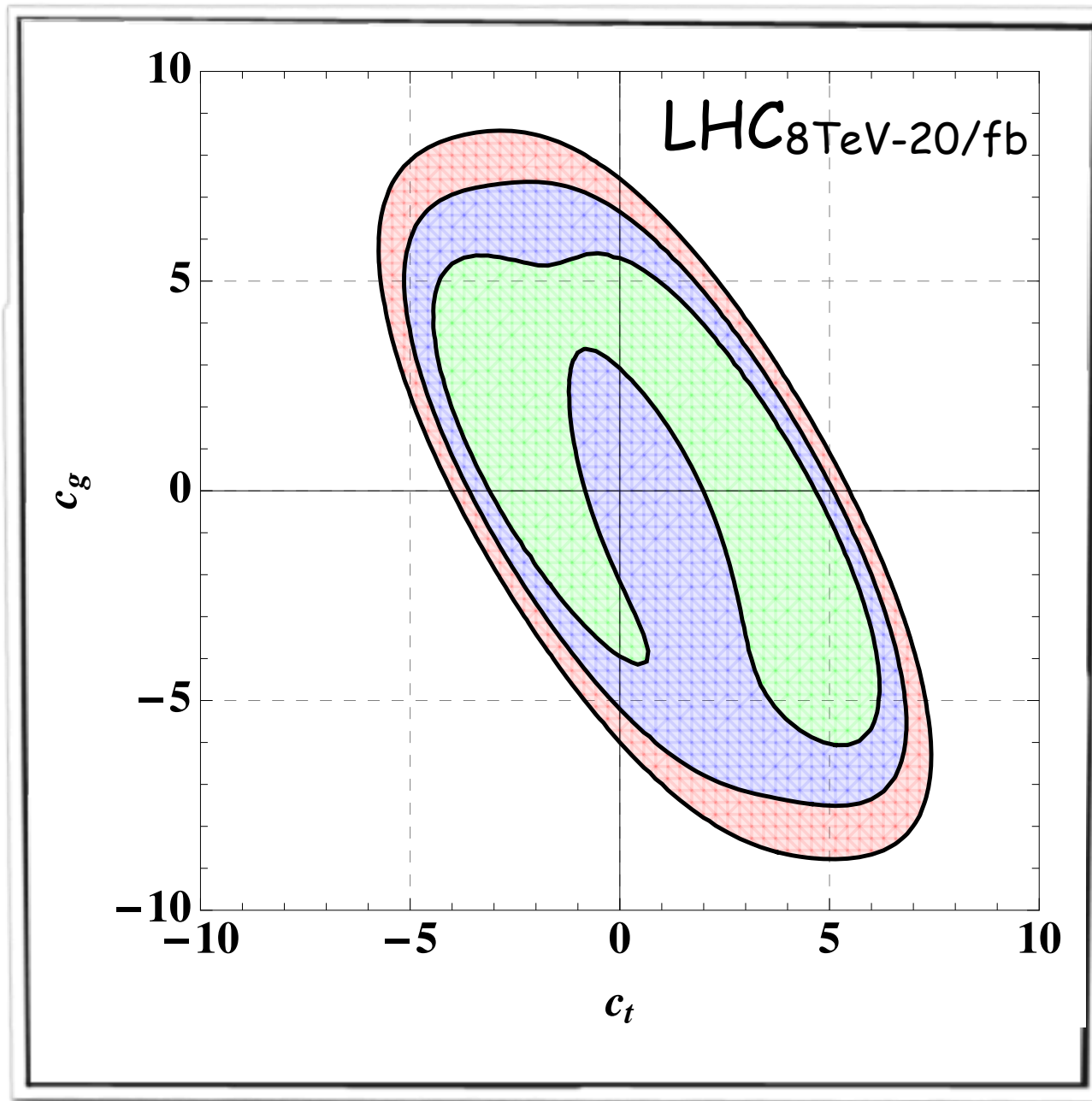
CMS interpretation in terms of bounds of the Higgs width is limited
data can be better used to measure the structure of the couplings at high \sqrt{s}

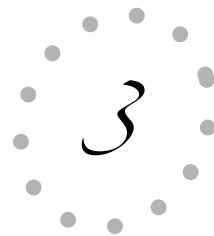


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Azatov, Grojean, Paul, Salvioni '14



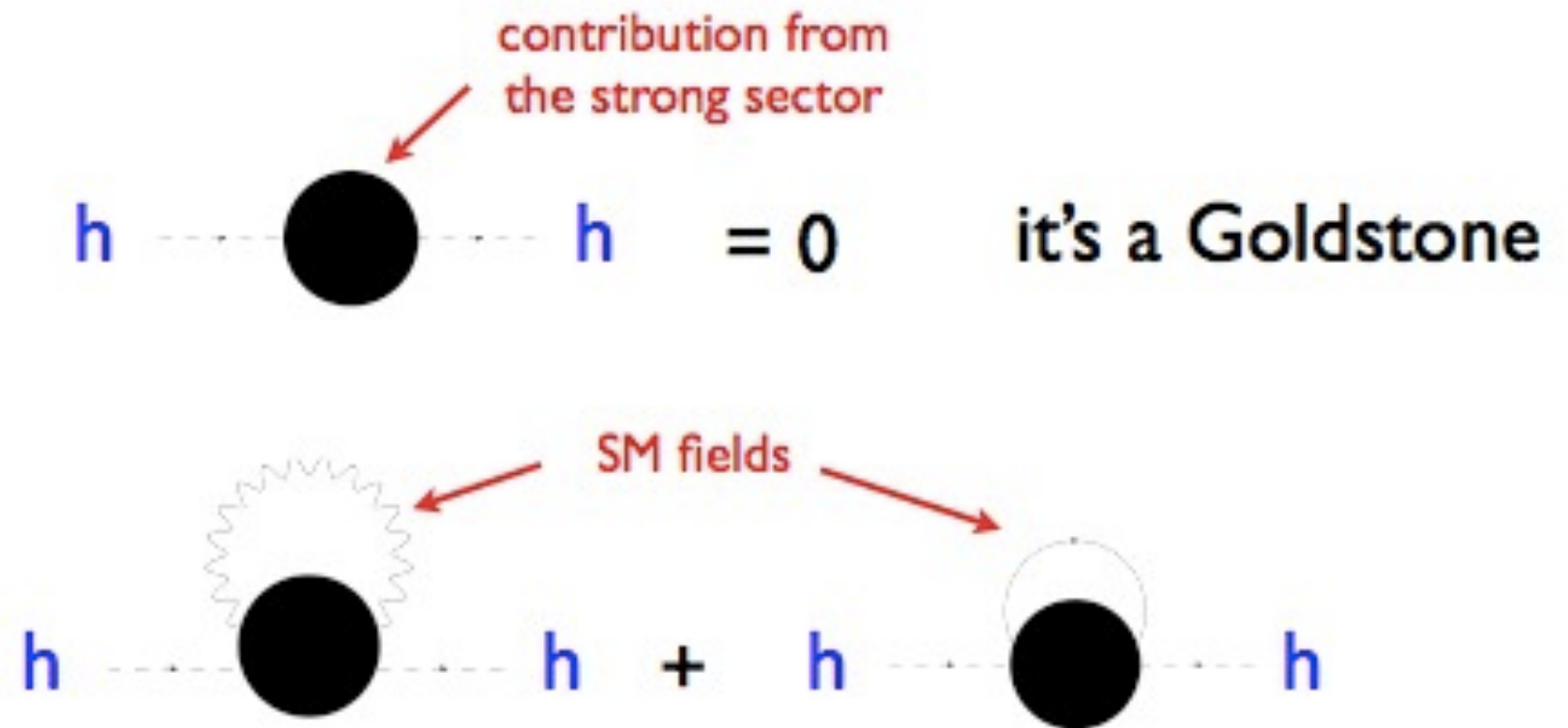


Direct searches of top partners

see also
De Curtis' talk

Light composite Higgs from "light" resonances

The interactions between the strong sector and the SM generate a potential for the Higgs



Impossible to compute the details of the potential from first principles
but using general properties on the asymptotic behavior of correlators
(saturation of Weinberg sum rules with the first few lightest resonances)

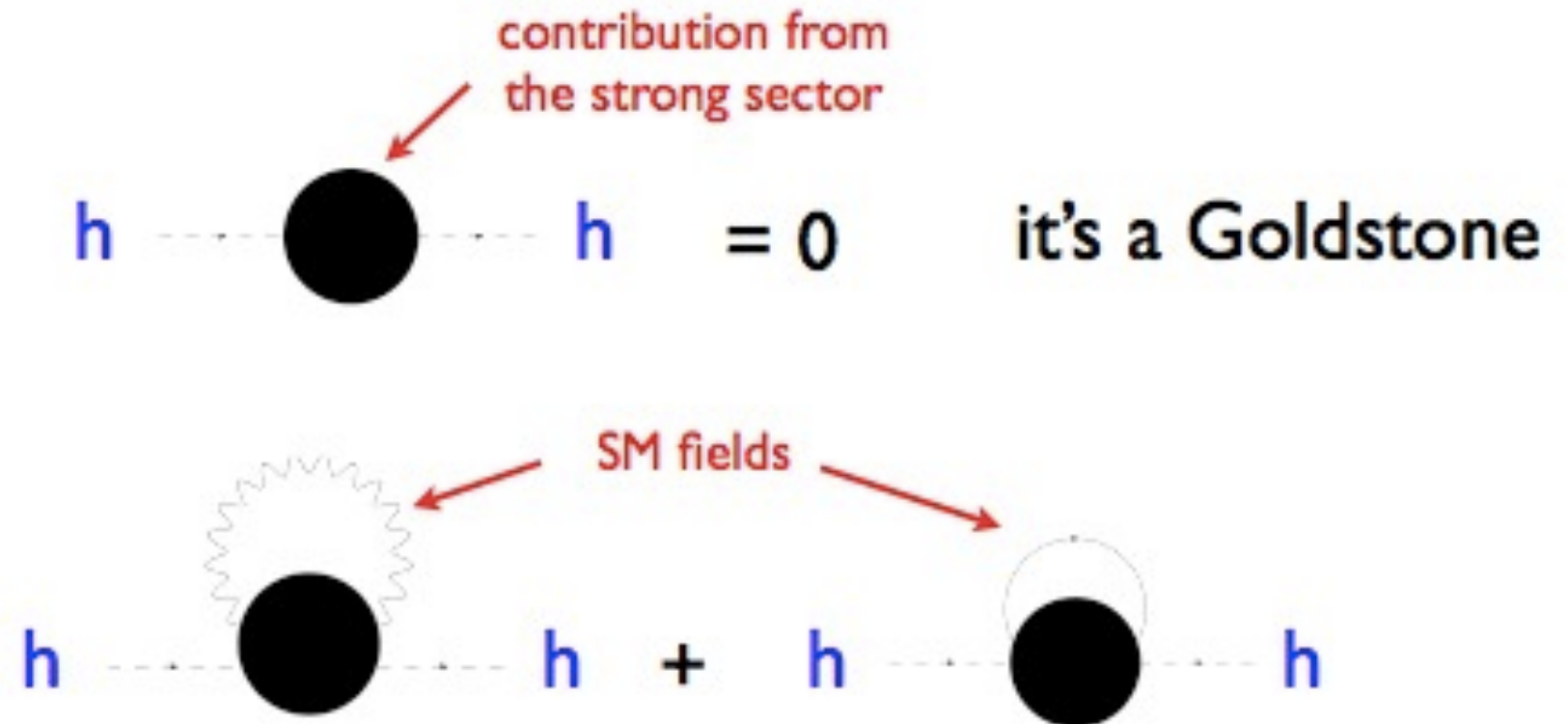
it is possible to estimate the Higgs mass

Pomarol, Riva '12

Marzocca, Serone, Shu '12

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Pomarol, Riva '12

$$m_h^2 \approx \frac{3}{\pi^2} \frac{m_t^2 m_Q^2}{f_{G/H}^2}$$



Marzocca, Serone, Shu '12

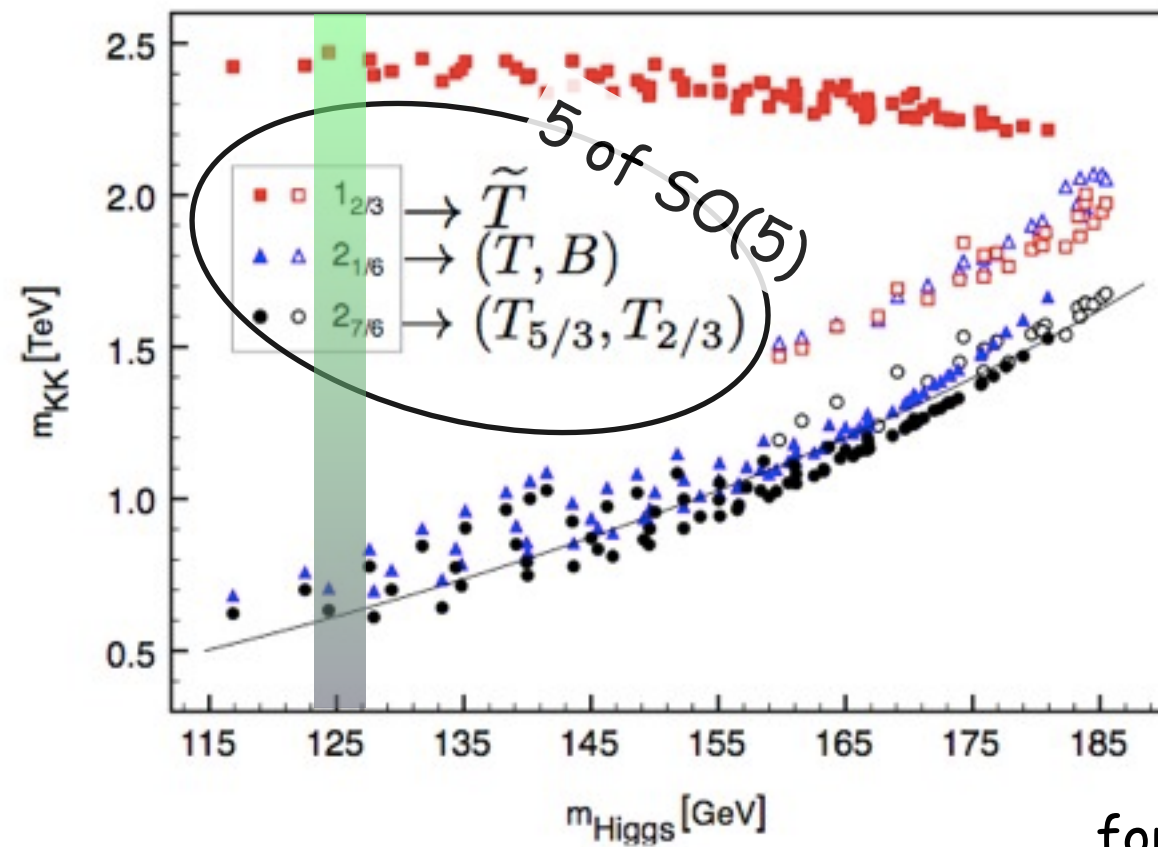
$$m_Q \lesssim 700 \text{ GeV} \left(\frac{m_h}{125 \text{ GeV}} \right) \left(\frac{160 \text{ GeV}}{m_t} \right) \left(\frac{f}{500 \text{ GeV}} \right)$$

fermionic resonances below $\sim 1 \text{ TeV}$
vector resonances $\sim \text{few TeV}$ (EW precision constraints)
 \sim for a natural ($< 20\%$ fine-tuning) set-up \sim

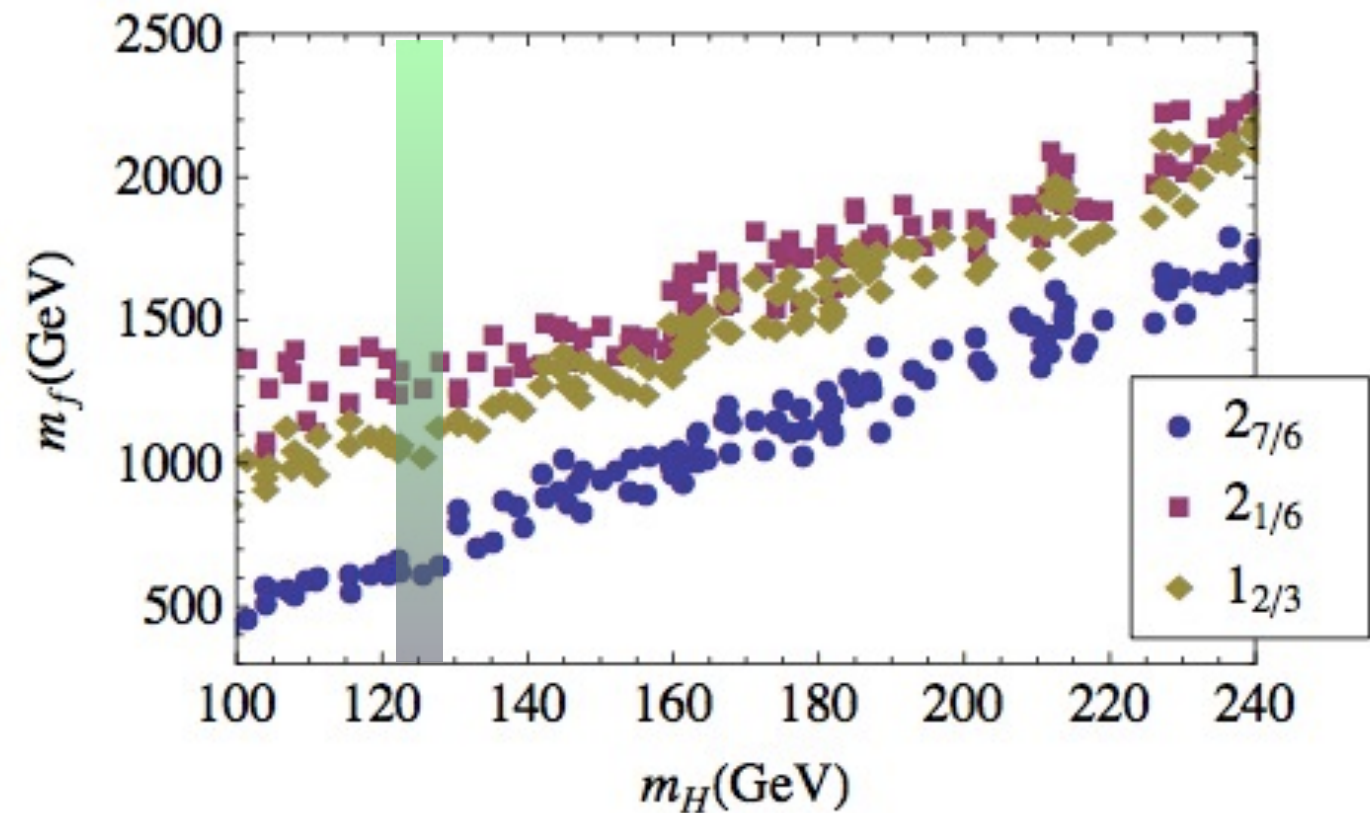
Light composite Higgs from "light" resonances

true spectrum in explicit realizations

Contino, Da Rold, Pomarol '06



De Curtis, Redi, Tesi '11



for similar results, see also

Matsedonskyi, Panico, Wulzer '12

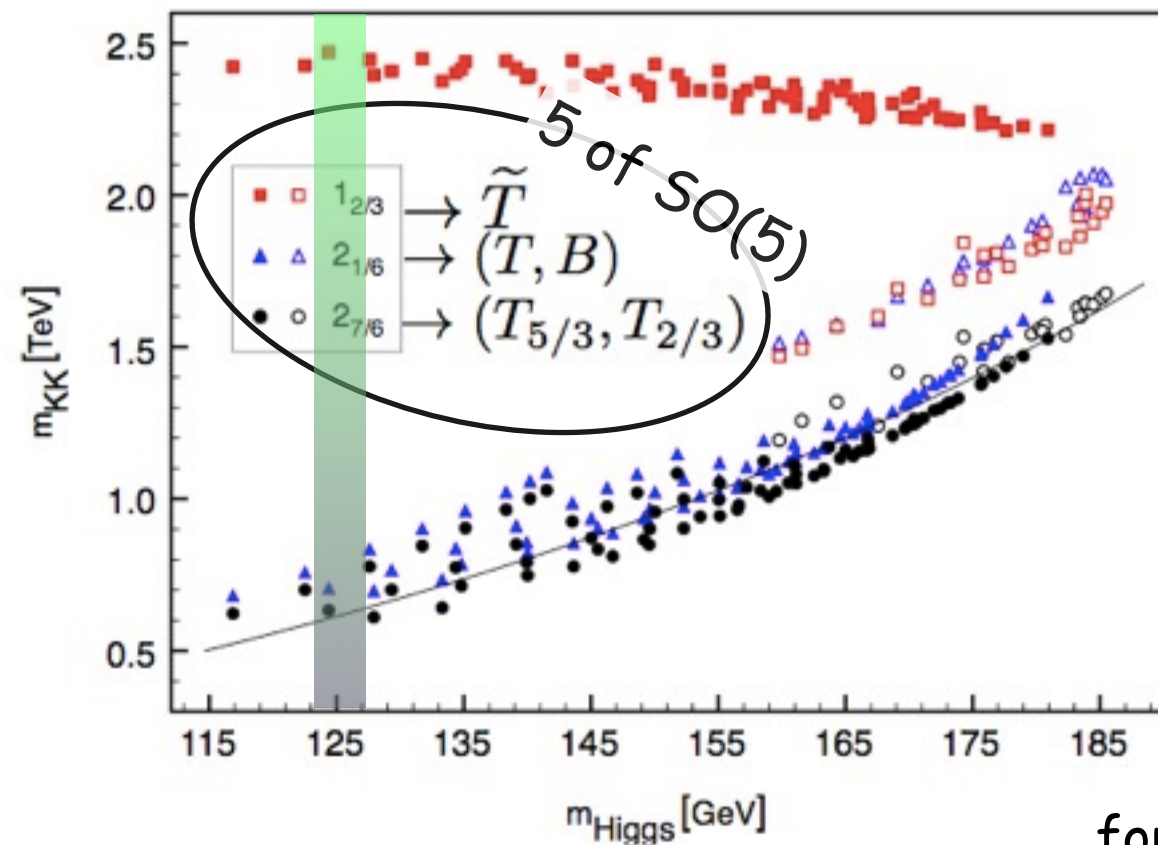
&

Marzocca, Serone, Shu '12

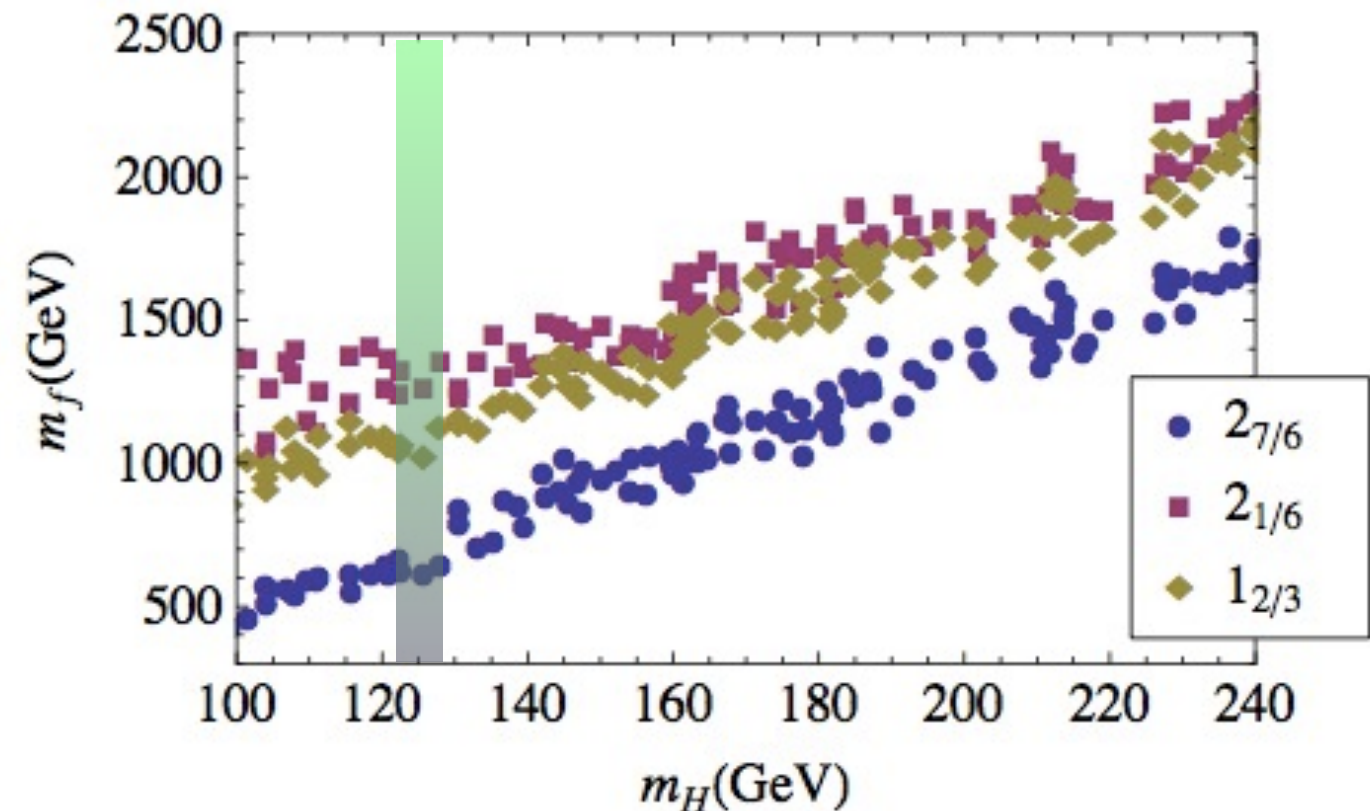
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&

Marzocca, Serone, Shu '12

Nice AdS/CFT interpretation

$$\text{Dim}[\mathcal{O}_\Psi] = \frac{3}{2} + |M_\Psi + \frac{1}{2}|$$

$M_\Psi = 1/2 \leftrightarrow \text{dim}[\mathcal{O}_\Psi] = 3/2 \leftrightarrow$ light free field decoupled from CFT

Rich phenomenology of the top partners

Search in same-sign di-lepton events

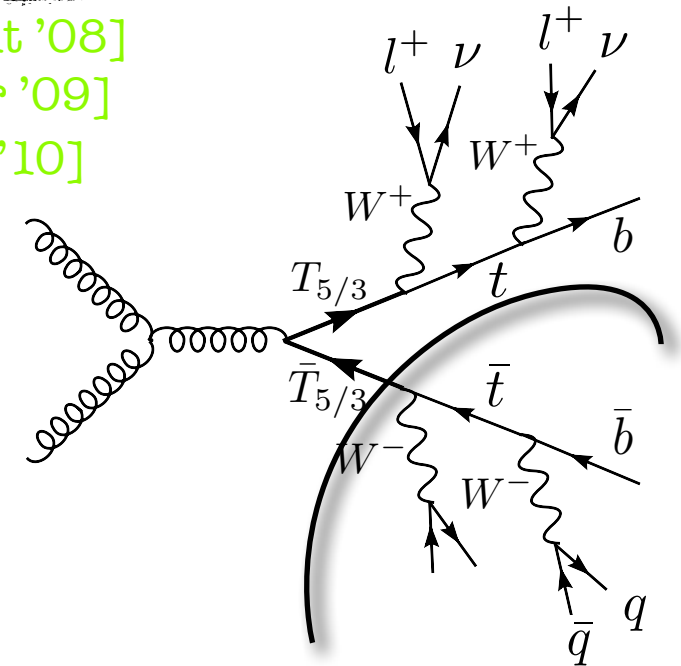
- $t\bar{t} + jets$ is not a background [except for charge mis-ID and fake e^-]
- the resonant (tW) invariant mass can be reconstructed

discovery potential (LHC_{14TeV})

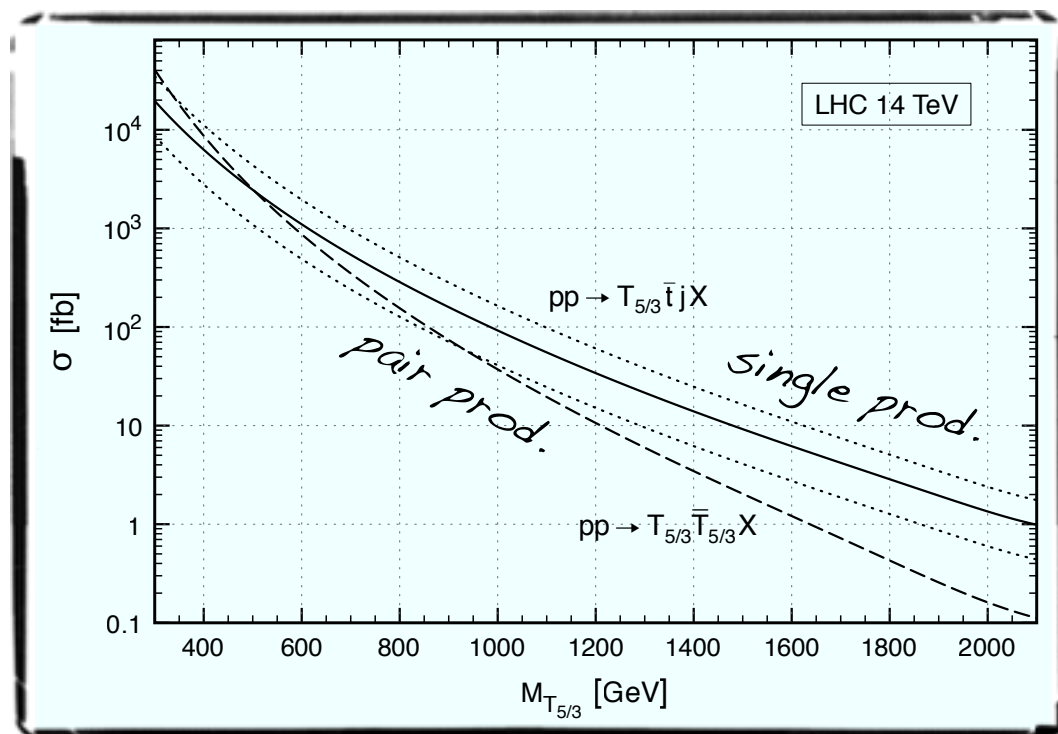
$M_{5/3} = 500 \text{ GeV}$ ($\sigma \times BR \approx 100/\text{fb}$) $\rightarrow 56 \text{ pb}^{-1}$

$M_{5/3} = 1 \text{ TeV}$ ($\sigma \times BR \approx 2/\text{fb}$) $\rightarrow 15 \text{ fb}^{-1}$

[Contino, Servant '08]
[Mrazek, Wulzer '09]
[Dissertori et al '10]

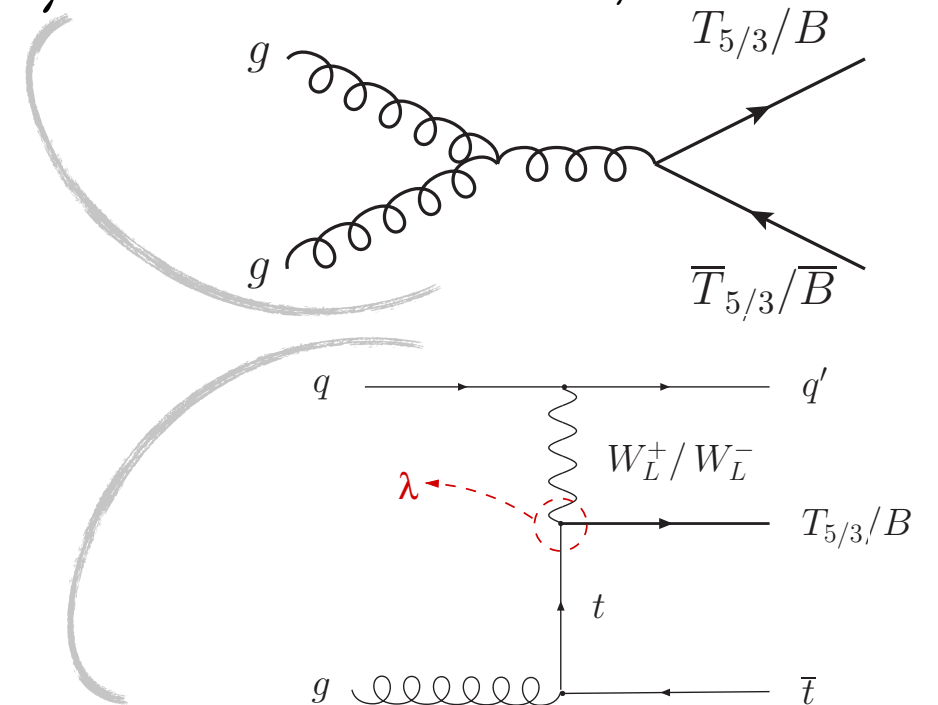


Dissertori, Furlan, Moortgat, Nef '09
4



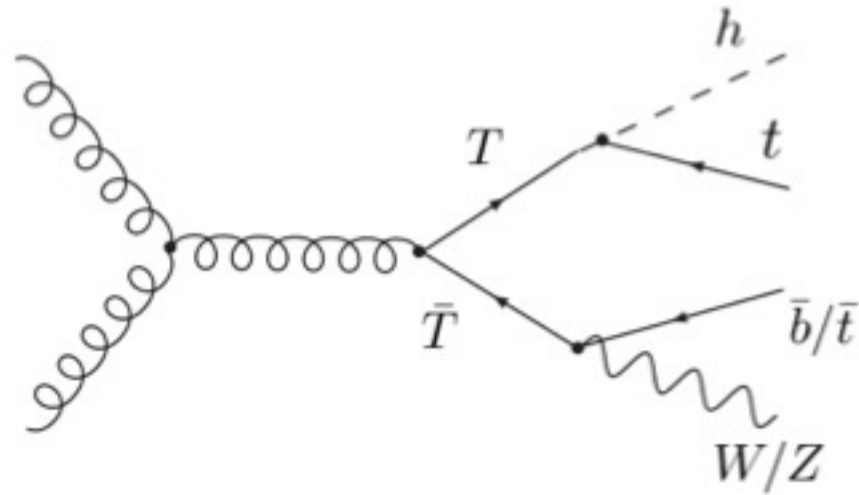
[Contino, Servant '08]

Pair production (model independent)



Single production (model dependent)

Rich phenomenology of the top partners



$l^\pm + 4b$ final state

Aguilar-Saavedra '09

$$T\bar{T} \rightarrow HtW^- \bar{b} \rightarrow HW^+ b W^- \bar{b}$$

$$H \rightarrow b\bar{b}, WW \rightarrow \ell\nu q\bar{q}',$$

$$T\bar{T} \rightarrow HtV\bar{t} \rightarrow HW^+ b VW^- \bar{b}$$

$$H \rightarrow b\bar{b}, WW \rightarrow \ell\nu q\bar{q}', V \rightarrow q\bar{q}/\nu\bar{\nu}$$

$l^\pm + 6b$ final state

Aguilar-Saavedra '09

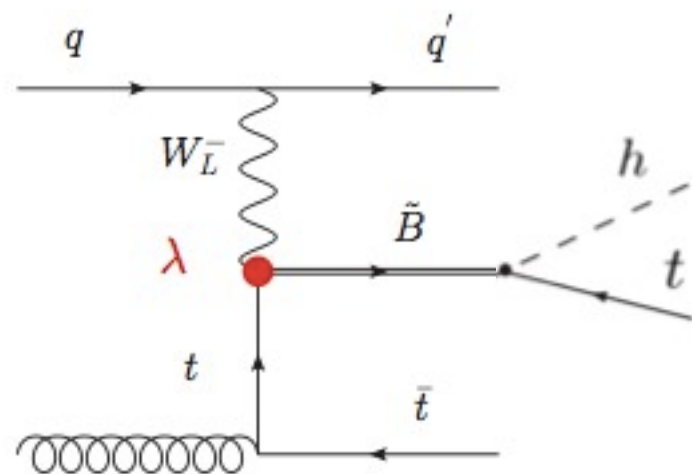
$$T\bar{T} \rightarrow HtH\bar{t} \rightarrow HW^+ b HW^- \bar{b}$$

$$H \rightarrow b\bar{b}, WW \rightarrow \ell\nu q\bar{q}'$$

$\gamma\gamma$ final state

Azatov et al '12

$$thbW/thtZ/thth, h \rightarrow \gamma\gamma$$



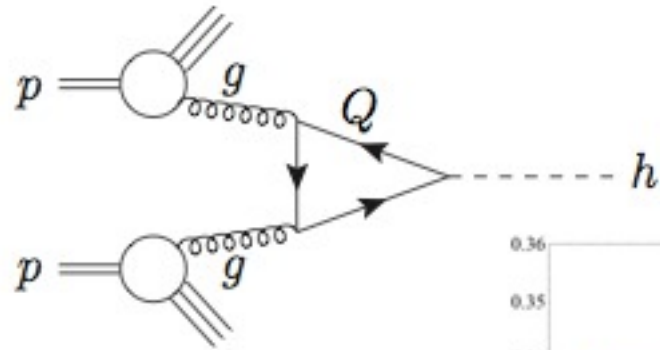
$l^\pm + 4b$ final state

Vignaroli '12

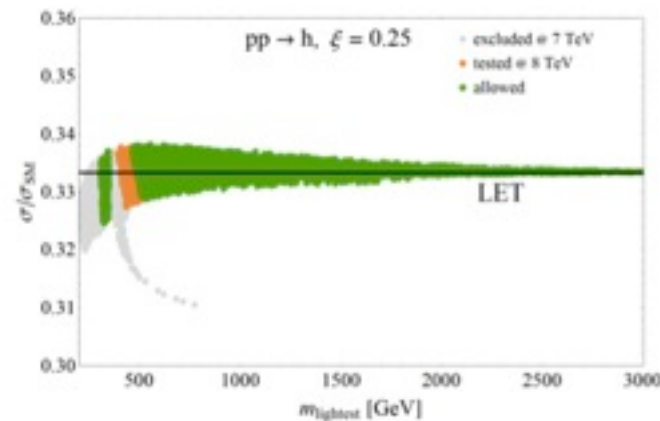
$$pp \rightarrow (\tilde{B} \rightarrow (h \rightarrow b\bar{b})b)t + X$$

Top partners & Higgs physics

~ current single higgs processes are insensitive to top partners ~



$$\sigma_{14\text{TeV}}^{\text{SM}} \approx 50 \text{ pb}$$



two competing effects that cancel:

- ☑ T's run in the loops
- ☑ T's modify top Yukawa coupling

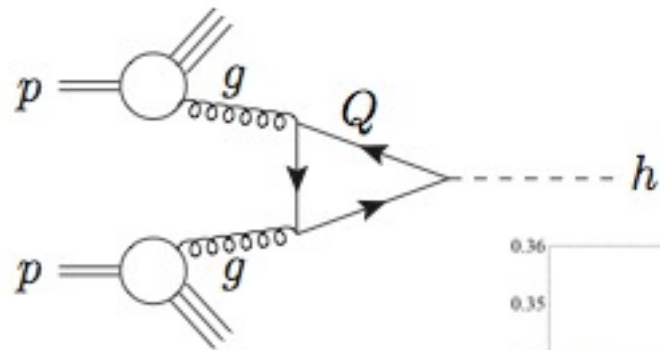
Falkowski '07

Azatov, Galloway '11

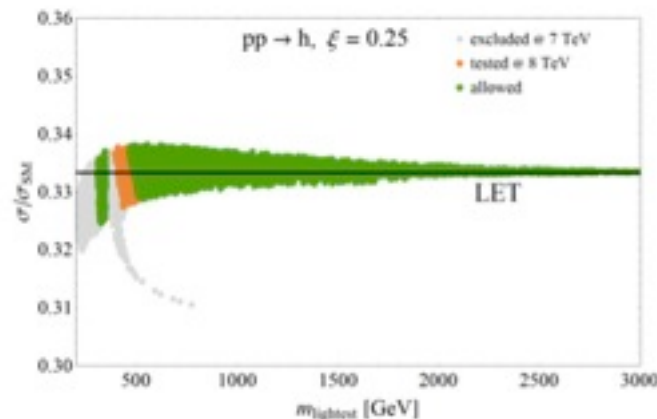
Delaunay, Grojean, Perez, '13

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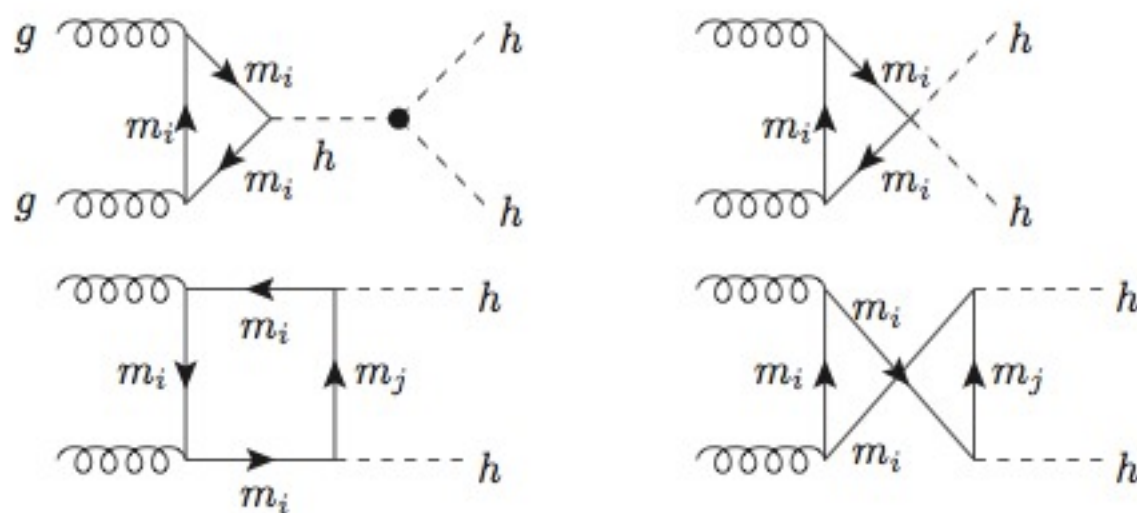
Falkowski '07

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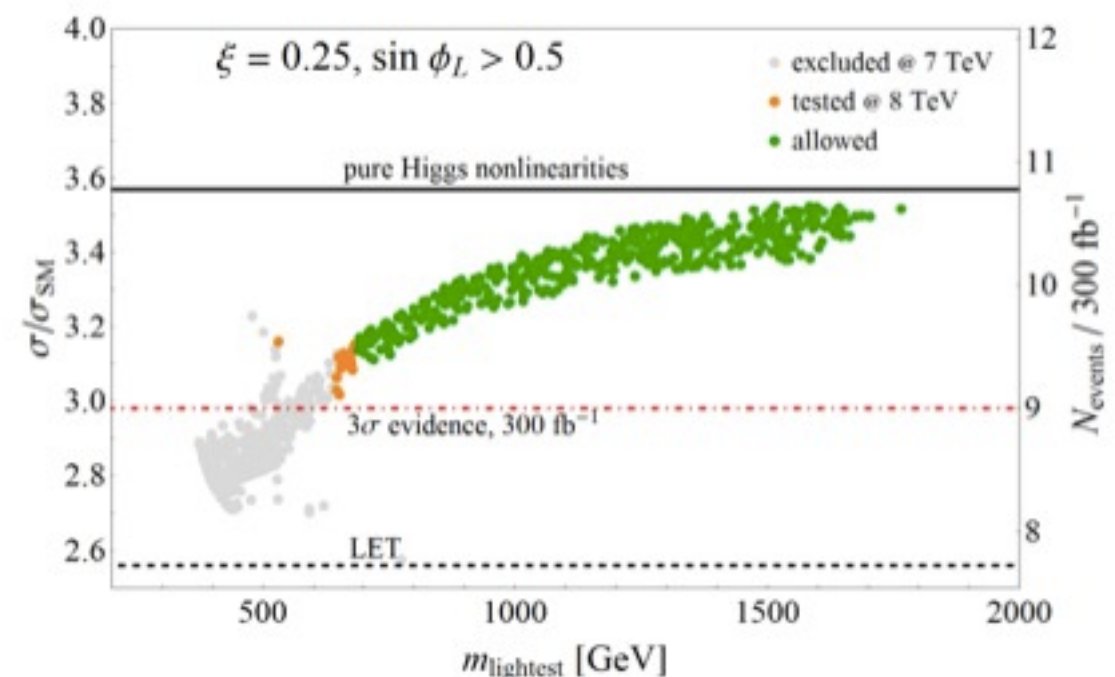
Delaunay, Grojean, Perez, '13

~ sensitivity in double Higgs production ~

Gillioz, Grober, Grojean, Muhlleitner, Salvioni '12



$$\sigma_{14\text{TeV}}^{\text{SM}} = 17.9\text{fb}$$



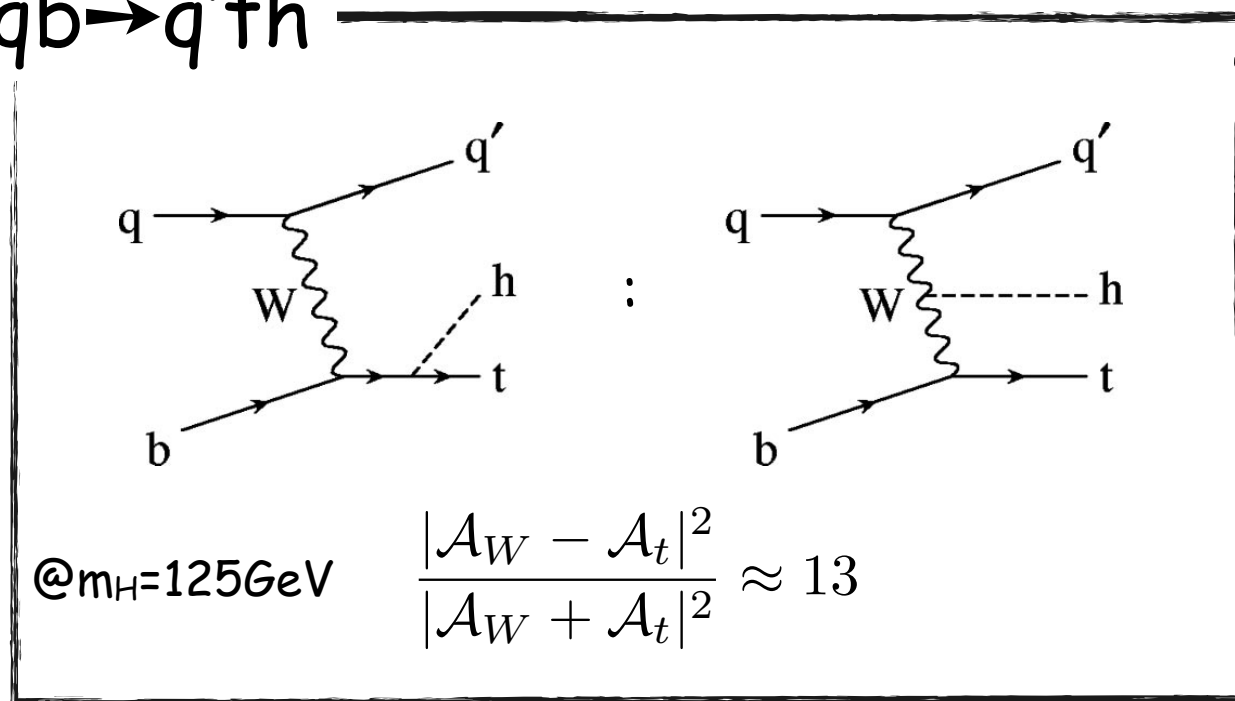
Top partners & Higgs physics

direct measurement of top-higgs coupling

htt is important but challenging channel

may be easier channel to look at

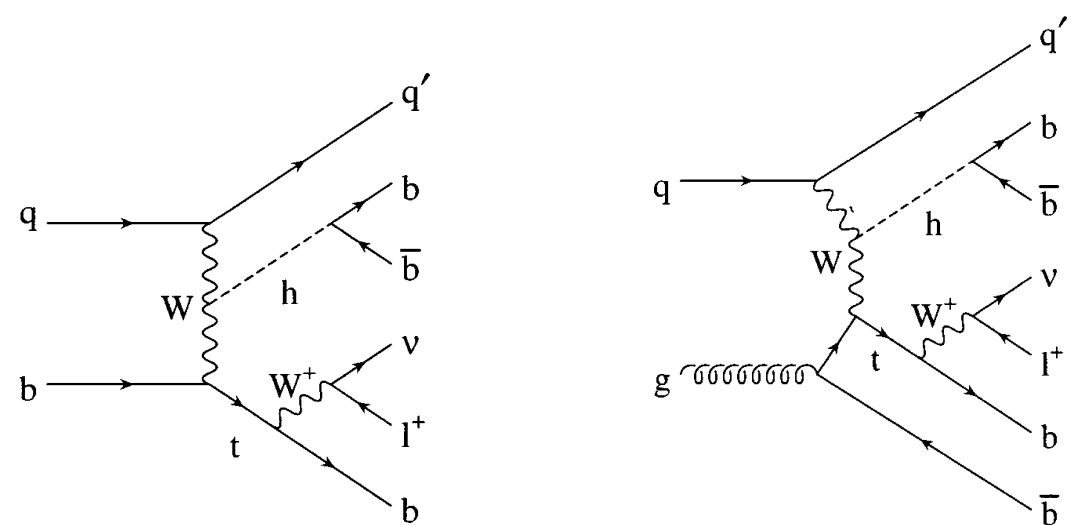
$qb \rightarrow q'th$



Farina, Grojean, Maltoni, Salvioni, Thamm '12

look at final states:

$3b + 1 \text{ fwd jet} + l^\pm + p^T$. $4b + 1 \text{ fwd jet} + l^\pm + p^T$.

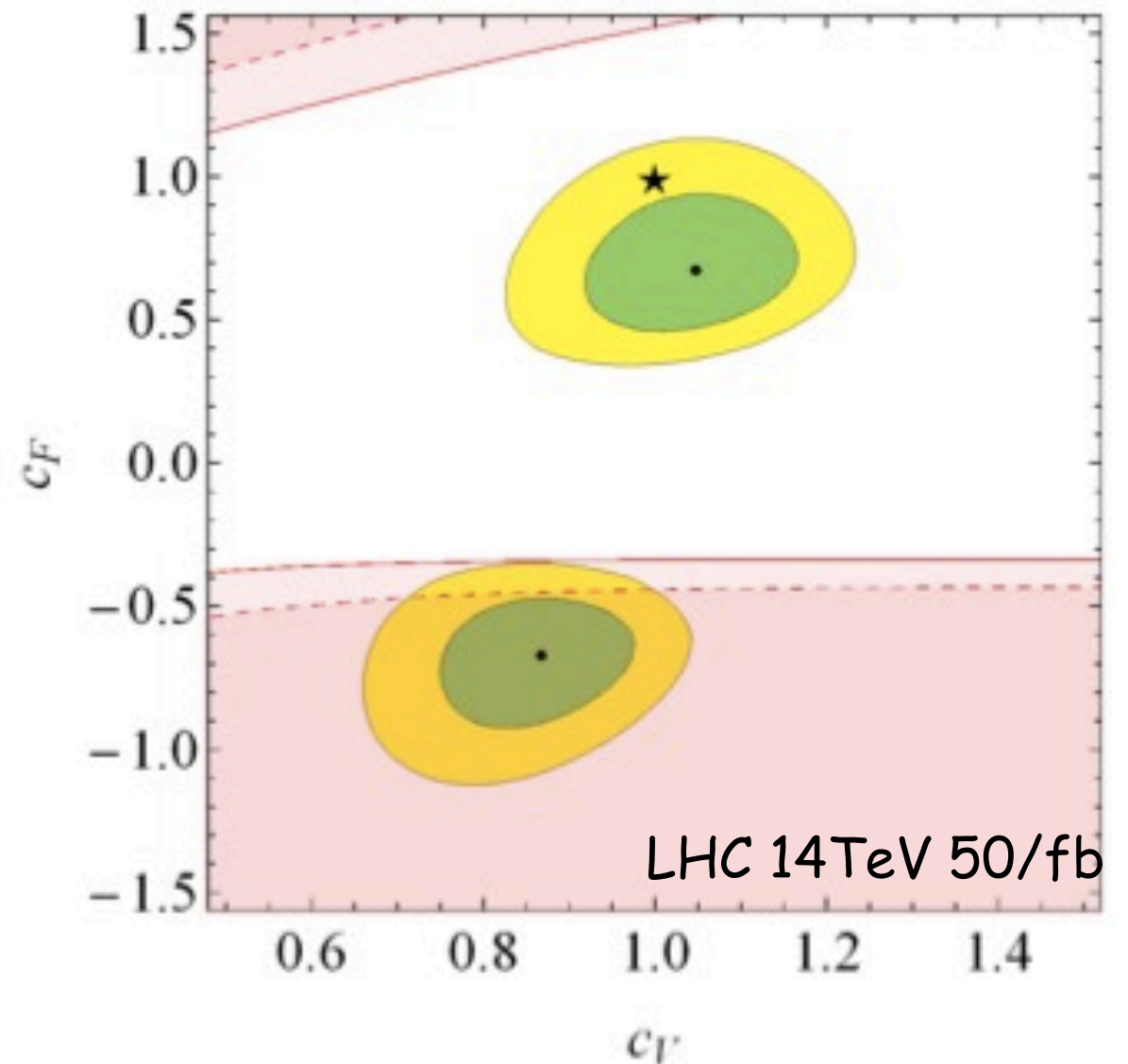
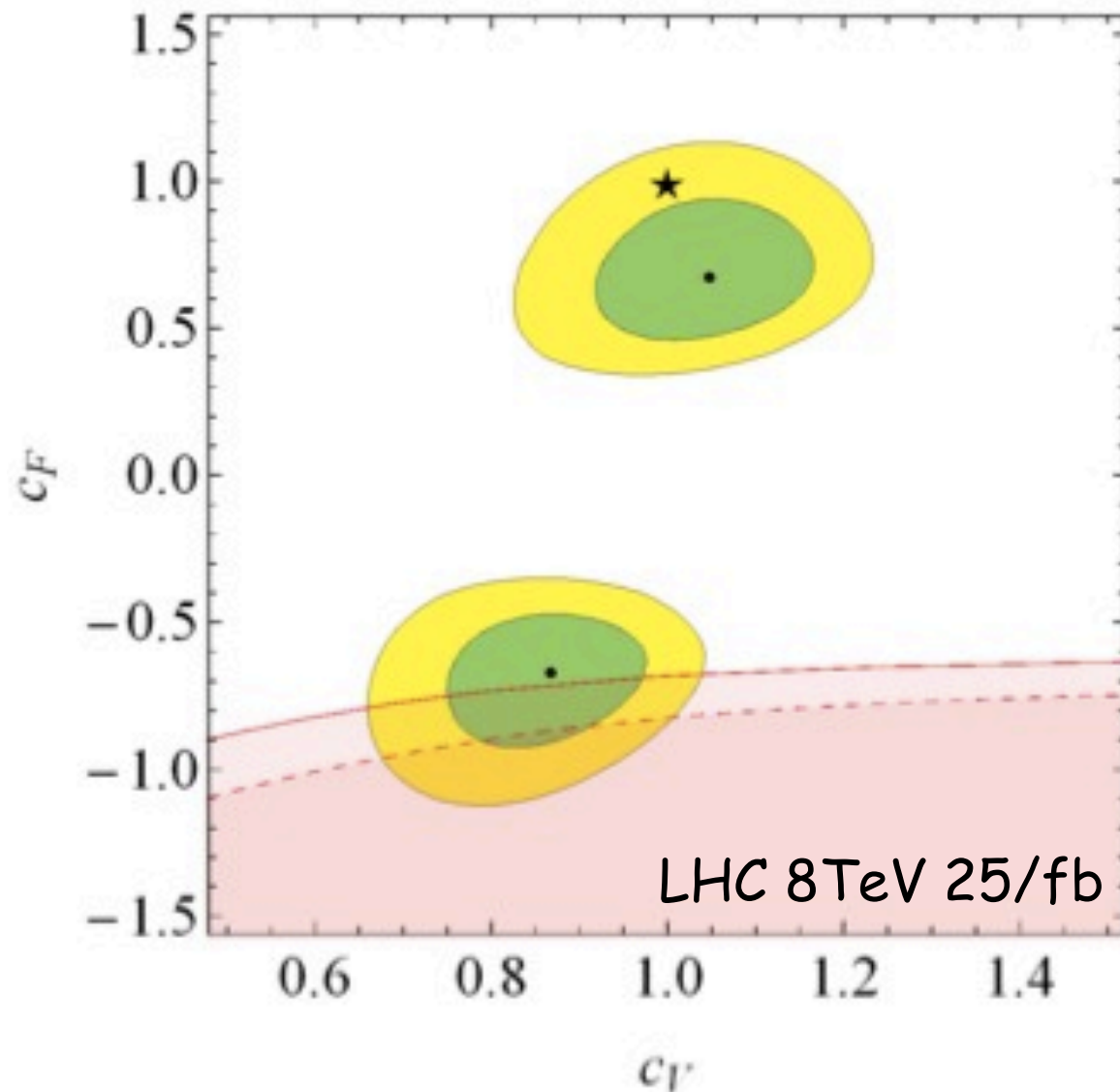


	$\sigma(pp \rightarrow tjh)$ [fb]		$\sigma(pp \rightarrow tjh\bar{b})$ [fb]	
	$c_F = 1$	$c_F = -1$	$c_F = 1$	$c_F = -1$
8 TeV	17.3	252.7	12.14	181.4
14 TeV	80.6	1042	59.6	828.5

Top partners & Higgs physics

direct measurement of top-higgs coupling

single-top in association with Higgs



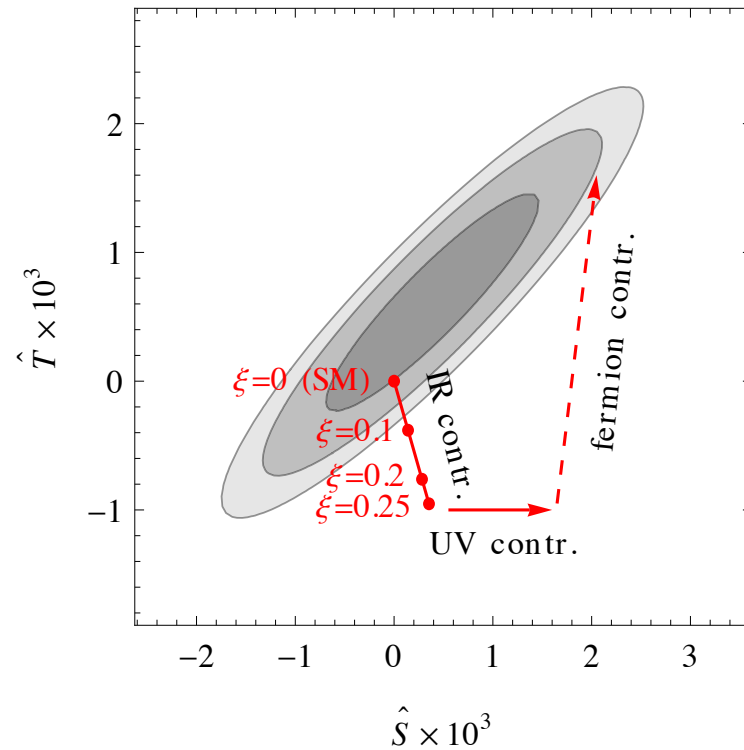
68% and 95% CL exclusion region vs current Higgs coupling fit

Farina, Grojean, Maltoni, Salvioni, Thamm '12

Top partners & EWPT

Grojean, Matsedonskyi, Panico '13

Oblique parameters



tree-level contribution

$$\Delta \hat{S} \simeq \frac{g^2}{g_*^2} \xi \simeq \frac{m_w^2}{m_*^2}$$

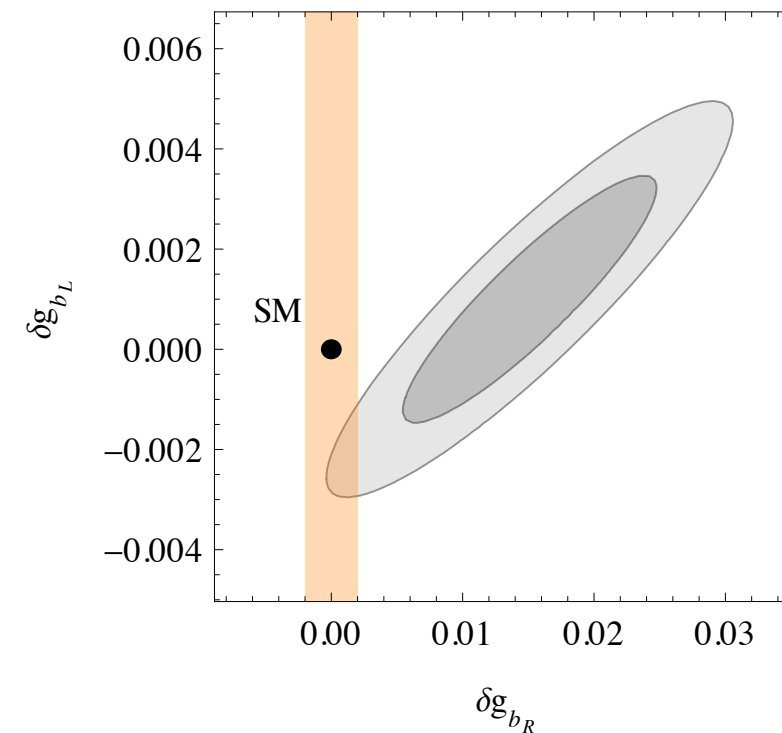
Higgs loop

$$\Delta \hat{S} = \frac{g^2}{192\pi^2} \xi \log \left(\frac{m_*^2}{m_h^2} \right) \simeq 1.4 \cdot 10^{-3} \xi \quad \Delta \hat{T} = -\frac{3g'^2}{64\pi^2} \xi \log \left(\frac{m_*^2}{m_h^2} \right) \simeq -3.8 \cdot 10^{-3} \xi$$

fermion loop

$$\Delta \hat{S}_{ferm}^{div} = \frac{g^2}{8\pi^2} (1 - 2c^2) \xi \log \left(\frac{m_*^2}{m_4^2} \right) \quad \Delta \hat{T} \simeq \frac{N_c}{16\pi^2} y_t^2 \xi \simeq 2 \cdot 10^{-2} \xi$$

Zb_Lb_L



tree-level contribution

$$\frac{\delta g_{b_L}}{g_{b_L}^{SM}} \sim \frac{y_L^2 f^2}{m^2} \frac{m_z^2}{m_*^2} \simeq 8 \cdot 10^{-4} \frac{f}{m} \left(\frac{4\pi}{g_*} \right)^2 \xi$$

fermion loop

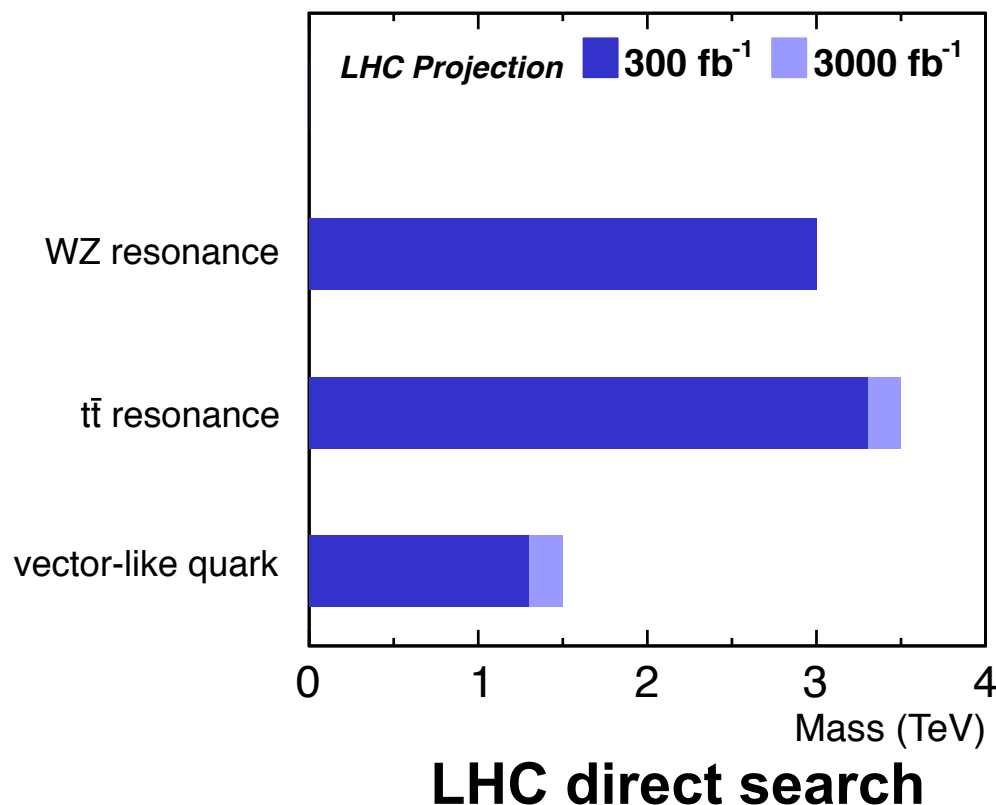
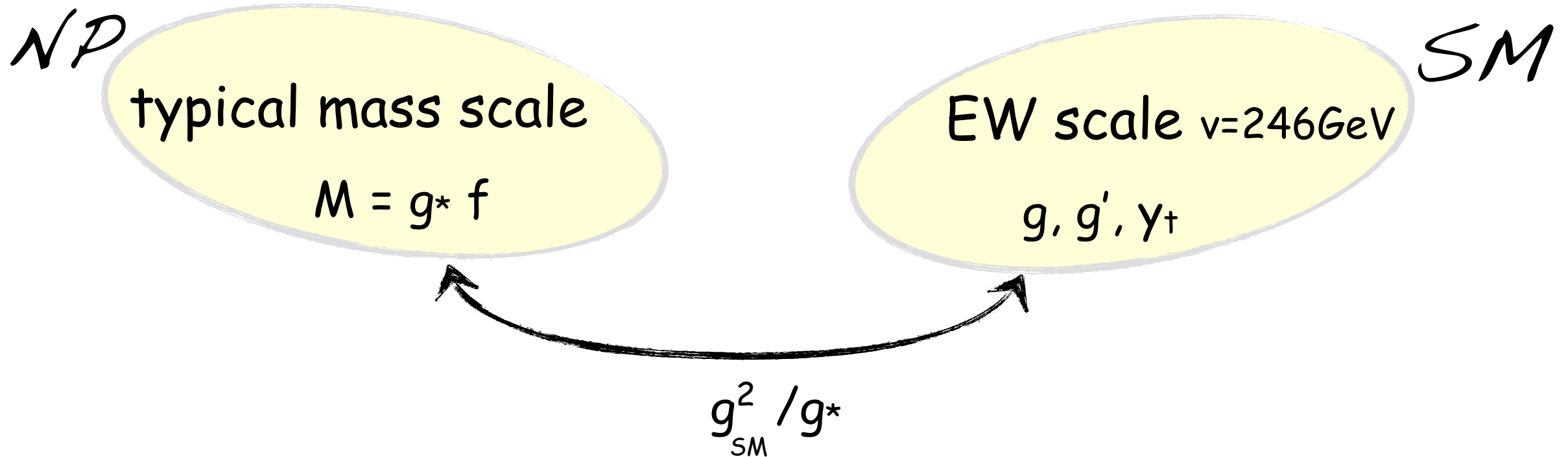
$$\frac{\delta g_{b_L}}{g_{b_L}^{SM}} \simeq \frac{y_t^2}{16\pi^2} \xi \log \left(\frac{m_*^2}{m_4^2} \right) \simeq 2 \cdot 10^{-2} \xi$$

$\xi < 0.1 \Rightarrow$ we might have to wait LHC-HL to see any new physics in Higgs data
BSM Higgs precision era

Conclusions: Higgs & New Physics

Precision /indirect searches (high lumi.) vs. direct searches (high energy)

Contino, Grojean, Pappadopulo, Rattazzi, Thamm '13



○ Precision Higgs study: $\xi \equiv \frac{\delta g}{g} = \frac{v^2}{f^2}$

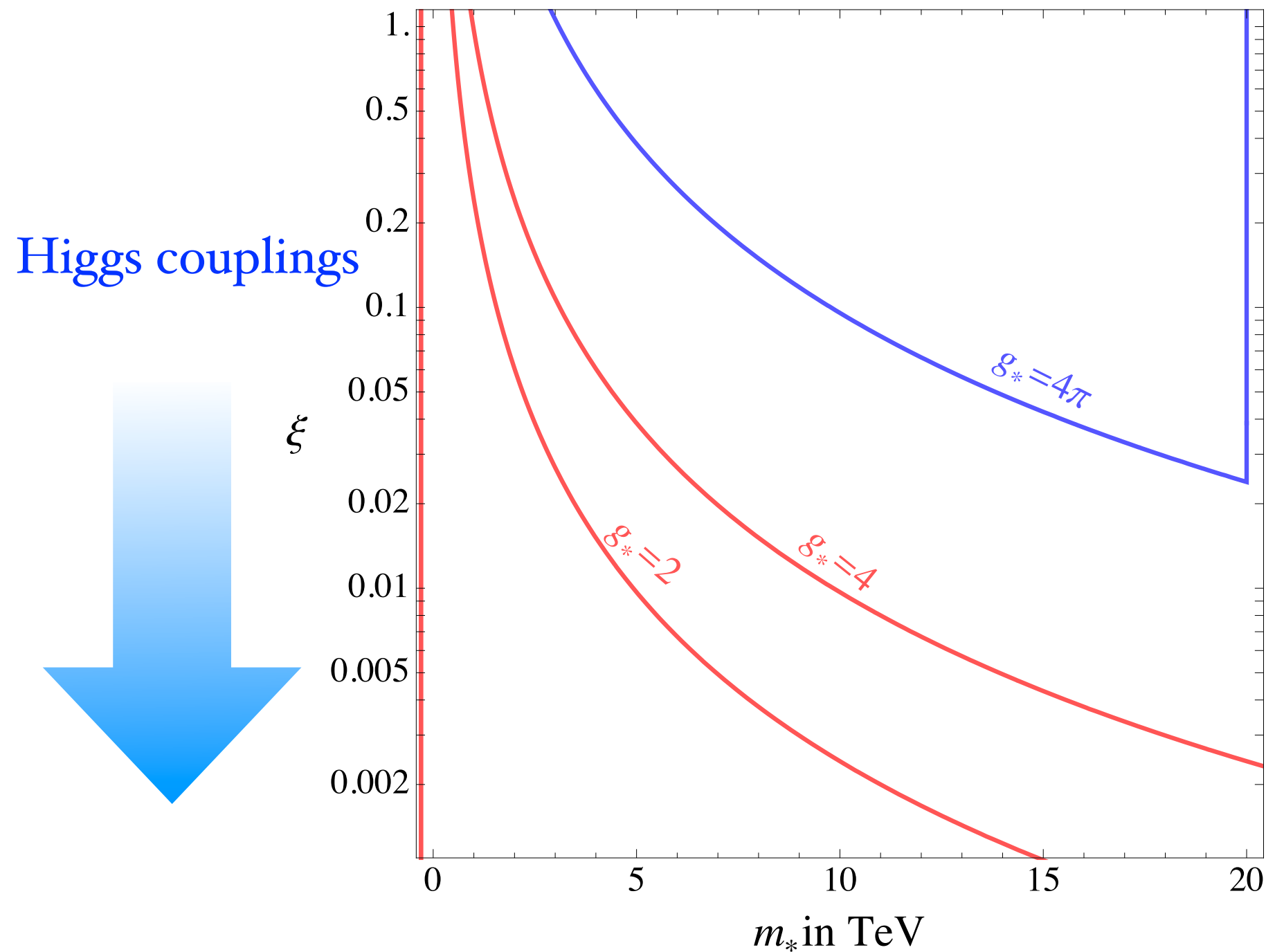
○ Direct searches for resonances: $m_\rho \approx g_* f$

Which one is doing best?
it depends on value of g_*

Conclusions: Higgs & New Physics

Precision /indirect searches (high lumi.) vs. direct searches (high energy)

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Rattazzi, BSM@100TeV, CERN '14

direct searches

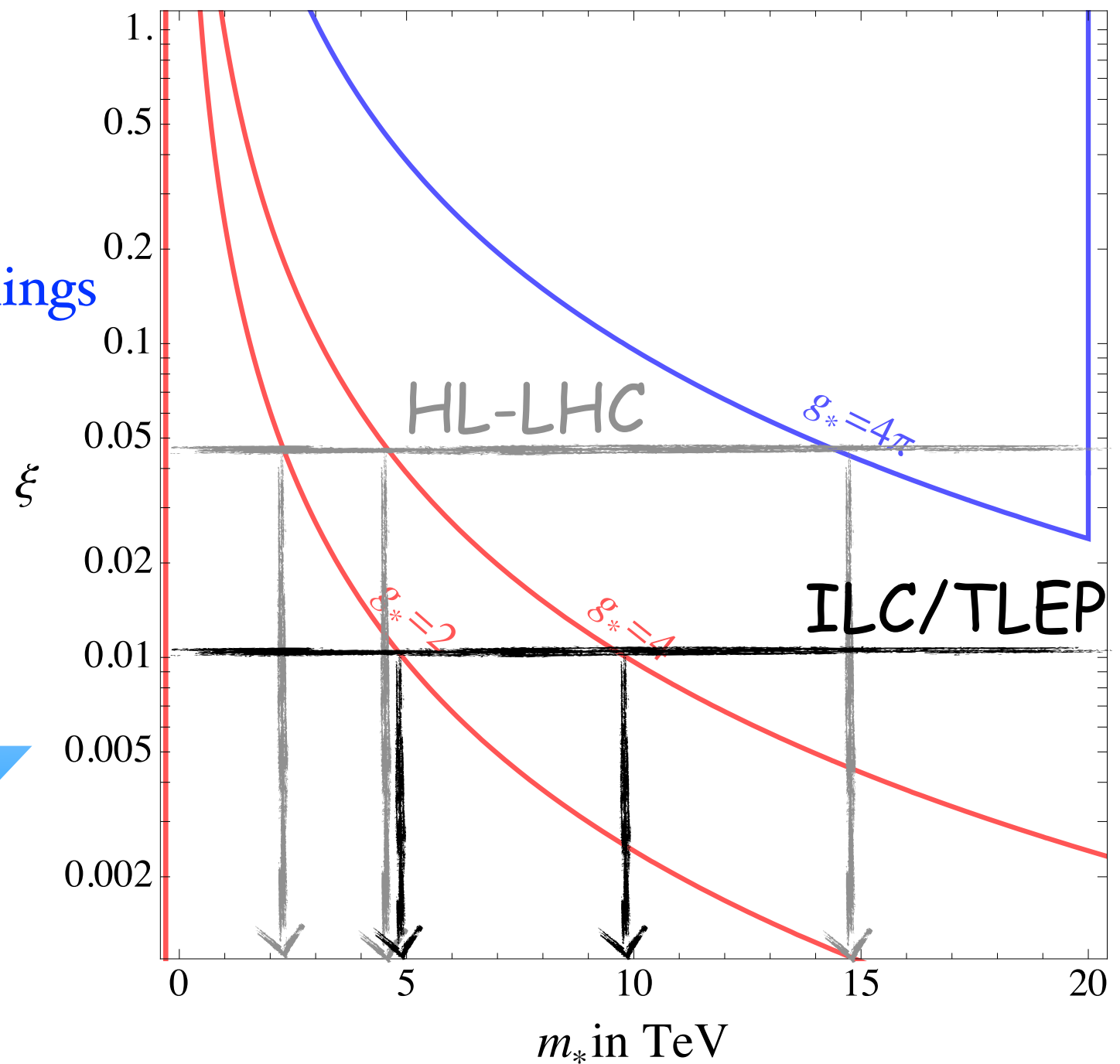
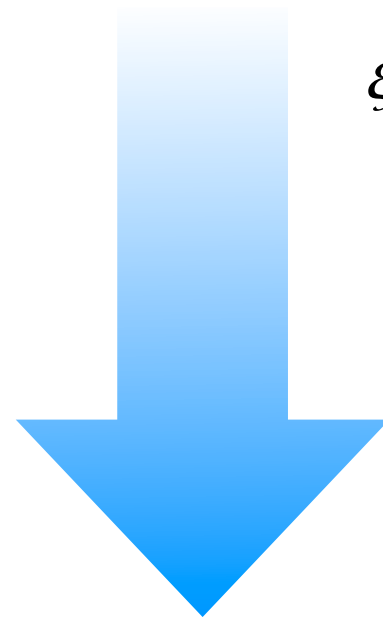
Conclusions: Higgs & New Physics

Precision /indirect searches (high lumi.) vs. direct searches (high energy)

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Higgs couplings

► nice complementarity between direct searches and precision Higgs physics



direct searches



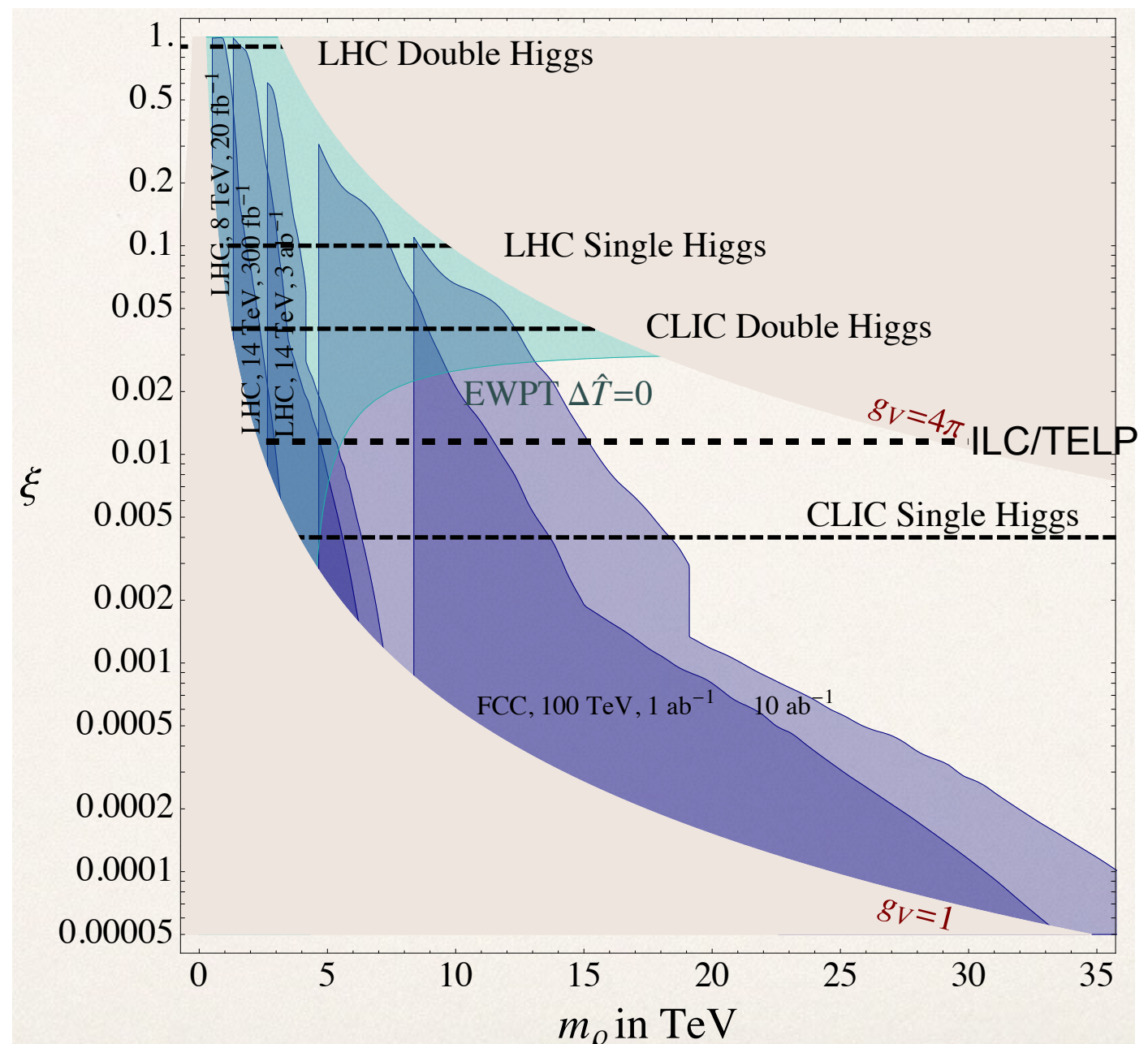
Rattazzi, BSM@100TeV, CERN '14

Conclusions: Higgs & New Physics

Precision /indirect searches (high lumi.) vs. direct searches (high energy)

► large region of parameter space already disfavored by EW precision data

► complementarity between direct searches @ hadron machine and indirect higgs measurements @ lepton machine



Contino, Grojean, Pappadopulo, Rattazzi, Thamm '13

Torre, Thamm, Wulzer '14

a deviation in Higgs couplings also teaches us on the maximum mass scale to search for!
e.g. 10% deviation $\Rightarrow m_V < 10\text{TeV}$ i.e. resonance within the reach of FCC-hh