

The classically conformal B-L extended Ma model

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The Classically conformal B-L extended Ma model solves 2 hierarchy problems

① Neutrino mass scale and EW scale

② EW(TeV) scale and Planck scale

Ma model

E. Ma, Phys.Rev.D73, 077301 (2006)

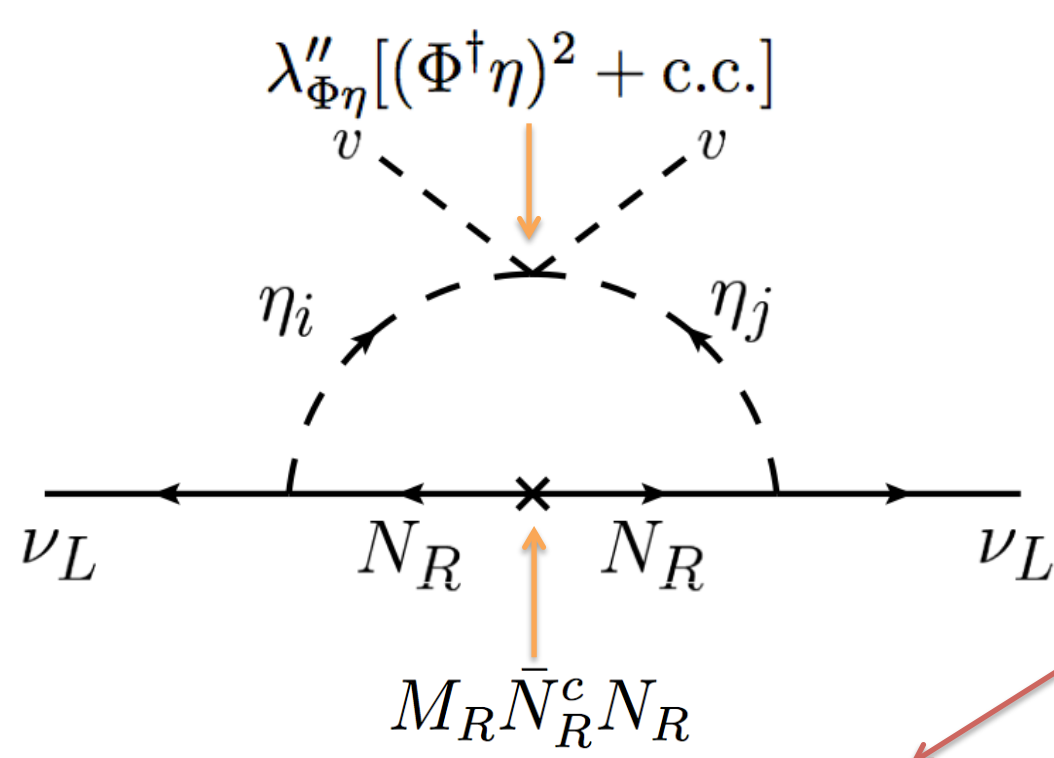
Ma model is minimal radiative see-saw model with DM

Fermion	L_L	e_R	N_R
$(SU(2)_L, U(1)_Y)$	$(2, -1/2)$	$(1, -1)$	$(1, 0)$
Z_2	+	+	-

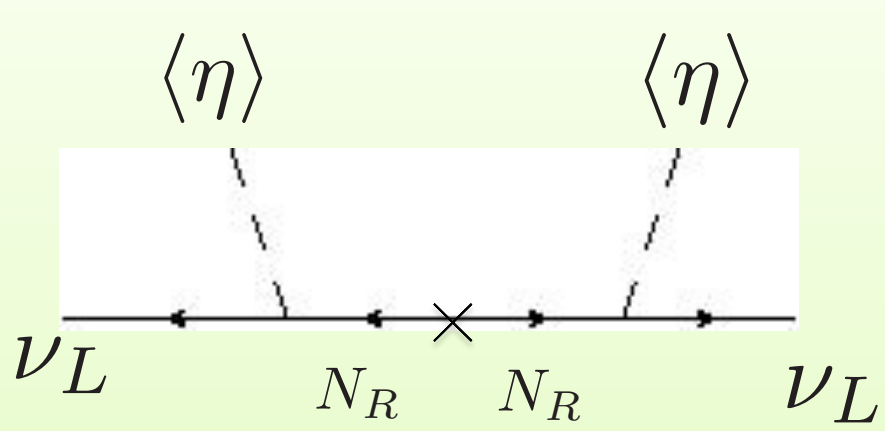
Boson	Φ	η
$(SU(2)_L, U(1)_Y)$	$(2, 1/2)$	$(2, 1/2)$
Z_2	+	-

If $\lambda''_{\Phi\eta}$ is zero, neutrino masses are zero

$$(\mathcal{M}_\nu)_{ab} = \frac{(y_\eta)_{ak}(y_\eta)_{bk}M_k}{(4\pi)^2} \left[\frac{m_R^2}{m_R^2 - M_k^2} \ln \frac{m_R^2}{M_k^2} - \frac{m_I^2}{m_I^2 - M_k^2} \ln \frac{m_I^2}{M_k^2} \right] \quad m_R^2 - m_I^2 = 2\lambda''_{\Phi\eta}v^2$$



If Majorana mass is forbidden, this diagram can't draw



If η has non-zero VEV, neutrino mass is generated by tree level

M2, C2

Yukawa interaction

$$-\mathcal{L}_Y = (y_\ell)_a \bar{L}_{La} \Phi e_{Ra} + (y_\eta)_a \bar{L}_{La} \eta^* N_{Ra} + \frac{1}{2} y_N \varphi \bar{N}_R^c N_R + \text{h.c.}$$

Potential

$$\mathcal{V} = \lambda_\Phi |\Phi|^4 + \lambda_\eta |\eta|^4 + \lambda_\varphi |\varphi|^4 + \lambda_{\Phi\eta} |\Phi|^2 |\eta|^2 + \lambda'_{\Phi\eta} |\Phi^\dagger \eta|^2 + \lambda''_{\Phi\eta} [(\Phi^\dagger \eta)^2 + \text{c.c.}] + \lambda_{\Phi\varphi} |\Phi|^2 |\varphi|^2 + \lambda_{\eta\varphi} |\eta|^2 |\varphi|^2,$$

Assumptions

1. These couplings are zero at Planck scale for simplicity
2. There are no intermediate scale

M1. Majorana mass

M2. $\lambda''_{\Phi\eta}$ is non-zero

M3. η is inert doublet

Classically conformal model

W.A. Bardeen, FERMILAB-CONF-95-391-T

Classically conformal theory with no intermediate scale can be an alternative solution to the naturalness problem

Hierarchy problem

$$m^2(\mu) = m_0^2 + \frac{3\Lambda^2}{8\pi^2 v^2} (m_Z^2 + 2m_W^2 + m^2 - 4m_t^2) + \mathcal{O}\left(\ln \frac{\Lambda}{\mu}\right) \quad \text{subtracted at UV scale}$$

Once subtracted, no longer appears

Classically conformal invariance

Natural boundary condition is no mass terms at Planck scale

$$m^2(\Lambda) = m_0^2 + \frac{3\Lambda^2}{8\pi^2 v^2} (m_Z^2 + 2m_W^2 + m^2 - 4m_t^2) = 0$$

C1. No mass term

C2. No intermediate scale

C3. CW type EWSB

If large new scale exists,

Higgs mass has large correction.

$$\delta m^2 \sim \lambda_{mix} m_{new}^2$$

Effective potential

$$V_{eff}(\phi) = \frac{\lambda_\Phi(\phi)}{4} \phi^4$$

$$\textcircled{1} \quad \frac{dV_{eff}}{d\phi} = \frac{M^4}{4} \left(\frac{d\lambda_\Phi}{dt} + 4\lambda_\Phi \right) = 0$$

$$\lambda_\Phi = -\frac{1}{4} \frac{d\lambda_\Phi}{dt} \sim -\frac{1}{64\pi^2} (96g^4 - Y_N^4)$$

λ_Φ is generated by loop effect

$$\textcircled{2} \quad \frac{d^2 V_{eff}}{d\phi^2} = -4\lambda_\Phi M^2 > 0 \quad \lambda_\Phi < 0$$

λ_Φ is negative

M1, C1

Majorana mass term is forbidden by Classically conformal invariance

B-L gauged extension

Majorana mass is generated by B-L symmetry breaking

$$\frac{1}{2} y_N \varphi \bar{N}_R^c N_R \rightarrow M_R \bar{N}_R^c N_R$$

Fermion	L_L	e_R	N_R
$(SU(2)_L, U(1)_Y)$	$(2, -1/2)$	$(1, -1)$	$(1, 0)$
$U(1)_{B-L}$	-1	-1	-1
Z_2	+	+	-

DM candidate

Boson	Φ	η	φ
$(SU(2)_L, U(1)_Y)$	$(2, 1/2)$	$(2, 1/2)$	$(1, 0)$
$U(1)_{B-L}$	0	0	2
Z_2	+	-	+

Inert doublet

C3

If B-L symmetry is broken, the SM Higgs doublet mass is generated through the mixing term in the scalar potential.

$$V \sim \lambda_\Phi \Phi^4 + \lambda_{\Phi\varphi} \Phi^2 \varphi^2$$

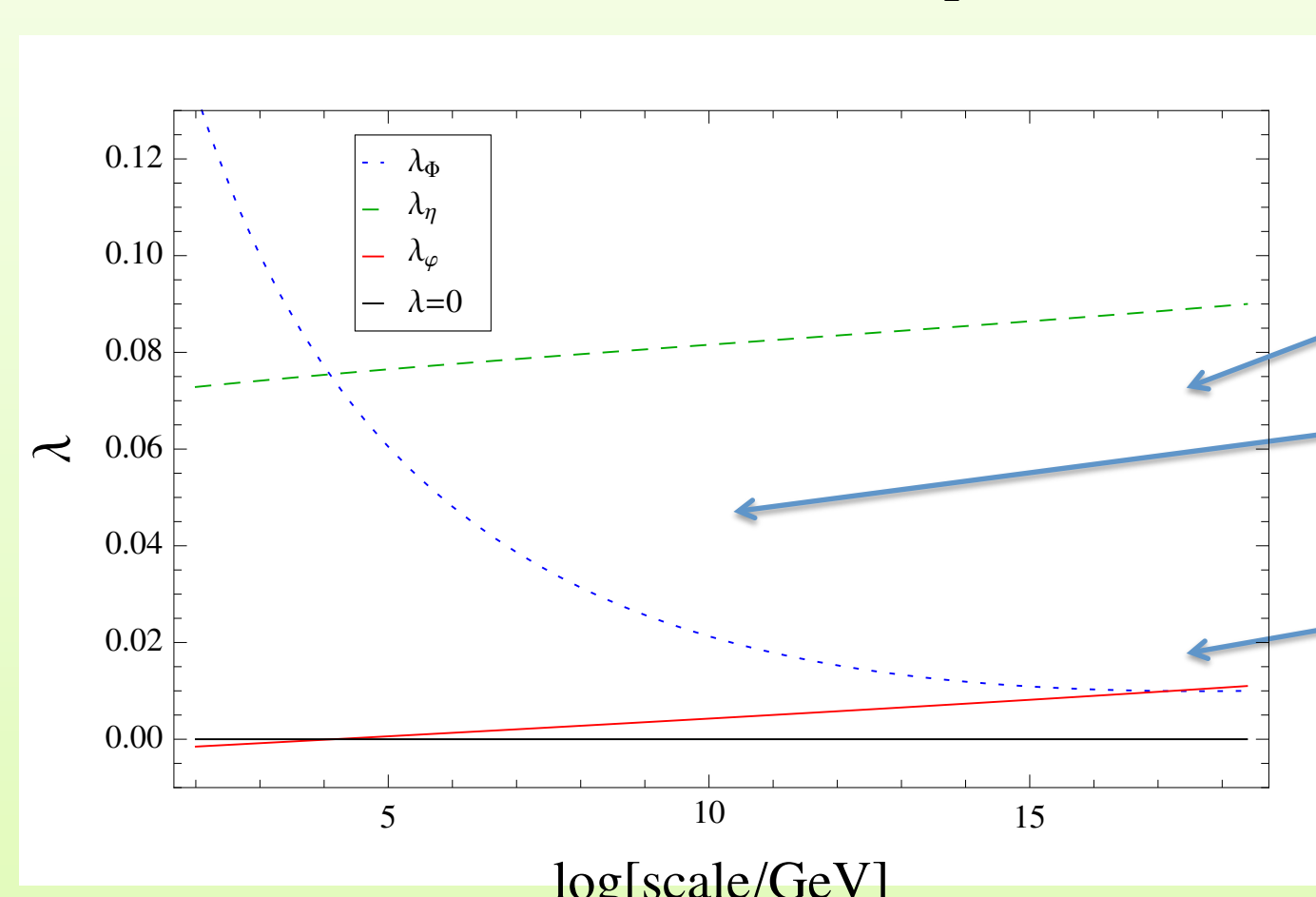
Φ has VEV.

$$V \sim \lambda_\Phi \Phi^4 + \lambda_{\Phi\varphi} \Phi^2 v_{B-L}^2$$

M3

$$g_{B-L} = 0.17 \quad Y_N = 0.2$$

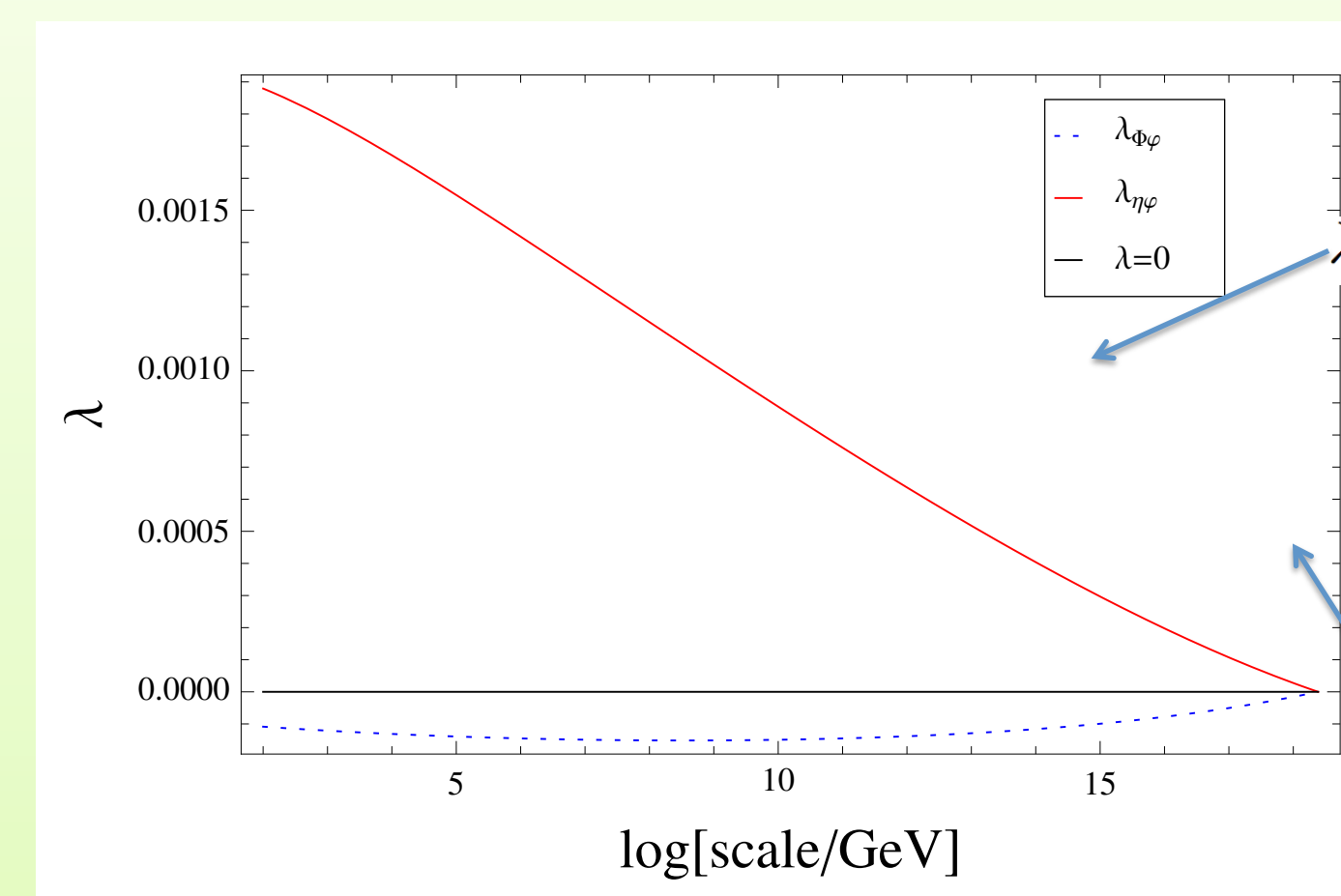
$$\lambda''_{\Phi\eta} = 10^{-9} \quad m_Z = 3.7 \text{ TeV}$$



λ_η should be positive for stability

$$\lambda_\Phi = -\frac{1}{4} \frac{d\lambda_\Phi}{dt} \sim -\frac{1}{64\pi^2} (96g^4 - Y_N^4)$$

$$\lambda_\Phi < 0$$



η mass term $\lambda_{\eta\varphi} v_{B-L}^2 \eta^2$

η is inert doublet

$$\lambda_{\eta\varphi} > 0$$

Higgs mass term

$$\lambda_{\Phi\varphi} v_{B-L}^2 \Phi^2$$

$$\lambda_{\Phi\varphi} < 0$$

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