

Mixed Light Sneutrinos as Thermal Dark Matter

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in collaboration with

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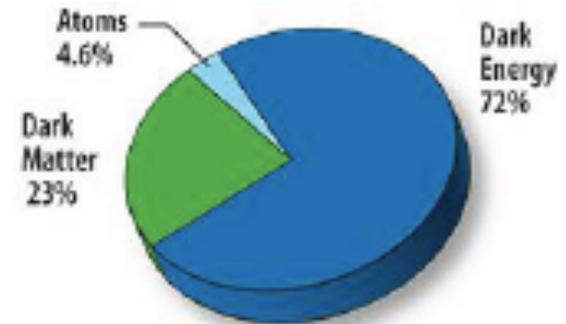
February 20, 2012, Toyama University

Outline

1. Motivation
2. Framework
3. Constraints on the model
4. Results in parameter space
 - 4-1. Direct detection
 - 4-2. Indirect detection
5. Collider signatures
6. Conclusions

1. Motivation

- Neutrino oscillations:
 - Nonzero neutrino masses are required
 - Neutrinos are massless in the SM
- Non-baryonic dark matter (DM)
 - $\Omega_{DM}h^2 = 0.1109 \pm 0.0059$
 - No candidate particle in the SM
- In Supersymmetric (SUSY) models with Dirac neutrinos
 - Right Handed (RH) neutrinos \rightarrow RH-sneutrino superpartners
 - Non-negligible $A_{\tilde{\nu}}$ induces sizable L-R mixing in sneutrino sector
 - \rightarrow Light sneutrino $\tilde{\nu}_1$ is the LSP
 - \rightarrow Relic density is in desirable range



[<http://map.gsfc.nasa.gov>]

[Arkani-hamed,Hall,Murayama,Weiner(2001);
Arina,Fornengo(2007);Thomas,Tucker-Smith,Weiner(2008)]

2. Framework

- MSSM is now extended with only Dirac masses for sneutrinos

$$\Delta\mathcal{L}_{\text{soft}} = m_{\tilde{N}_i}^2 |\tilde{N}_i|^2 + A_{\tilde{\nu}_i} \tilde{L}_i \tilde{N}_i H_u + \text{h.c.}$$

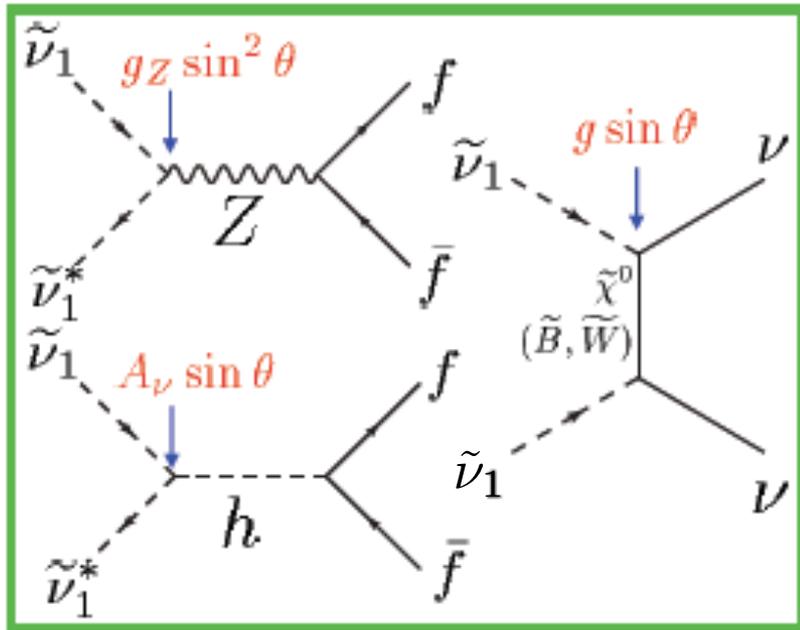
- Sneutrino mass matrix

$$\mathcal{M}_{\tilde{\nu}}^2 = \begin{pmatrix} m_{\tilde{L}}^2 + \frac{1}{2}m_Z^2 \cos 2\beta & \frac{1}{\sqrt{2}}A_{\tilde{\nu}} v \sin \beta \\ \frac{1}{\sqrt{2}}A_{\tilde{\nu}} v \sin \beta & m_{\tilde{N}}^2 \end{pmatrix}$$

$$\begin{aligned} \tilde{\nu}_1 &= \cos \theta_{\tilde{\nu}} \tilde{\nu}_R - \sin \theta_{\tilde{\nu}} \tilde{\nu}_L, & \theta_{\tilde{\nu}} &= \frac{1}{2} \sin^{-1} \left(\frac{\sqrt{2}A_{\tilde{\nu}} v \sin \beta}{m_{\tilde{\nu}_2}^2 - m_{\tilde{\nu}_1}^2} \right) \\ \tilde{\nu}_2 &= \sin \theta_{\tilde{\nu}} \tilde{\nu}_R + \cos \theta_{\tilde{\nu}} \tilde{\nu}_L, & & \text{where } m_{\tilde{\nu}_1} < m_{\tilde{\nu}_2} \end{aligned}$$

- $A_{\tilde{\nu}}$ is not proportional to small neutrino Yukawa couplings
→ not negligible → RH sneutrino mass term by $A_{\tilde{\nu}}$
- Light sneutrino $\tilde{\nu}_1$ is mostly $\tilde{\nu}_R$ by assuming $m_{\tilde{N}} < m_{\tilde{L}}$

Important processes

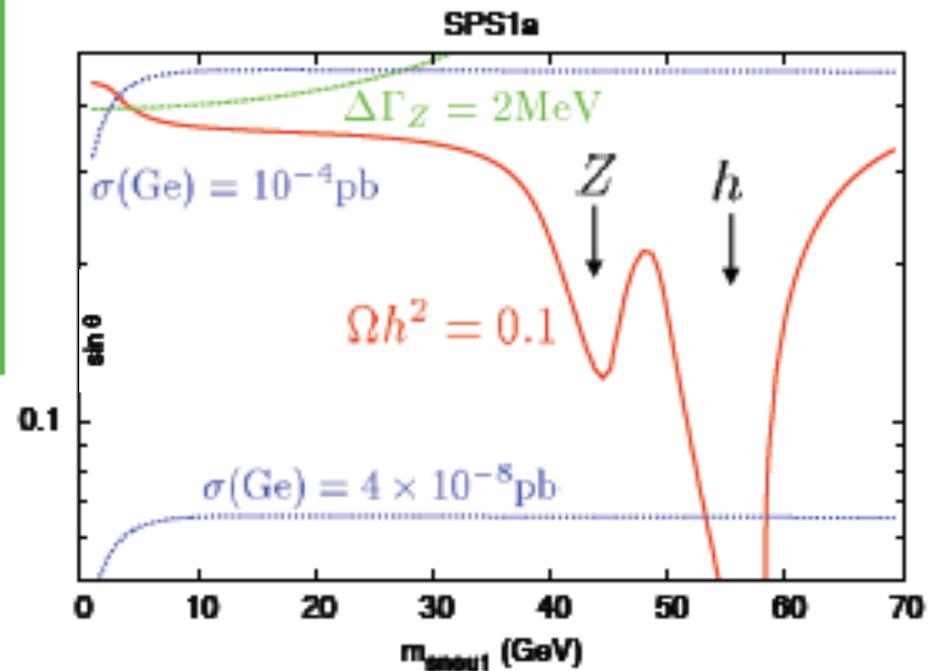


- **Computation:**

LanHEP \rightarrow CalcHEP

\rightarrow micrOMGEAs

[Belanger, Boudjema, Pukhov, Semenov]



Particle spectrum

- Assumptions

- * $M_2 \simeq 2M_1 \simeq M_3/3$ at the weak scale
- * use $m_{\tilde{\nu}_1}$, $m_{\tilde{\nu}_2}$ and $\theta_{\tilde{\nu}}$, as input parameters,
- * compute the $m_{\tilde{L}}$, $m_{\tilde{N}}$ and $A_{\tilde{\nu}}$
- * one-generation case, only the tau-sneutrino is light
- * three-generation case, complete degeneracy between the slepton gen

- Additional radiative corrections to Higgs at 2-loop and sneutrinos at 1-loop

- Main contribution to radiative corrections to sneutrino masses

$$\Delta M_{\tilde{\nu}RR}^2(Q)|_{\text{app}} = \frac{A_{\tilde{\nu}}^2}{8\pi^2} \left(\log \frac{m_{\tilde{\nu}_2}^2}{Q^2} - 1 \right)$$

- Higgs mass \rightarrow negative correction

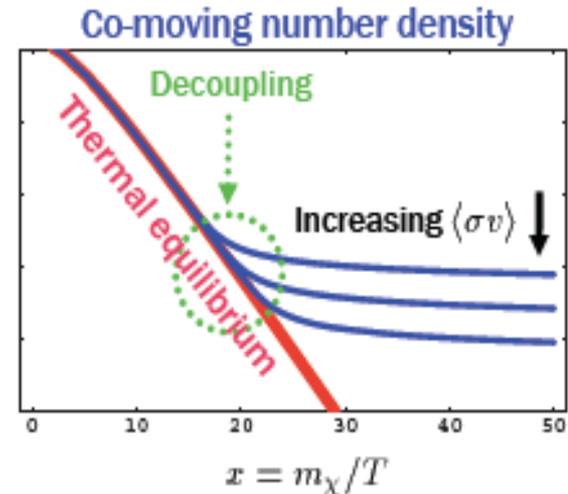
$$\Delta \lambda_2^{(\tilde{\nu})} = -\frac{1}{16\pi^2} \sum_{i=1}^{N_f} \frac{|A_{\nu}|^4}{(m_{\tilde{\nu}_2}^2 - m_{\tilde{\nu}_1}^2)^2} \left(\frac{m_{\tilde{\nu}_2}^2 + m_{\tilde{\nu}_1}^2}{m_{\tilde{\nu}_2}^2 - m_{\tilde{\nu}_1}^2} \log \frac{m_{\tilde{\nu}_2}^2}{m_{\tilde{\nu}_1}^2} - 2 \right)$$

3. Constraints

- Relic abundance of neutrino:

After annihilation rate dropped below expansion rate, number of density per co-moving volume is almost fixed

$$\Omega_\chi h^2 \simeq 0.04 \times \left(\frac{\langle \sigma v \rangle}{1 \text{ pb}} \right)^{-1} \left(\frac{x_F}{22} \right) \left(\frac{g_*}{90} \right)^{-1/2}$$



Standard freeze-out feature

- Collider constrains:

Invisible width of Z-boson with light sneutrino $m_{\tilde{\nu}} < m_Z/2$, and LEP2 constraints

$$\Delta\Gamma_Z = \sum_{i=1}^{N_f} \Gamma_\nu \frac{\sin^4 \theta_{\tilde{\nu}}}{2} \left(1 - \left(\frac{2m_{\tilde{\nu}}}{m_Z} \right)^2 \right)^{3/2} < 2 \text{ MeV}$$

- Direct detection (DD):

Spin-independent DD through Z or Higgs exchange assuming point-like nucleus

$$\sigma_{\tilde{\nu}_1 N}^{\text{SI}} = \frac{4\mu_\chi^2 (Zf_p + (A-Z)f_n)^2}{\pi A^2}$$

For Xenon $A = 131, Z = 54$
while for Germanium $A = 76, Z = 32$

Annihilation channels

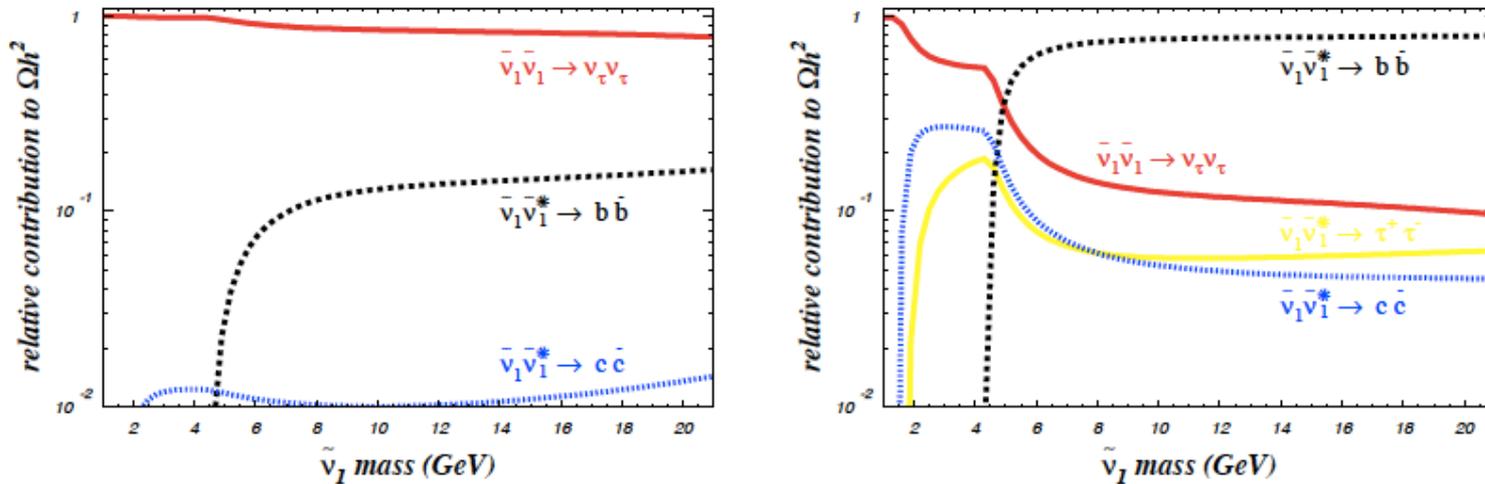
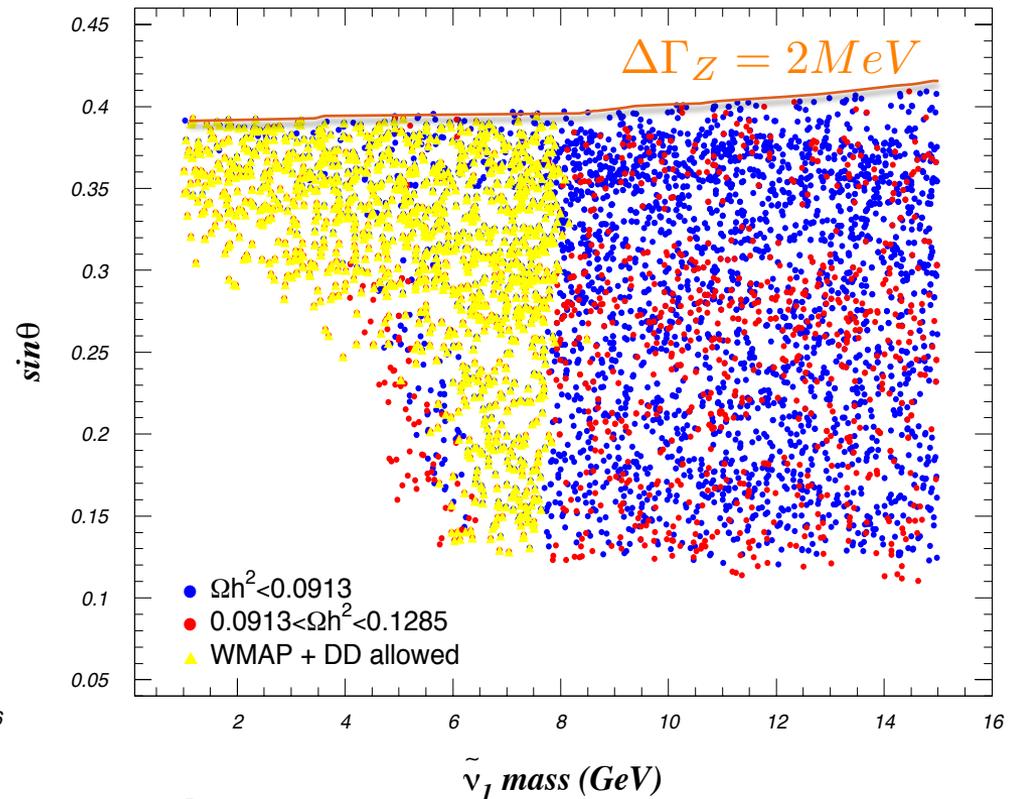
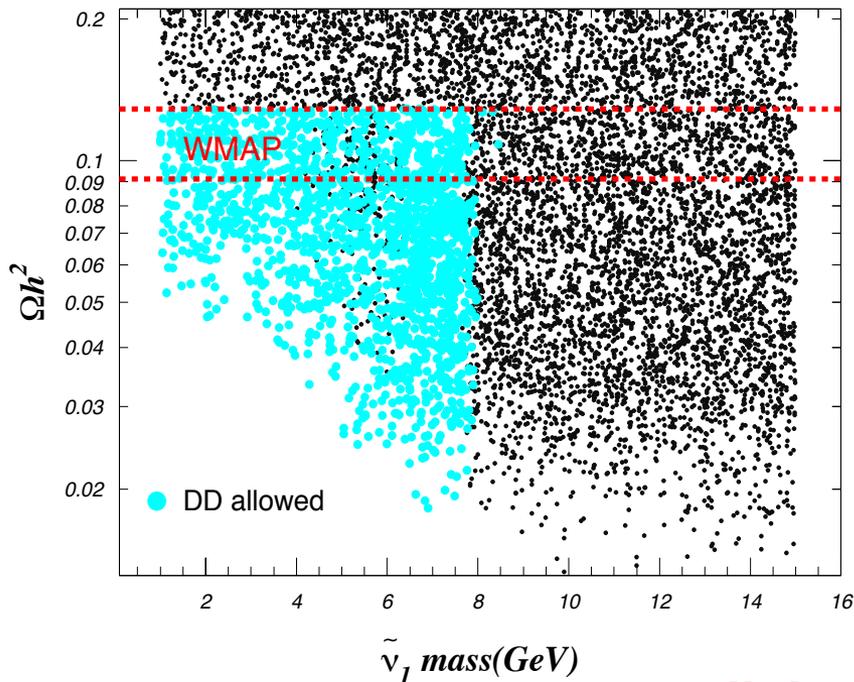


Figure : Relative contributions of different annihilation channels as a function of the $\tilde{\nu}_1$ mass, on the left for $m_{\tilde{\nu}_2} = 200$ GeV and $\sin \theta_{\tilde{\nu}} = 0.35$, on the right for $m_{\tilde{\nu}_2} = 500$ GeV and $\sin \theta_{\tilde{\nu}} = 0.22$

Parameter scan with constraints

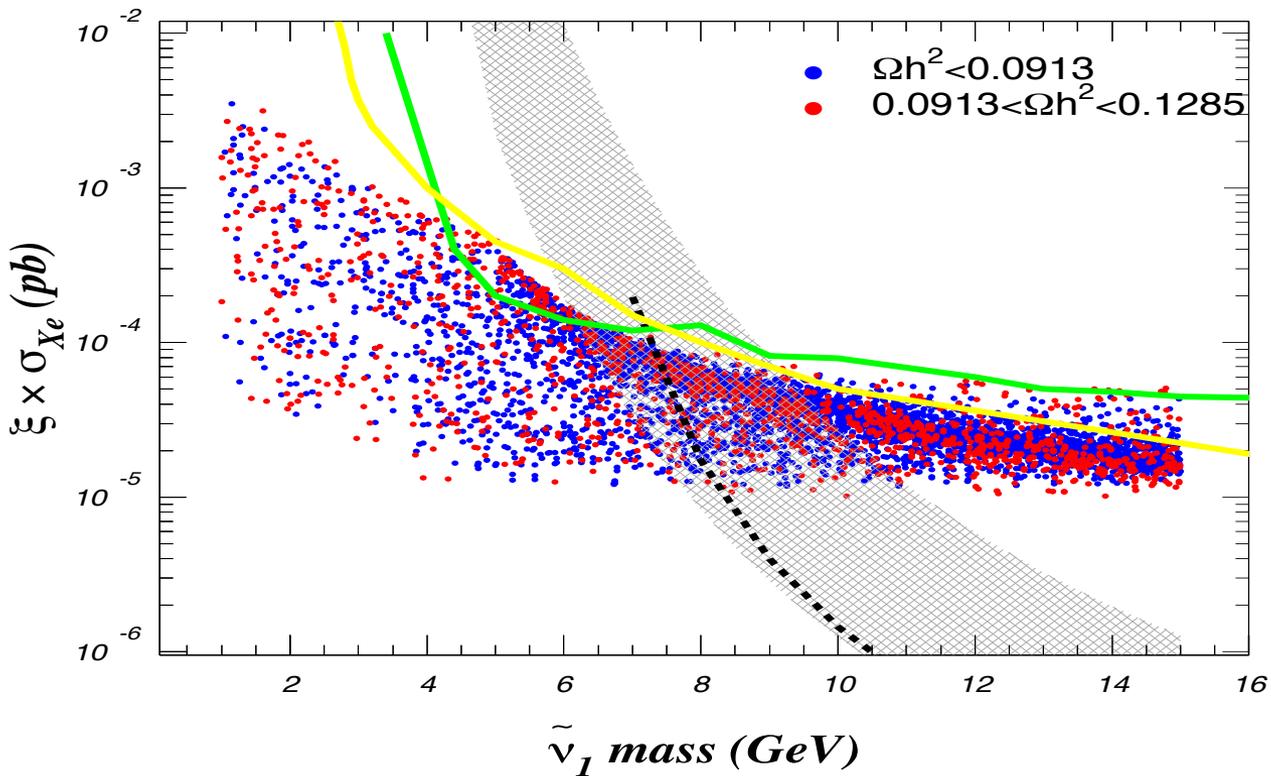
- Apply WMAP bound $\Omega h^2 < 0.1285$ and collider constraints
- Scan over $m_{\tilde{\nu}_1}, m_{\tilde{\nu}_2}, \sin\theta, M_2$



one light sneutrino case

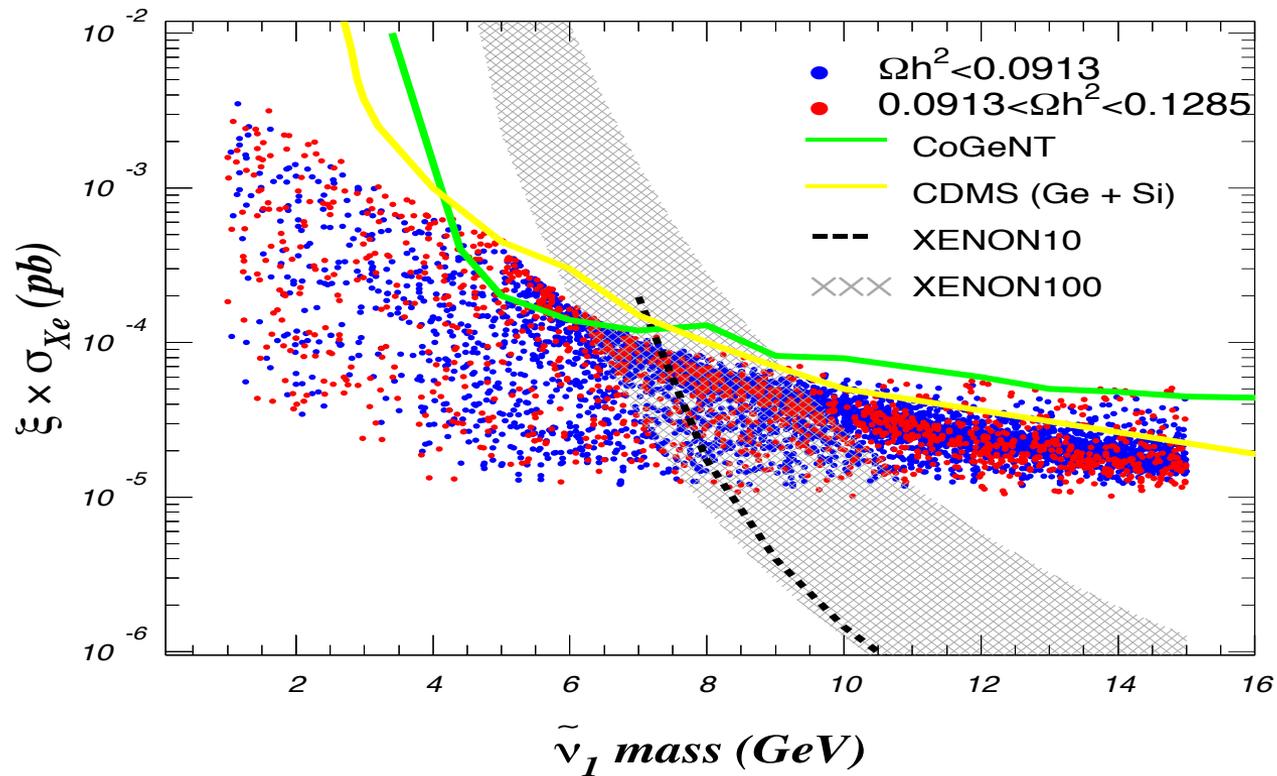
4-1. Direct detection one light sneutrino case

- Apply WMAP bound $\Omega h^2 < 0.1285$ and collider constraints $\Delta\Gamma_Z = 2MeV$, LEP2



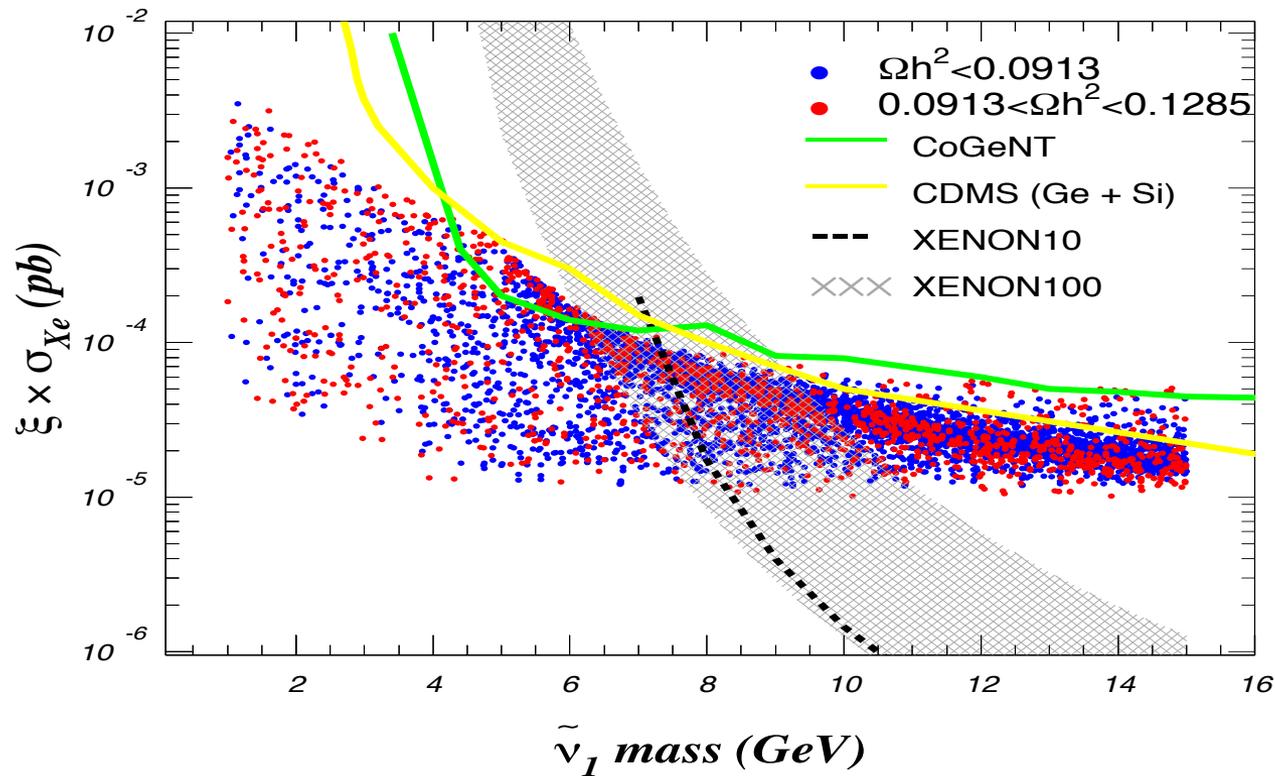
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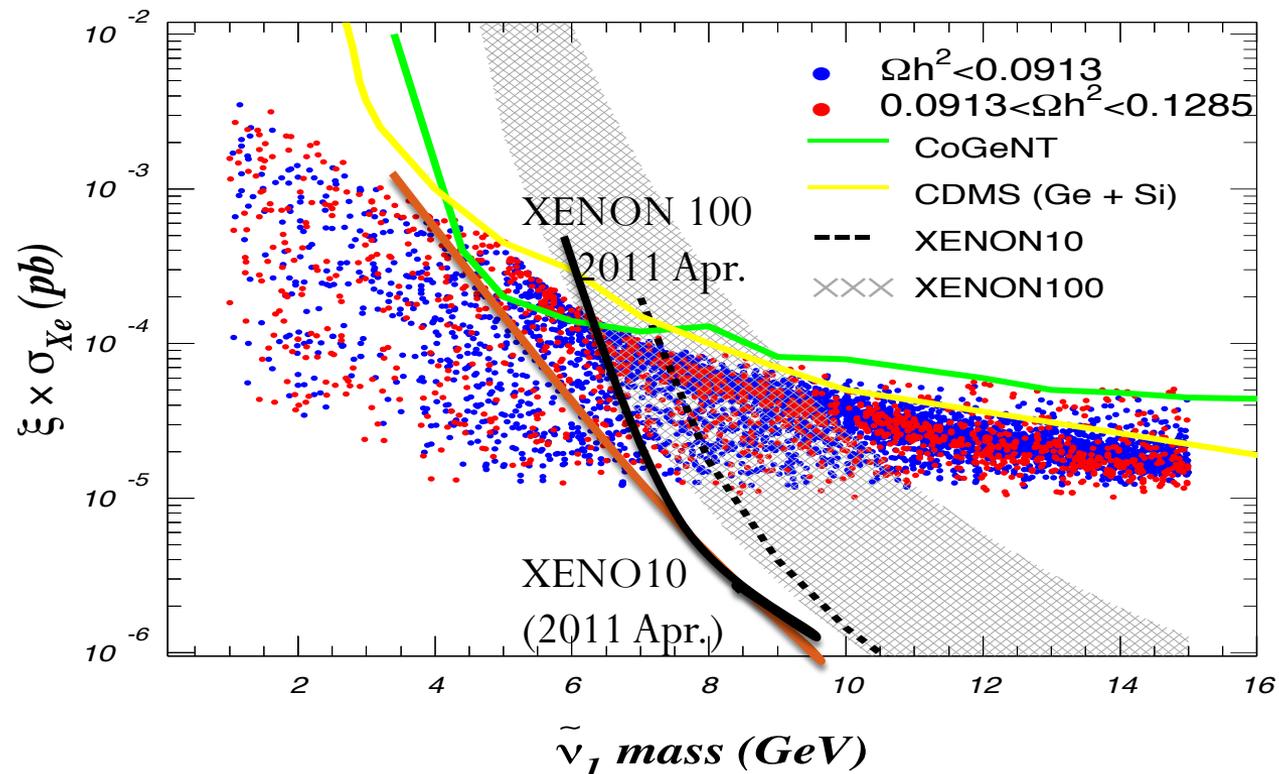
4-1. Direct detection one light sneutrino case

- Apply WMAP bound $\Omega h^2 < 0.1285$ and collider constraints $\Delta\Gamma_Z = 2MeV$
- $\xi = \Omega h^2 / 0.11$ for $\Omega h^2 < 0.0913$, and 1 otherwise
- $v_0 = 220 \pm 20 \text{ km/s}$, $500 < v_{esc} < 600 \text{ km/s}$, $\rho = 0.3 \text{ GeV/cm}^3$



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- DD exp. favored regions below 8 GeV



4-1. Direct detection three degenerate light sneutrinos

- Relic density can increase by factor up to 3: $\langle\sigma v\rangle \propto (\sum g_i^2 \sigma_{ij}) / \sum g_i^2$

- Z- width on mixing angle is stronger

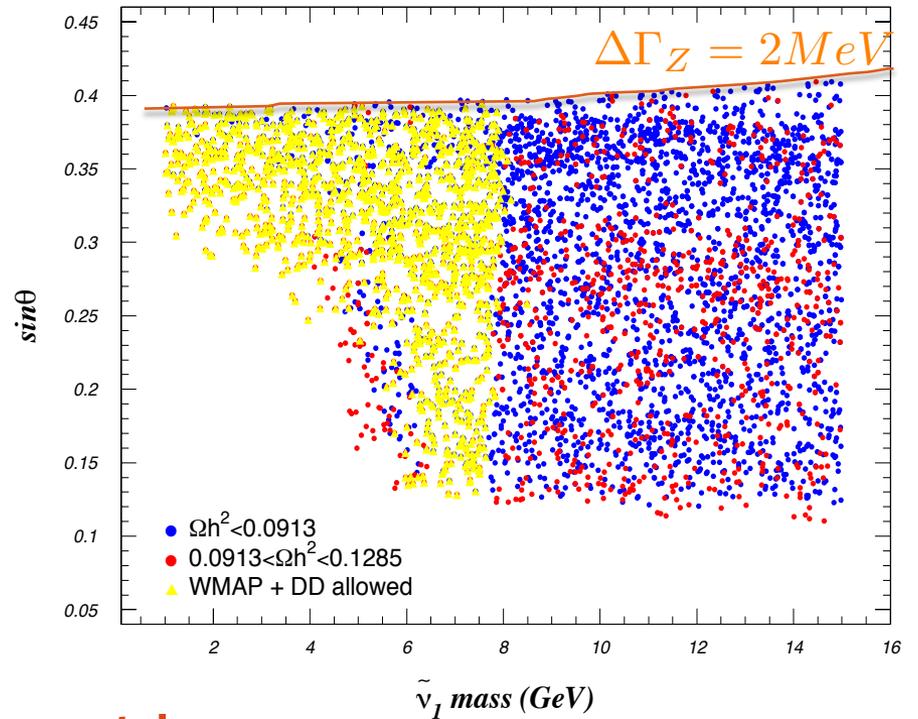
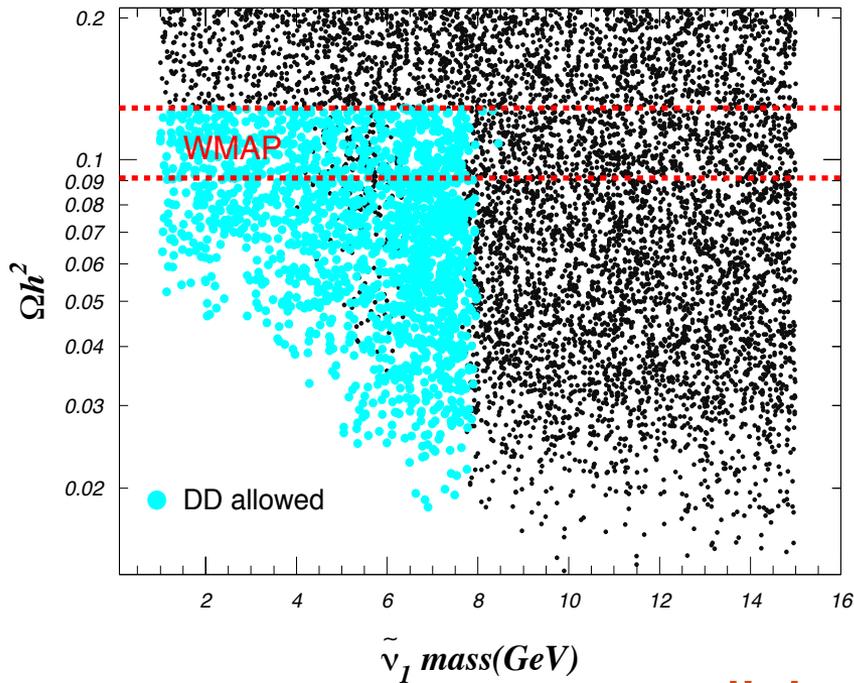
$$\Delta\Gamma_Z = \sum_{i=1}^{N_f} \Gamma_\nu \frac{\sin^4 \theta_{\bar{\nu}}}{2} \left(1 - \left(\frac{2m_{\bar{\nu}}}{m_Z} \right)^2 \right)^{3/2} < 2 \text{ MeV}$$

⇒ Tension between these two facts above

three light sneutrinos case

Relic density increases up to 3 times

Z-width constrains mixing angle upper bound to be smaller

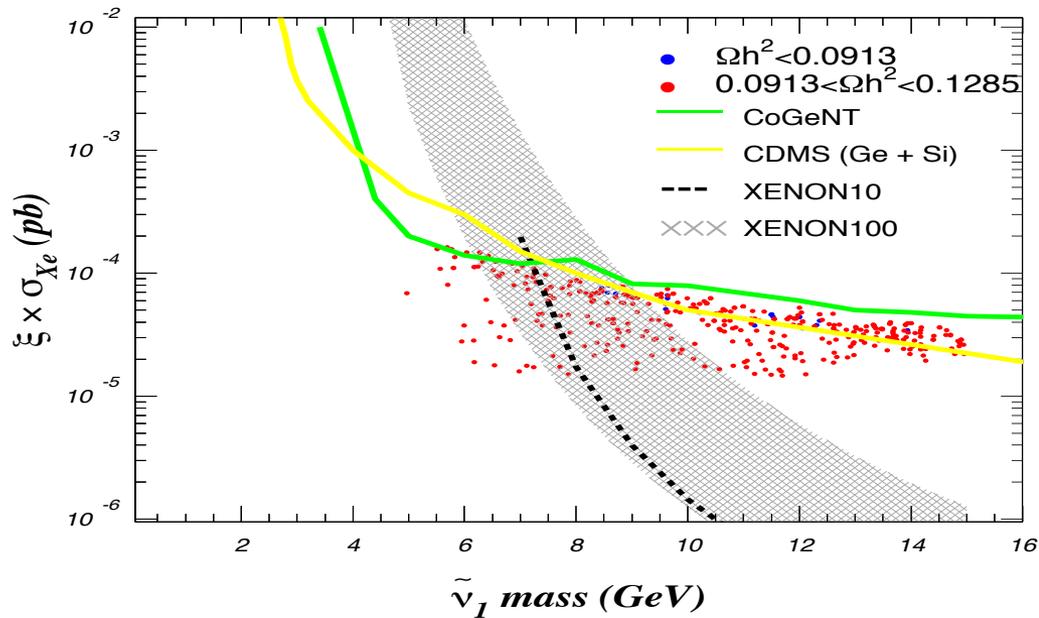


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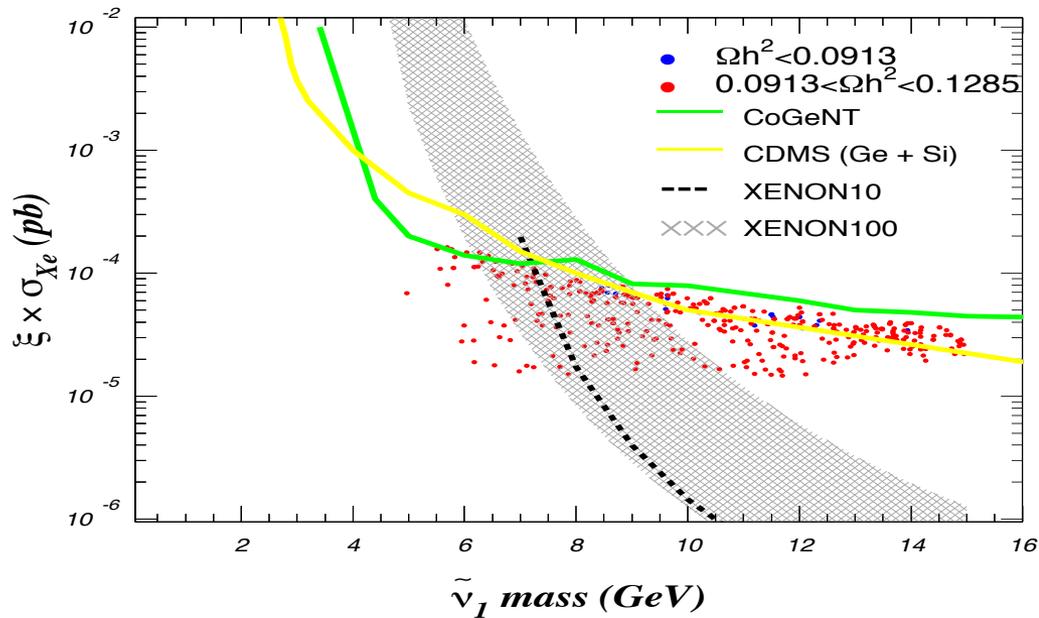
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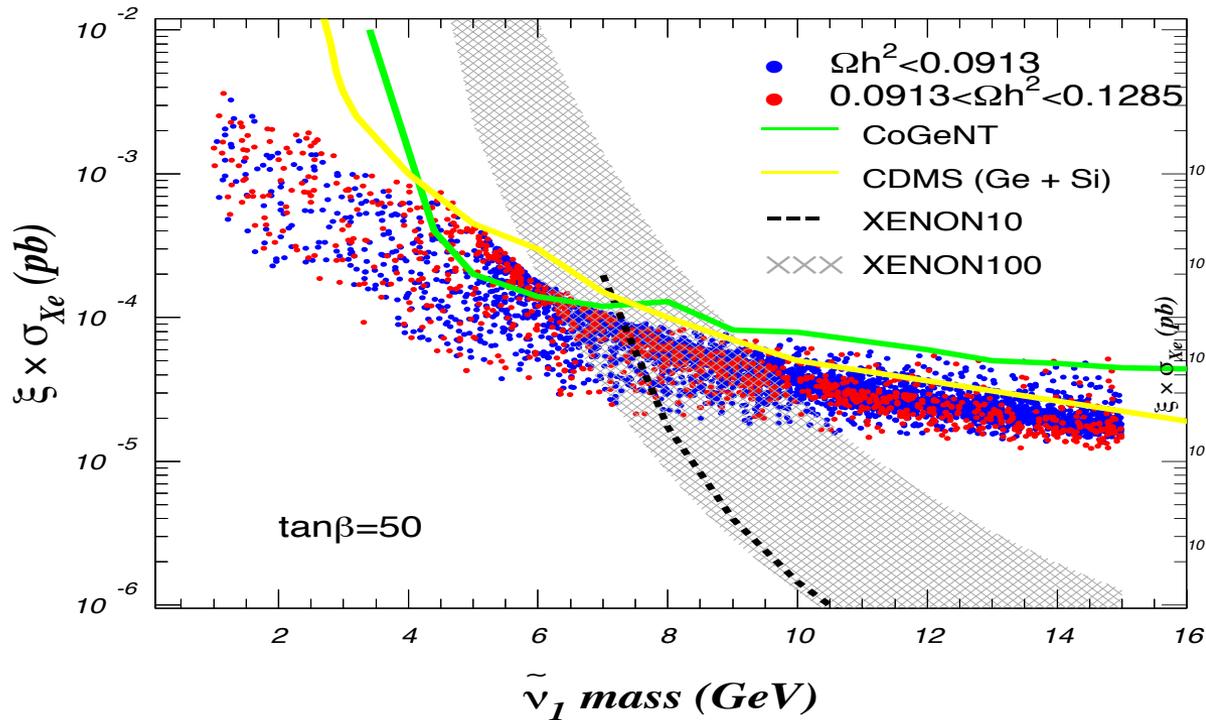
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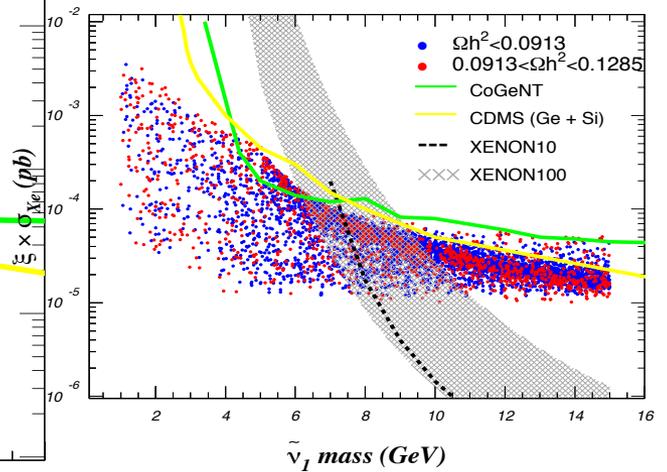
- Mass splitting between different flavours of sneutrinos can be induced by small mass splittings in soft terms → back to the case of one light sneutrino

4-1. Direct detection one light sneutrino with $\tan\beta = 50$

* $\tan\beta = 50$: low $m_{\tilde{\nu}_2}$ removed



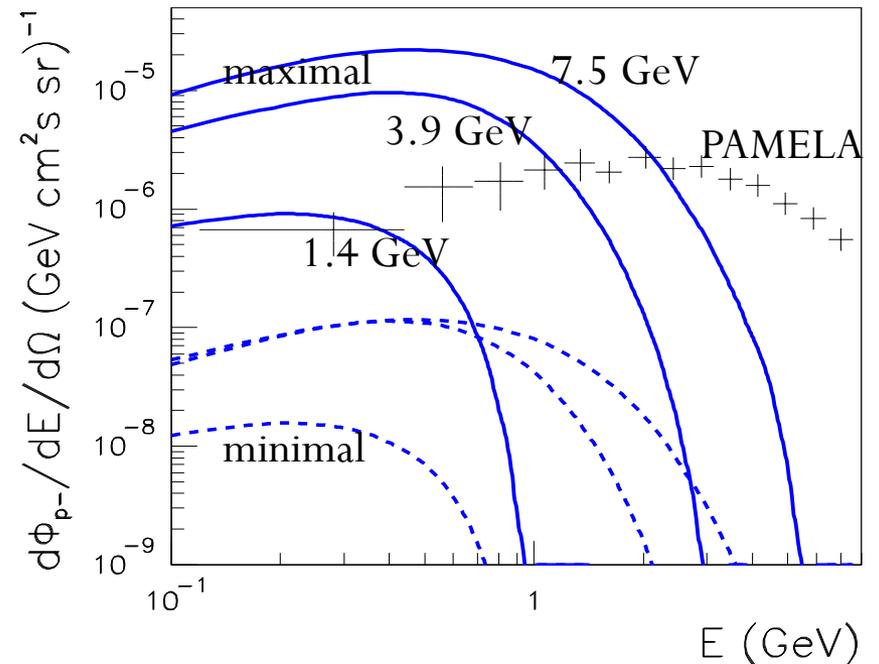
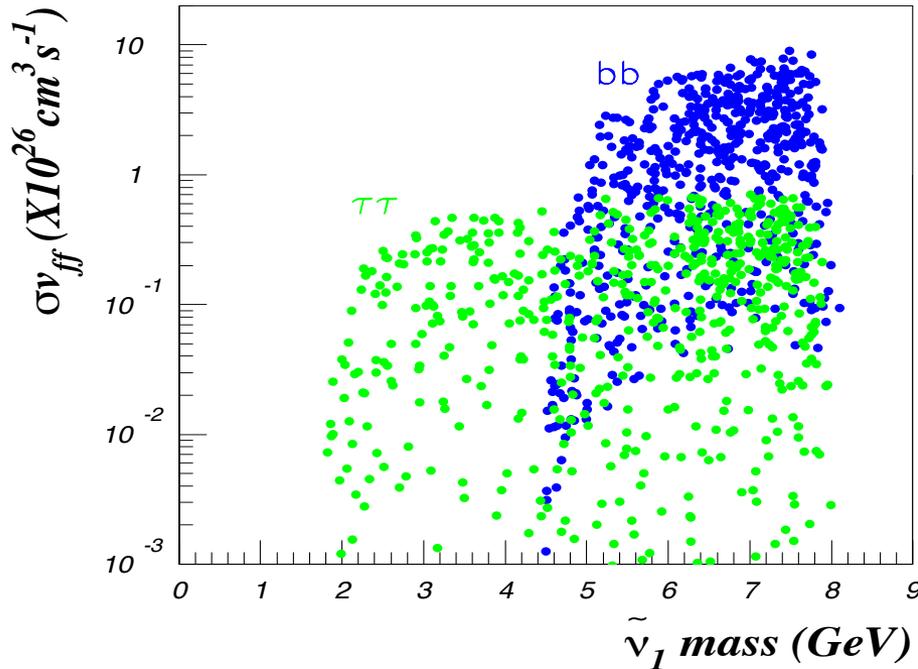
$\tan\beta = 10$



4-2. Indirect detection

Photons, Antiprotons and Positrons

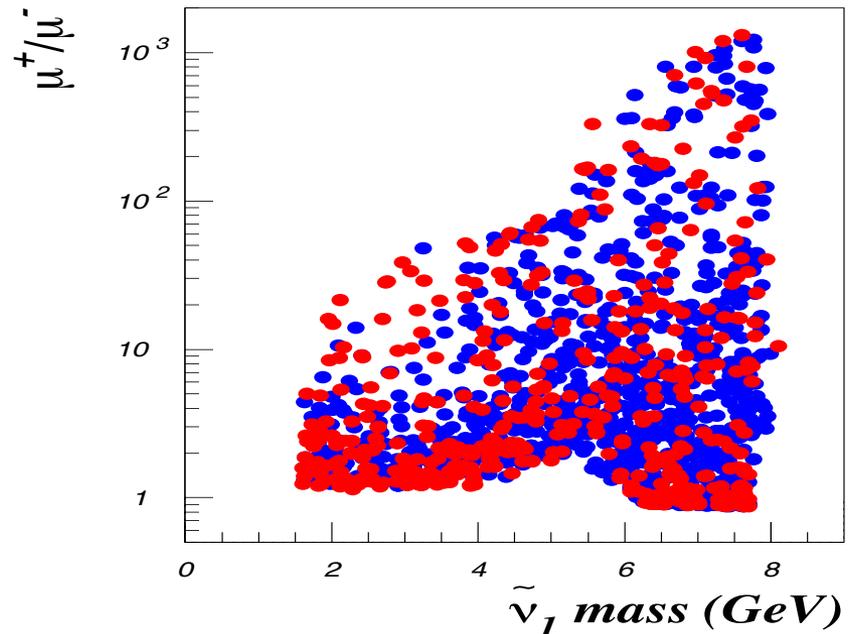
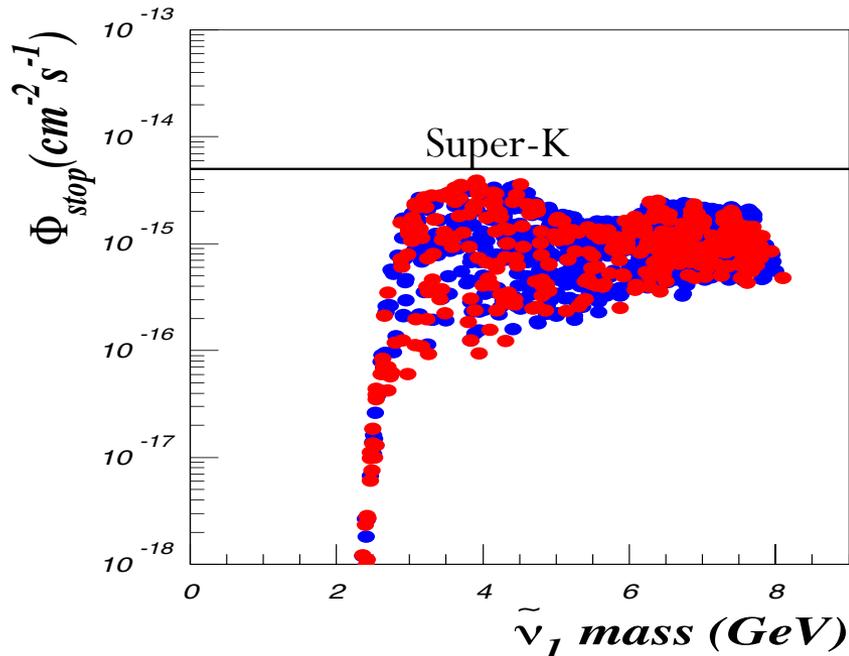
- Annihilation channels into charged fermions leave signature in photons
- NFW DM profile
- Fermi-LAT can probe photon signals
- Room for dark matter contribution to antiproton, positron fluxes detected by PAMELA



4-2. Indirect detection

Neutrinos from annihilation in the Sun

- Stopped muon flux detectable at Super-K with threshold 1.6 GeV
- Antimuon flux is larger than muon flux in our model
→ additional test

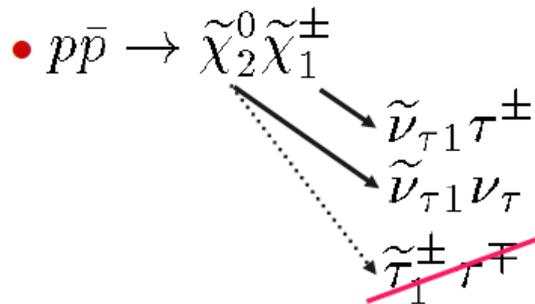


5. Collider signatures

LHC

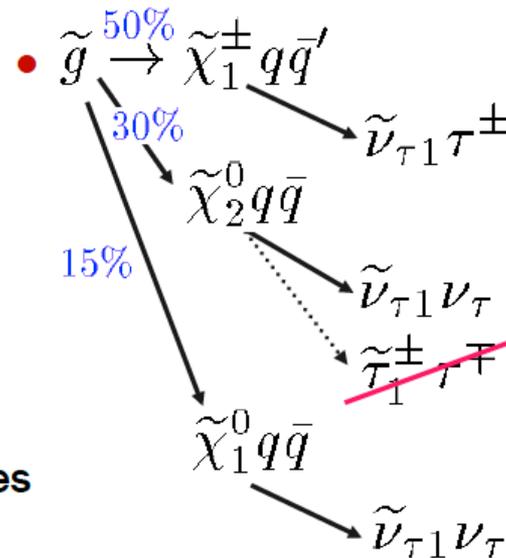
- To reconcile GeV sneutrino scenario with DD exp.constraints, lighter gauginos and heavier sleptons are desirable

- Lepton modes**



- Higgs decay into invisible modes**

$\text{Br}(h \rightarrow \tilde{\nu}_1 \tilde{\nu}_1^*) > 90\%$ in most cases



5. Collider signatures

ILC $\sqrt{s} = 500$ GeV

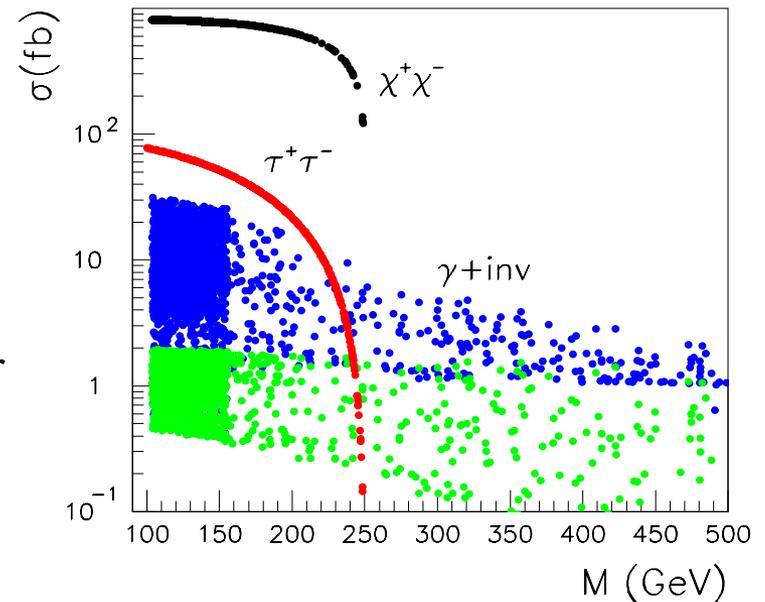
- Main SUSY productions

$$\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp, \tilde{\chi}_1^0 \tilde{\chi}_2^0, \tilde{\chi}_2^0 \tilde{\chi}_2^0, \tilde{\nu}_1 \tilde{\nu}_1^*, \tilde{\tau}_1^+ \tilde{\tau}_1^-$$

- $e^+e^- \rightarrow Zh$ is main production for light Higgs

- $\text{Br}(h \rightarrow \tilde{\nu}\tilde{\nu}^*) \sim 99\%$

- Single photon production from $\tilde{\nu}_1 \tilde{\nu}_1^*$ (dominant) and $\tilde{\chi}_i^0 \tilde{\chi}_j^0$



$$e^+e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \quad M = m_{\tilde{\chi}_1^\pm}$$

$$e^+e^- \rightarrow \tilde{\tau}_1^+ \tilde{\tau}_1^- \quad M = m_{\tilde{\tau}_1}$$

Conclusions

- * In SUSY models with Dirac neutrino mass, $A_{\tilde{\nu}}$ induces a sizable mixing between LH-RH sneutrinos \rightarrow lighter sneutrino dark matter for below ~ 10 GeV
 - * Mixing should be large enough to give efficient pair-annihilation but not be too large in order not to exceed DD limits or Z invisible decay width
- Cross-section for the spin-independent elastic scattering on nuclei is predicted to be $\sigma^{\text{SI}} > 10^{-5}$ pb, an order of magnitude of present limits for DM masses around 5-10 GeV, favored region by CoGeNT (perhaps CRESST)
- * Positron and Antiproton fluxes and Photon flux can be probed by PAMELA and FermiLAT respectively, through sneutrino pair annihilations into $b\bar{b}$ and $\tau^+\tau^-$
 - * Only Super-K has low enough threshold to have sensitivity to neutrinos from light sneutrino annihilation by detecting stopped muons in it
- At LHC light Higgs and neutralinos decay exclusively and lighter chargino give a single lepton while at ILC SUSY production contribute to single photon x-section